

DÉPARTEMENT Sciences de la Planète et de l'Univers



# Characterising turbulence in clusters with Xrays

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#### WHIM and Cluster Outskirts: Lost and Found Baryons in the Local Universe









1

# WHICH PROCESSES AND WHICH SCALES (1)?

#### Injection of turbulent motions

#### Observation of cluster merger



Image credit: X-ray: NASA/CXC/SAO/van Optical: NASA/STScI; Radio: NSF/NRAO/VLA

#### Simulation of cluster merger







Image credit: X-ray: Chandra X-ray observatory

# WHICH PROCESSES AND WHICH SCALES (2)?

3

### Quantify + characterise Turbulence



Kolmogorov Spectrum

 $k_f$ 

 $k_{\eta}$ 



#### Detect



XARM (X-Ray Astronomy Recovery Mission)



Athena Observatory

X-Ray Facilities



# WHICH PROCESSES AND WHICH SCALES (3)?

#### Dissipation of turbulent motions

simplest case

Isotropic Spitzer viscosity + Unmagnetised cluster plasma

Range of dissipation scales  $k_d pprox [10-50] \; {
m kpc}$  (ZuHone et. al 2016)

# But ICM is weakly magnetised

Since  $l_{
m free,ion} \ll r_{
m Larmor}$ 



Viscosity becomes anisotropic (Braginskii 1965)

And

Given an isotropically tangled magnetic field



Viscosity reduced by a factor 5 (Nulsen and McNamara 2013)



Cosmological Context



A cluster is a sphere at Hydrostatic Equilibrium

$$\frac{dP_{th}}{dr} = -\frac{GM(r)\rho(r)}{r^2}$$

In the scope of this approximation

Nowadays

$$M_X = -\frac{k_B T(r)r}{\mu m_p G} \left(\frac{\mathrm{d}ln(\rho(r))}{\mathrm{d}ln(r)} + \frac{\mathrm{d}ln(T(r))}{\mathrm{d}ln(r)}\right)$$

However : The measure mass due to lensing

$$M_{lens} = M_{true} = (1-b)M_X$$
  $b$  : mass bias

Add non-thermal processes  $\longrightarrow \frac{d(P_{th}+P_{n-th})}{dr} = -\frac{GM(r)\rho(r)}{r^2}$ 



$$M_X + \frac{k_B T(r)r}{\mu m_p G} \frac{P_{n-th}}{P_{th}} \frac{\mathrm{d}ln(P_{th}(r))}{\mathrm{d}ln(r)} = -\frac{k_B T(r)r}{\mu m_p G} \left(\frac{\mathrm{d}ln(\rho(r))}{\mathrm{d}ln(r)} + \frac{\mathrm{d}ln(T(r))}{\mathrm{d}ln(r)}\right)$$

# Candidate for (part of) the mass bias

 $\propto -b$ 

Overestimate of the "true mass"



#### STUDY THE DETECTABILITY OF TURBULENCE TROUGH THOSE NEW GENERATION X-RAY SATELLITES

Observable

-2

-1

600-400-200

0

arcmin

(km/s)

2018

al.

2

Roncarelli

2



Where  $\omega(r_{\perp},r_z)$  is the normalised emissivity



How do I retrieve  $k_{inj}$  and  $k_{diss}$  from the observable?

# What is the relation from $P_{3D}(k_{3d})$ to $P_{2D}(k_{\perp})$ ? (1)

# BUT

Due to the sensitivity of the current X-ray telescopes

Observation limited to the central regions

$$P_{2D}(k_{\perp}; x, y) = P_{3D}(k_{\perp}) \int P_{\omega}(k_z; x, y) \, dk_z$$

(Churazov et al. 2012)

(Zhuravleva et al. 2012)

(ZuHone et al. 2016)



Old Cold Front consequent of passing by cluster that agitated the ICM

(ZuHone et al. 2011)

Potential Kelvin Helmholtz instability in the sloshing cold front (Walker et al. 2017) What is the relation from  $P_{3D}(k_{3d})$  to  $P_{2D}(k_{\perp})$ ? (2)

With new X-Ray satellite

Better sensitivity

From 
$$\bar{v}_z(r_\perp) = \int v_{3d}(r_\perp, r_z) \omega(r_\perp, r_z) \mathrm{d}r_z$$

Full 3D Fourier transform

 $\left(\!\left(P_{2D}(\vec{k}_{\perp})\right)_{S}\right) = \int d^{2}\vec{k}_{\perp_{1}} dk_{z_{1}} P_{3D}\left(\sqrt{k_{\perp_{1}}^{2} + k_{z_{1}}^{2}}\right) P_{\omega}(\vec{k}_{\perp} - \vec{k}_{\perp_{1}}, k_{z_{1}})$ 

Emissivity weighted line of sight velocity power spectrum = Centroid velocity power

spectrum

Velocity power spectrum (described by a Kolmogorov spectrum)

Emissivity power spectrum



# Computation of the emissivity power spectrum

Taking 
$$n_e = \left(1 + \left(rac{r}{r_c}
ight)^2
ight)^{-3eta/2}$$

 $r_c$  the cluster core radius eta=2/3

We can define  $\omega(r_{\perp},r_z)=rac{n_e^2(r_{\perp},r_z)}{\int n_e^2(r_{\perp},r_z)~dz}$ 

In the coordinate system  $(r_{\perp}=\sqrt{x^2+y^2}, \theta, z)$ 

We got 
$$\omega(r_{\perp},z) = rac{2}{\pi} rac{(r_c^2 + r_{\perp}^2)^{3/2}}{(r_c^2 + r_{\perp}^2 + z^2)^2}$$

Taking the squared amplitude of its Fourier transform

$$P_{\omega}(k_{\perp},k_z) = |\tilde{\omega}(k_{\perp},k_z)|^2 = \frac{1}{2\pi} \left( \int_{r_c}^{+\infty} (|c^2k_z| + c) \ J_0\left(k_{\perp}(c^2 - r_c^2)^{1/2}\right) e^{-|ck_z|} dc \right)^2$$









 $J_{(}$ 

$$\int_{0}^{\infty} P_{\omega}(k_{\perp},k_{z}) dk_{z} = \int_{0}^{\infty} |\tilde{\omega}(k_{\perp},k_{z})|^{2} dk_{z} = \int_{0}^{\infty} \frac{1}{2\pi} \left( \int_{r_{c}}^{\infty} (|c^{2}k_{z}|+c) J_{0}\left(k_{\perp}(c^{2}-r_{c}^{2})^{1/2}\right) e^{-|ck_{z}|} dc \right)^{2} dk_{z}$$

Pretimica and

#### Variations with the spectral index





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13

#### Variations with the core radius



# What about dissipation scales?



# Unlikely to probe the dissipation scales







### The Structure Function

$$SF(r_{\perp}) = \left\langle \left| \bar{v}_z(\chi_{\perp} + r_{\perp}) - \bar{v}_z(\chi_{\perp}) \right|^2 \right\rangle_{r_{\perp}}$$

Prediction

 $= |\tilde{\bar{v}}_z(k_\perp)|^2 -$ 



#### Reconstruction



Mock observation of the velocity centroid for Astro-H satellite resolution

16

# Conclusion and Prospects

#### We have shown that:

Dinstiguishe between various dissipation scale

Dinstiguishe between various injection scale



For Kolmogorov spectrum

#### BUT

Maybe not the best statistics



Look at the Structure Function

Cluster density is irregular



Add density fluctuation to the model

Take into account the instruments response



Statistical and systematics errors





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Orsay

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