

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{v} \times \mathbf{B}) + \eta \nabla^2 \mathbf{B} \quad \text{Induction Equation}$$
$$\frac{D\rho}{Dt} + \rho \nabla \cdot \mathbf{v} = 0 \quad \text{Mass Conservation}$$
$$\rho \frac{D\mathbf{v}}{Dt} = -\nabla p + \mathbf{j} \times \mathbf{B} + \rho \mathbf{g} + \text{Viscous Terms} \quad \text{Motion}$$
$$\frac{\rho^\gamma}{\gamma-1} \frac{D}{Dt} \left( \frac{p}{\rho^\gamma} \right) = \nabla \cdot \left( \kappa_{||} \nabla T \right) - \rho^2 Q(T) + H(s, t, \mathbf{B}, \rho, T) \quad \text{Energy}$$
$$\rho = \frac{R_p T}{\mu} \quad \text{Gas Law}$$
$$\nabla \cdot \mathbf{B} = 0 \quad \text{Gauss' Law}$$

# Turbulence (with and without $\mathbf{B}$ )

Sébastien GALTIER

IAS - UPS - IUF



# Navier-Stokes equations



*Claude Navier (1785-1836) ; George Stokes (1819-1903).*

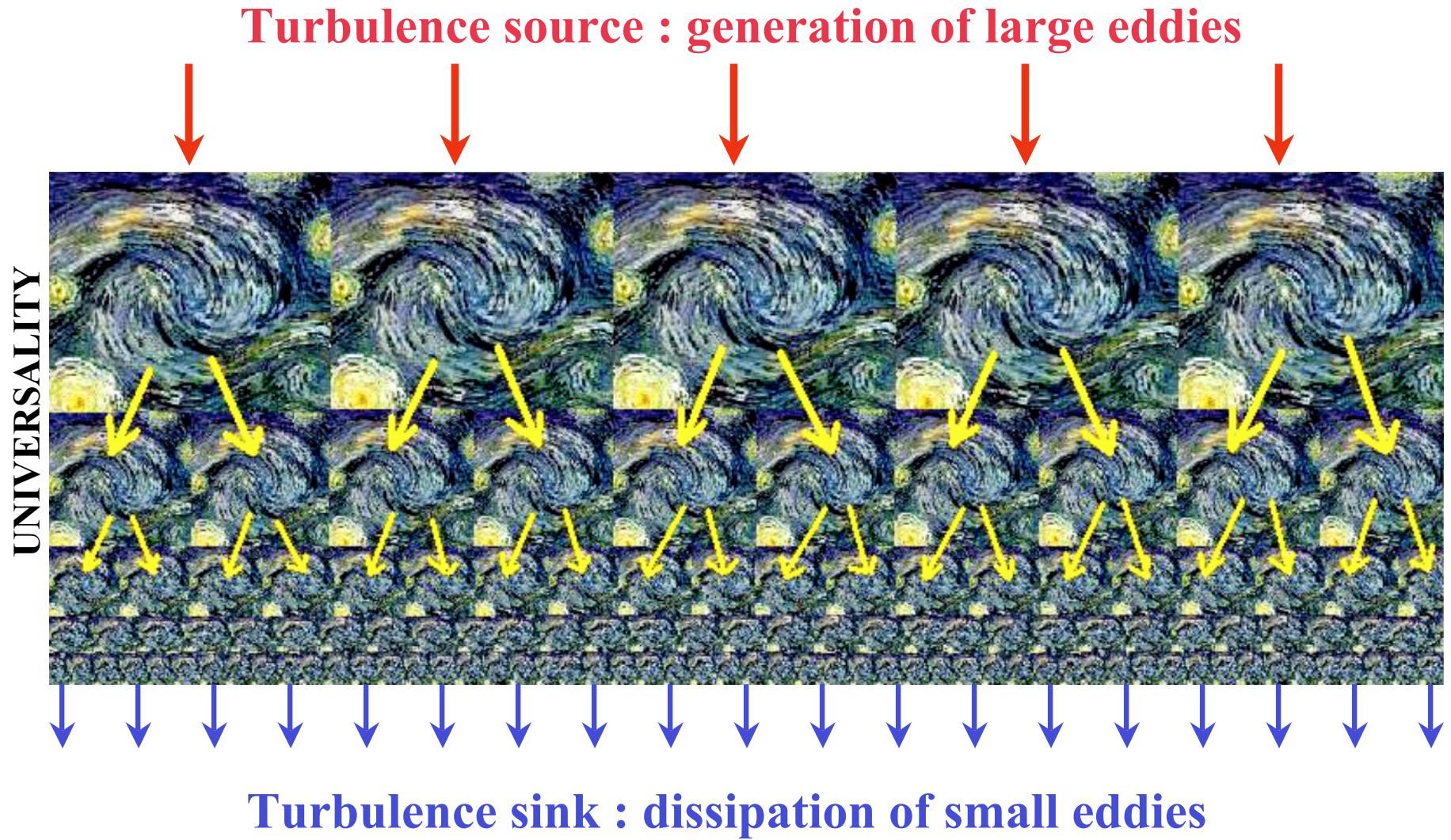
$$\partial \mathbf{v} / \partial t + \mathbf{v} \cdot \nabla \mathbf{v} = -\nabla P + \nu \Delta \mathbf{v} ,$$

$$\nabla \cdot \mathbf{v} = 0 ,$$

Reynolds number:  $R_e = U L / \nu$

$$R_e > 1000 \gg 1$$

# Concept of cascade



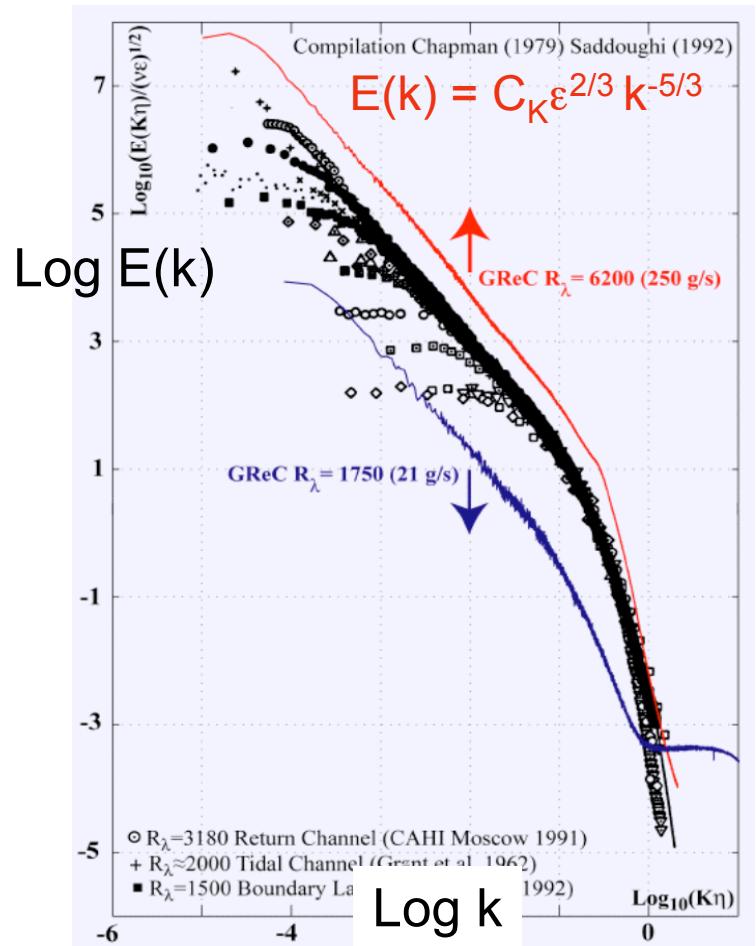
# Homogeneous turbulence spectrum

$$\begin{cases} R_{ij} = \langle u_i(\mathbf{x}) u_j(\mathbf{x}+\mathbf{r}) \rangle = R_{ij}(\mathbf{r}) \\ \phi_{ij}(\mathbf{k}) = (1/2\pi)^3 \iiint_{R^3} R_{ij}(\mathbf{r}) e^{-i\mathbf{k}\cdot\mathbf{r}} d\mathbf{r} \end{cases} \rightarrow E(\mathbf{k}) = (1/2) \phi_{ii}(\mathbf{k})$$



Here

Kolmogorov's spectrum



# Solar wind turbulence

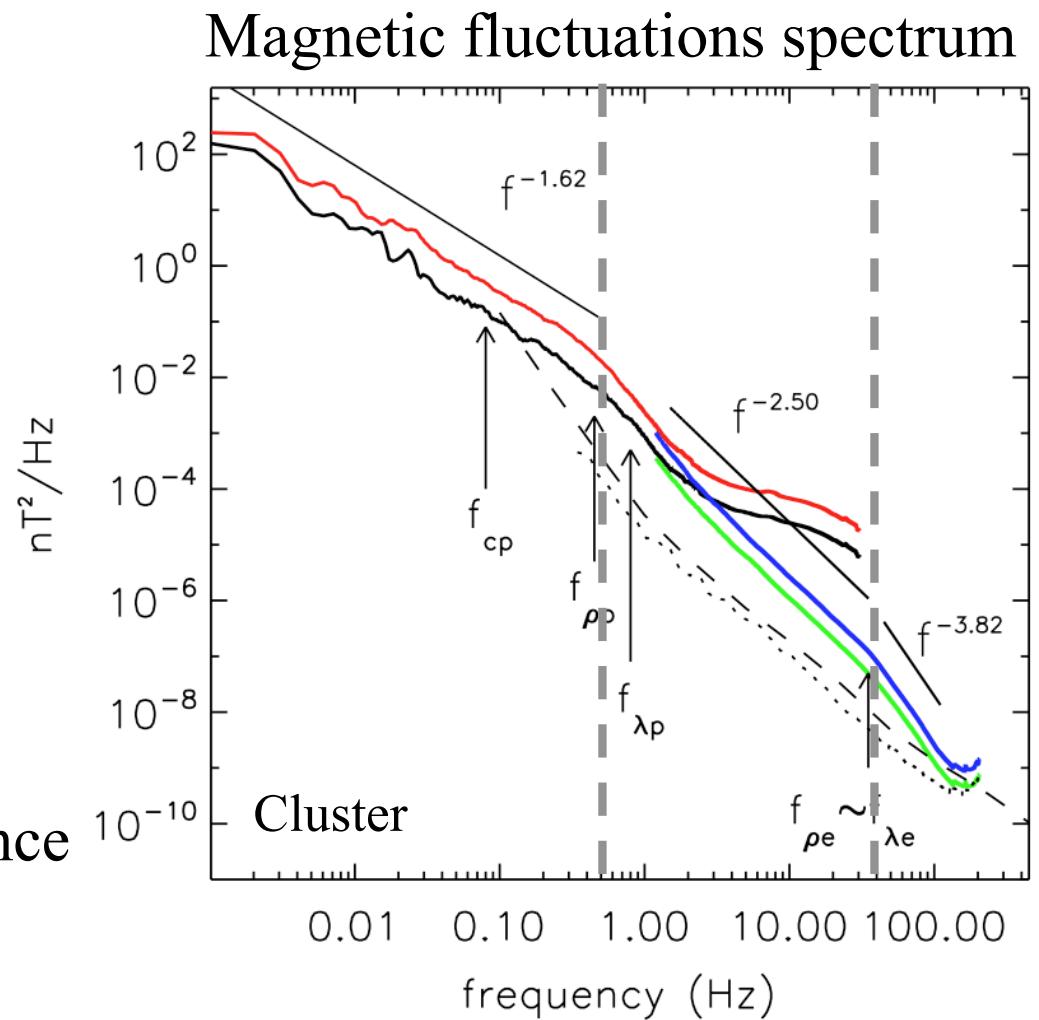
Taylor hypothesis:

$$k \approx 2\pi f / V_{SW}$$

Possible models:

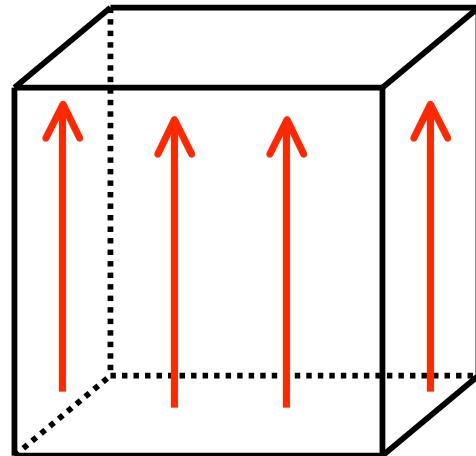
$f < 1/2\text{Hz}$  : **MHD** turbulence

$f > 1/2\text{Hz}$  : **Hall MHD** turbulence



[Sahraoui et al., PRL, 2009]

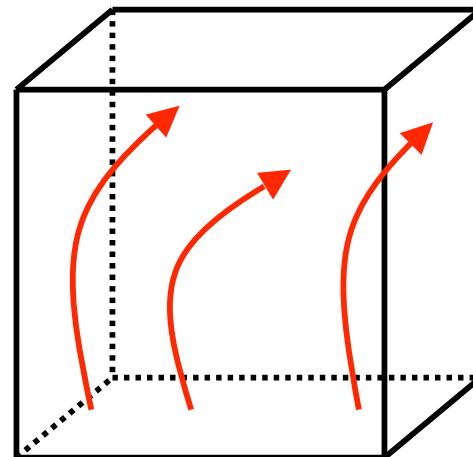
# Regimes of turbulence in MHD



*MHD is non invariant under  
Galilean transformation*

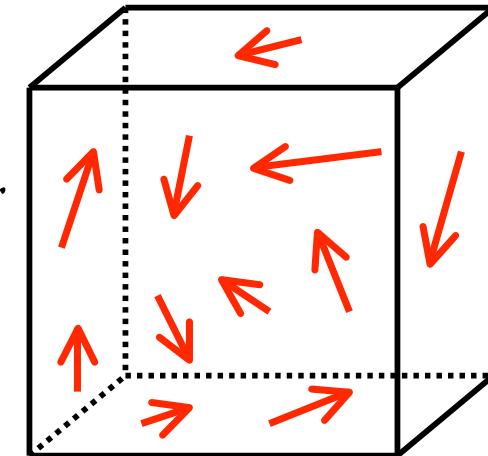
$$\mathbf{B} = \mathbf{B}_0 + \epsilon \mathbf{b}$$

Wave turbulence  
( $\epsilon \ll 1$ )



$$\mathbf{B} = \mathbf{B}_0 + \mathbf{b}$$

Anisotropic strong turbulence



$$\mathbf{B} = \mathbf{b}$$

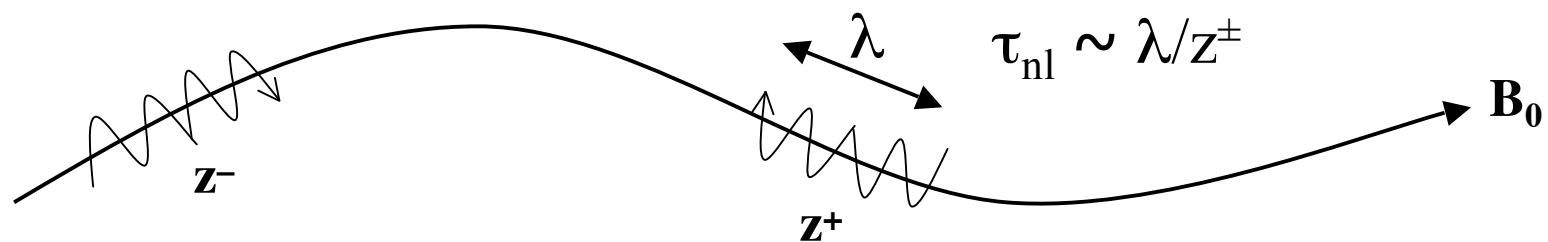
Isotropic strong  
turbulence

# Concept in MHD turbulence

The main difference between neutral fluids and MHD  
is the presence of **Alfvén waves** [Alfvén, Nature, 1942]

Wavepackets interact nonlinearly on a crossing time:  $\tau_A \sim \lambda / B_0$

$$\left\{ \begin{array}{l} \partial_t \mathbf{z}^\pm \mp \mathbf{B}_0 \cdot \nabla \mathbf{z}^\pm = - \mathbf{z}^\mp \cdot \nabla \mathbf{z}^\pm - \nabla P_* \\ \nabla \cdot \mathbf{z}^\pm = 0 \end{array} \right. \quad \mathbf{z}^\pm = \mathbf{u} \pm \mathbf{b} \text{ are the Elsässer fields}$$



Need  $N$  **collisions** between wavepackets :  $\tau_{tr} \sim N\tau_A \sim (\tau_{nl}/\tau_A)^2 \tau_A$

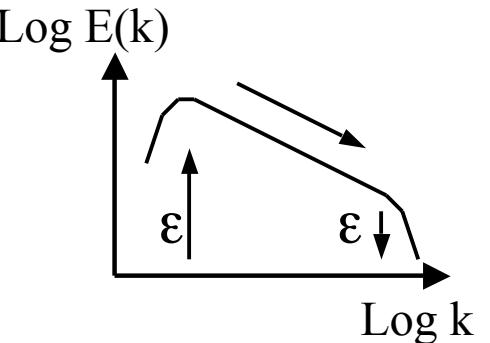
[Iroshnikov, Sov. Astron., 1964; Kraichnan, Phys. Fluids, 1965]

# Theories for isotropic MHD

[Iroshnikov, Sov. Astron., 1964; Kraichnan, Phys. Fluids, 1965]

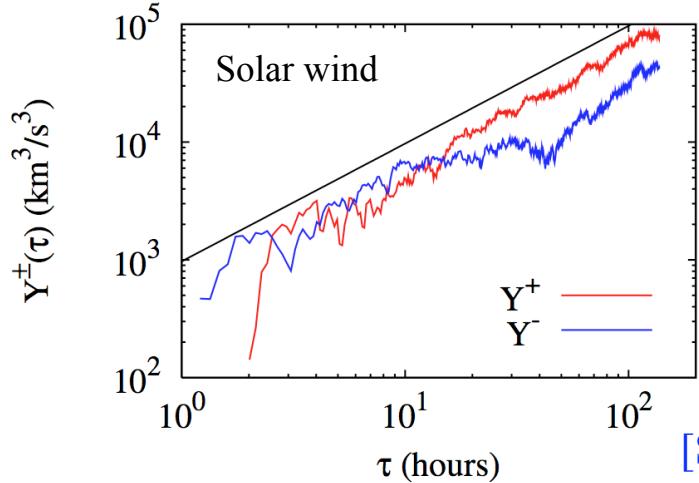
$$\varepsilon \sim z_\lambda^4 / (\lambda B_0) \sim (E_k k)^2 k / B_0 \sim E_k^2 k^3 / B_0$$

$$E(k) \sim (\varepsilon B_0)^{1/2} k^{-3/2}$$



- 1) Problem: **isotropy** is assumed !!
- 2) Very **close** to the Kolmogorov prediction ( $k^{-5/3}$ )
- 3) Heuristic result **not compatible** with the 4/3's exact law

Homogeneous isotropic medium, at  $\text{Re}, \text{Rm} \rightarrow +\infty$



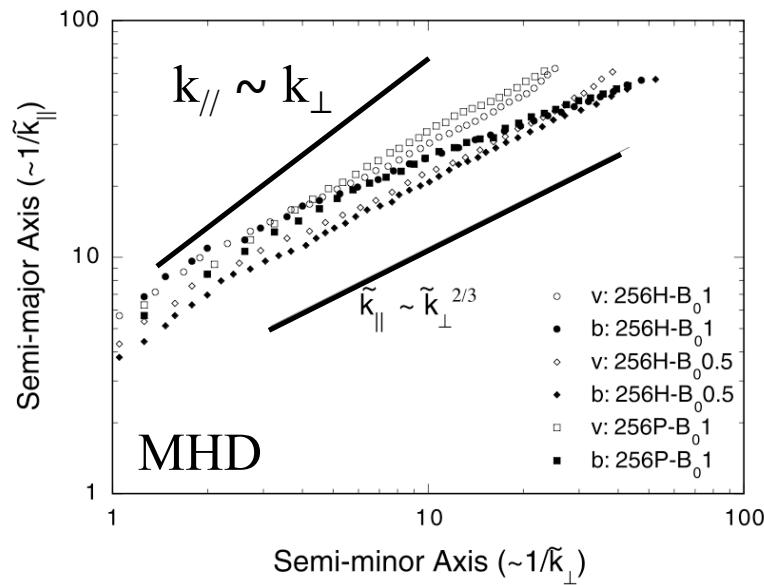
[Politano & Pouquet, PRE, 1998]

$$-\frac{4}{3}\varepsilon^\pm r = \underbrace{\langle \delta z_L^\mp (\delta z^\pm)^2 \rangle}_{Y^\pm} \Rightarrow E(k) \sim k^{-5/3}$$

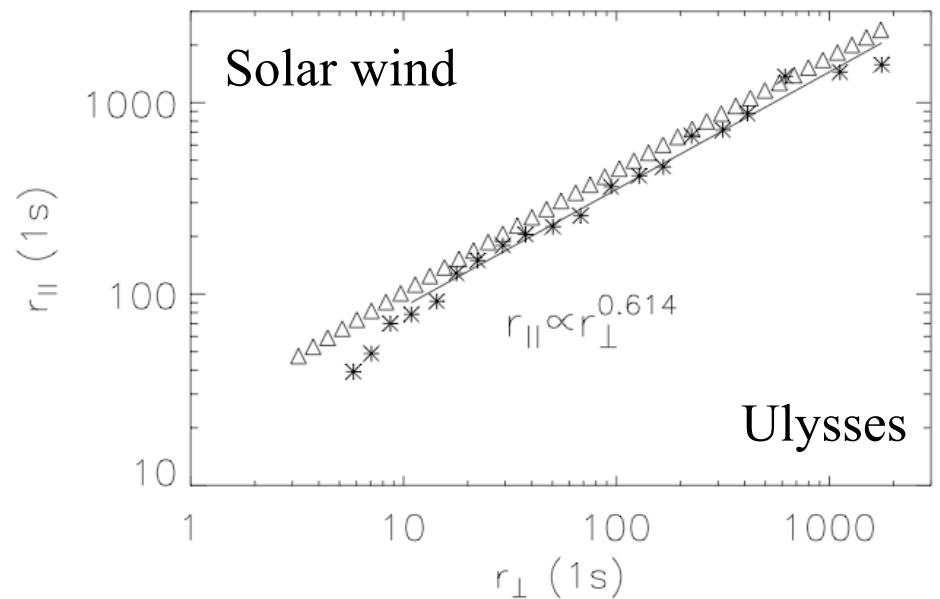
[Sorriso-Valvo et al., PRL, 2007]

# Theories for anisotropic MHD

- Critical balance :  $\tau_{nl} \sim \tau_A$  at **all** scales [Goldreich & Sridhar, ApJ, 1995]
  - phenomenology
  - $E(k_\perp) \sim k_\perp^{-5/3}$  and  $E(k_\parallel) \sim k_\parallel^{-2}$
  - increase of anisotropy at small scales :  $k_\parallel \sim k_\perp^{2/3}$



[Cho & Vishniac, ApJ, 2000]

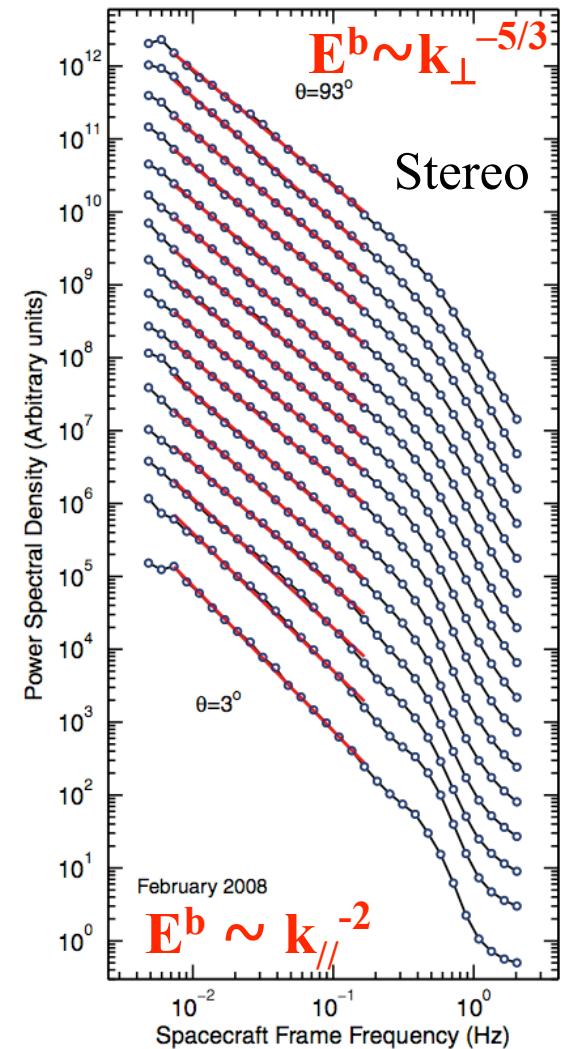
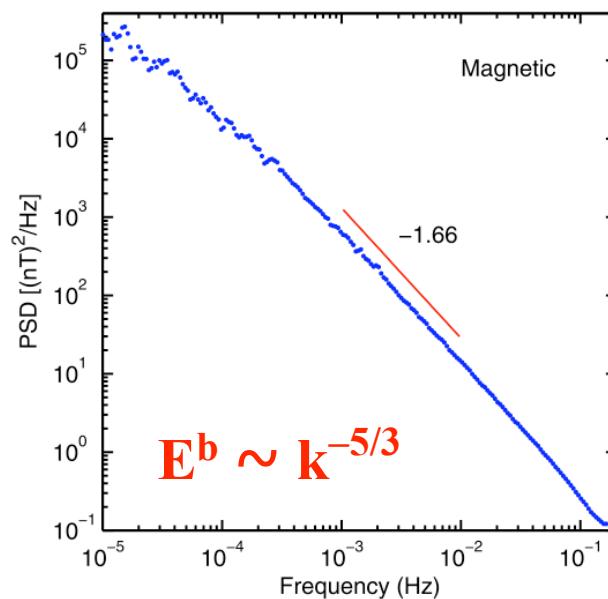
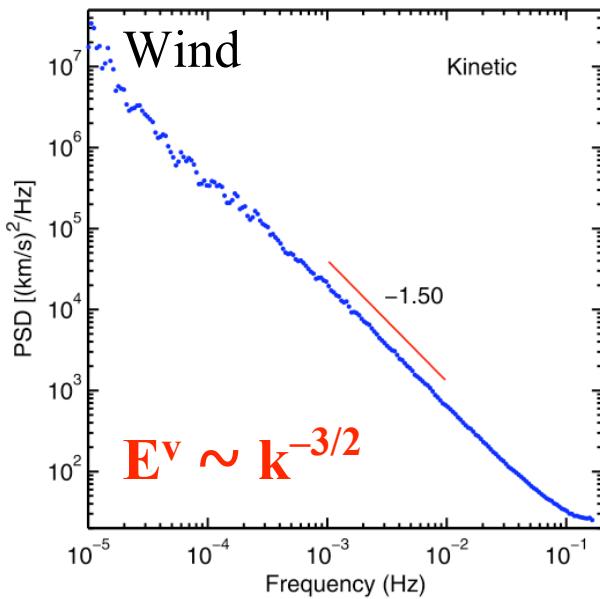


[Luo & Wu, ApJ, 2010]

# Solar wind energy spectra

Taylor hypothesis:  $k \approx 2\pi f / V_{SW}$

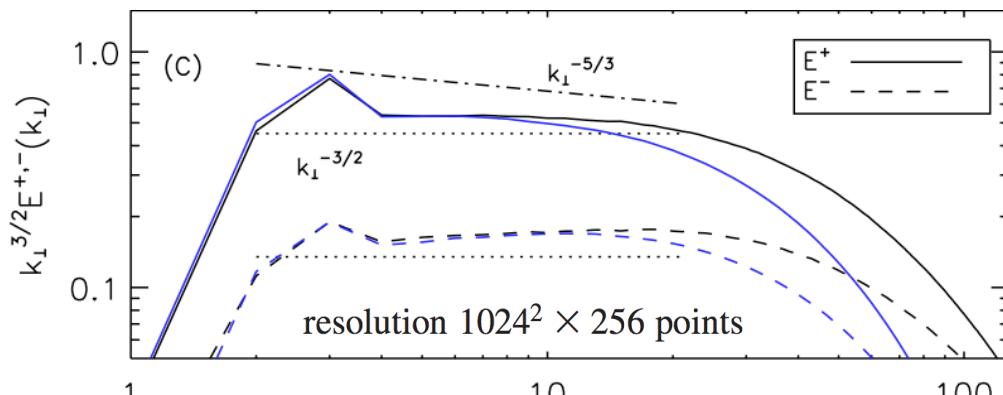
No unique solution !



[Podesta, J Geophys. Res., 2006; Podesta, ApJ, 2009]

# Theories for anisotropic MHD

- Dynamic alignment :  $\mathbf{v} \cdot \mathbf{b} = \pm v b \cos \theta$  [Boldyrev, PRL, 2006]
  - phenomenology
  - $\theta \sim \lambda_{\perp}^{1/4}$  and  $E(k_{\perp}) \sim k_{\perp}^{-3/2}$
  - anisotropic relation :  $k_{\parallel} \sim k_x^{1/2}$



[Perez & Boldyrev, PRL, 2009]

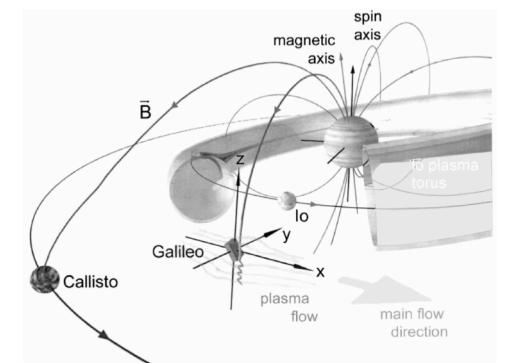
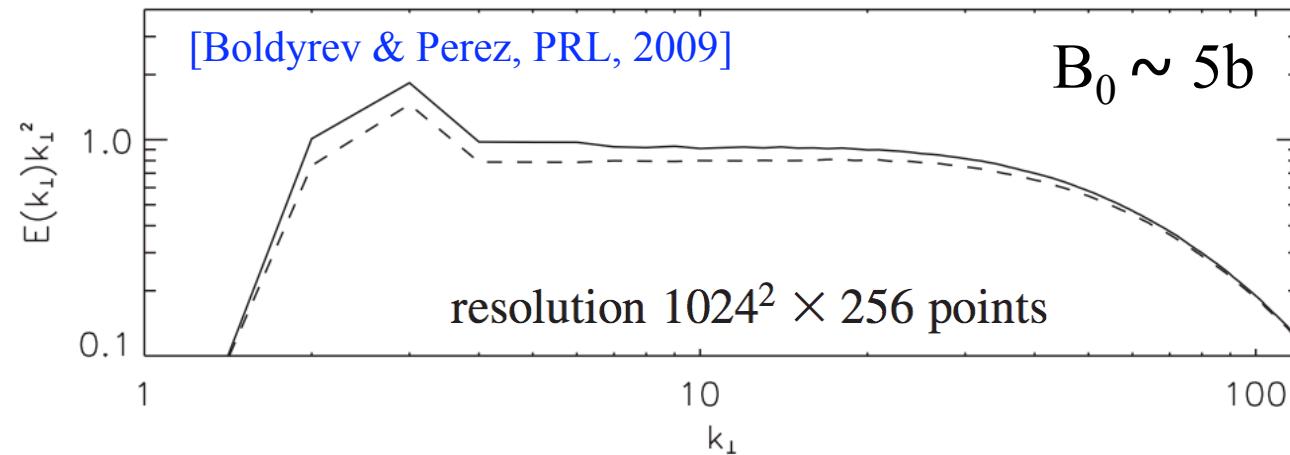
But DNS reveals that :

$$E(k_{\perp}, k_{\parallel}=0) \sim k_{\perp}^{-3/2}$$

[Bigot et al., PRE, 2008]

# Theories for anisotropic MHD

- Alfvén wave turbulence ( $B_0 \rightarrow +\infty$ ;  $\tau_A \ll \tau_{nl}$ ) [SG et al., J. Plasma Phys., 2000]
  - **exact** asymptotic closure
  - $E(k_\perp) \sim k_\perp^{-2}$  : exact solution!
  - **observed** in Jupiter's magnetosphere  
[Saur et al., A&A, 2002]

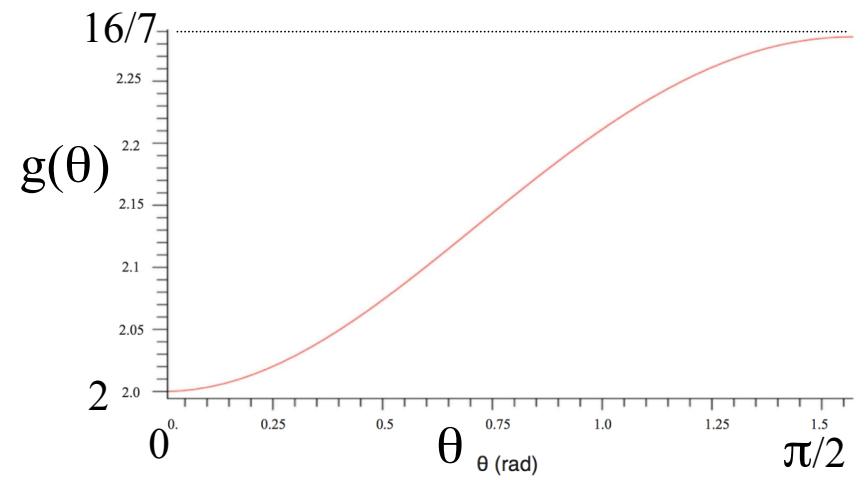
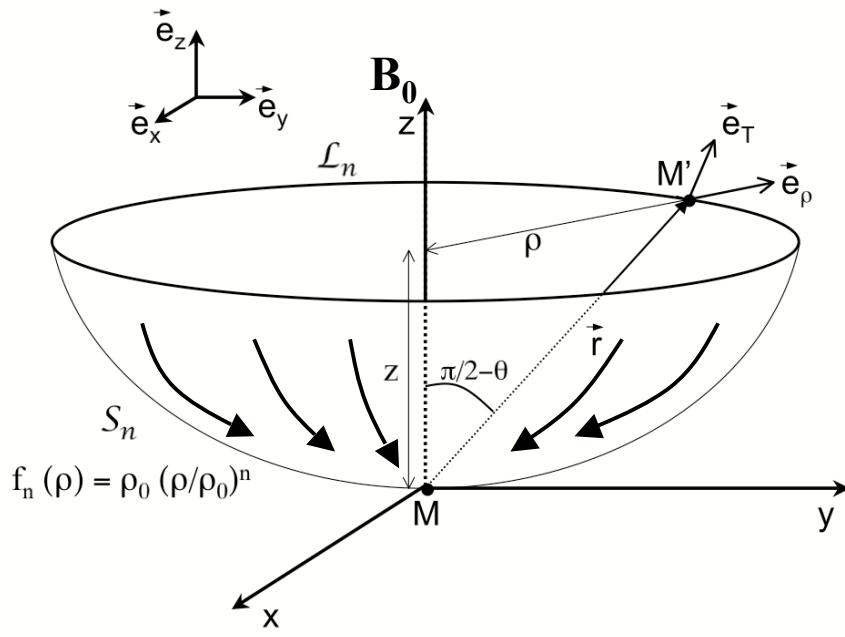


# Theories for anisotropic MHD

Critical balance:  $k_{\parallel} \sim k_{\perp}^{2/3}$   $\Leftrightarrow$  space correlation is **foliated**  
[SG, 2010]

$$-g(\theta)\varepsilon^{\pm}r \mathbf{e}_T = \mathbf{F}_T^{\pm}(\mathbf{r})$$

$$\mathbf{F}^{\pm}(\mathbf{r}) = \langle \delta \mathbf{z}^{\mp} (\delta \mathbf{z}^{\pm})^2 \rangle$$



# Beyond MHD : Hall MHD

- **Exact** law for (isotropic) Hall MHD :

[SG, PRE, 2008]

$$E \sim k^{-5/3}$$

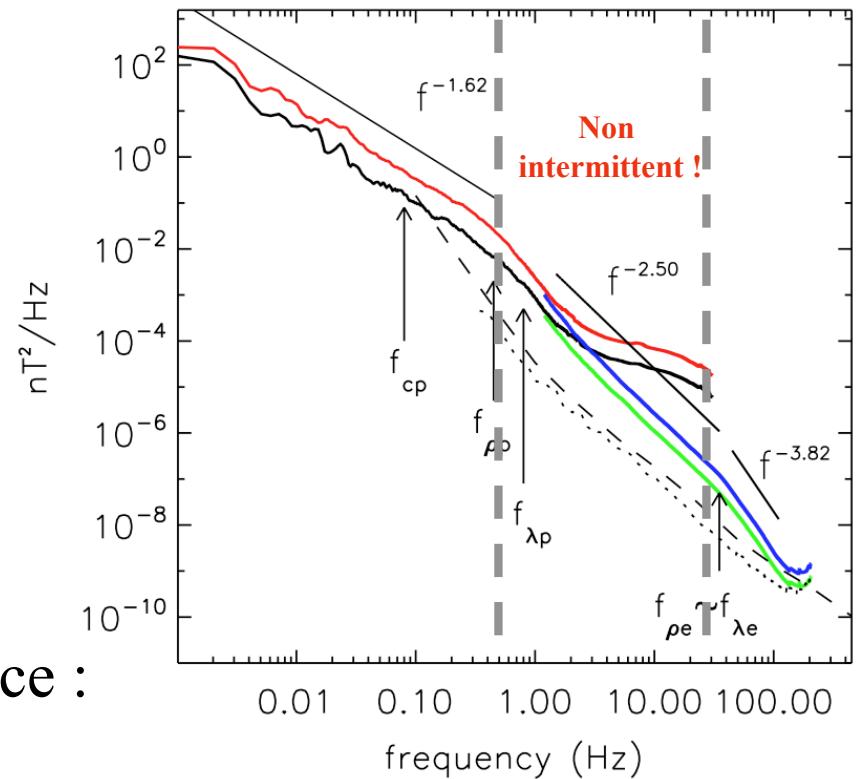
$$-\frac{4}{3}\boldsymbol{\varepsilon}^T \boldsymbol{r} = \underbrace{B_{\parallel ii}^{vvv} + B_{\parallel ii}^{bbb} - 2B_{\parallel ii}^{bvb}}_{+ 4d_I \langle [(\mathbf{J} \times \mathbf{b}) \times \mathbf{b}']_{\parallel} \rangle} E^b \sim k^{-7/3}$$

- **Whistler** (anisotropic) wave turbulence :

$$\rightarrow E^b(k_{\perp}) \sim k_{\perp}^{-2.5} : \text{exact solution}$$

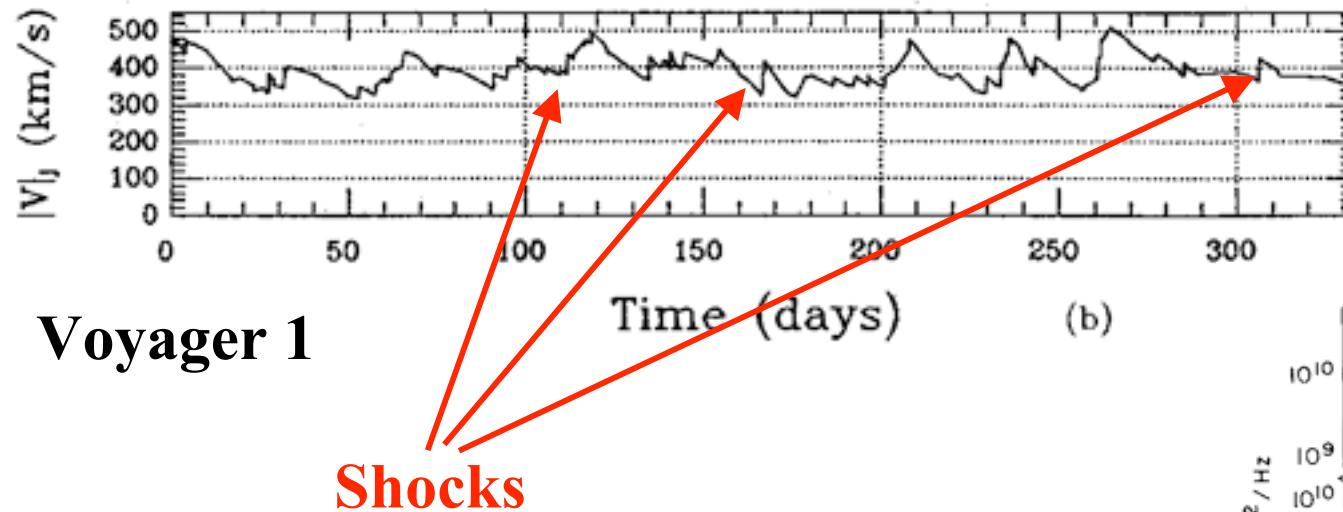
[SG, J. Plasma Phys., 2006]

[Kiyani et al., PRL, 2009]



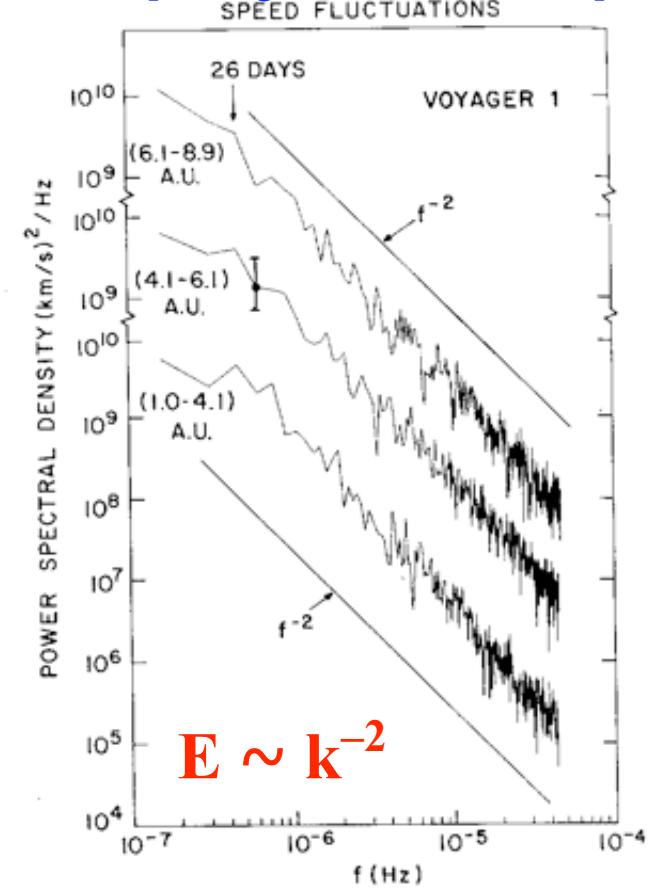
# Compressible turbulence

[Roberts et al., JGR, 1987]



Solar wind  
(>1AU)

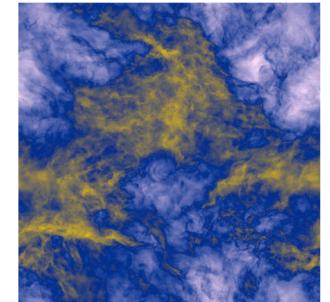
[Burlaga et al., JGR, 1987]



There is **no exact** law for  
3D compressible turbulence !



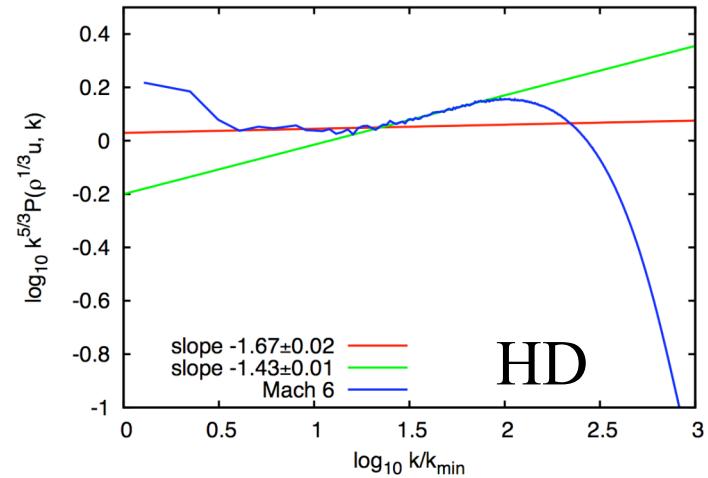
# Compressible turbulence



- 3D M-HD simulations (adiabatic case):

[Kritsuk et al., ApJ, 2007; ASP, 2009]

$$v = \rho^{1/3} u \quad \rightarrow \quad E^v(k) \sim k^{-5/3}$$



**2048<sup>3</sup>**

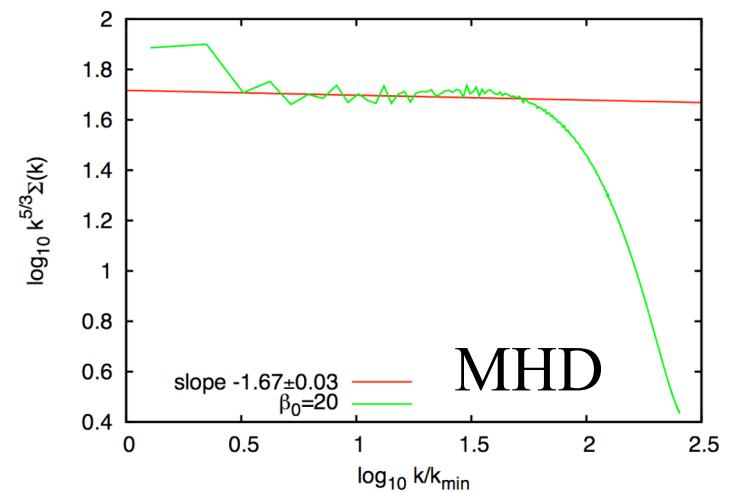
- First analysis (HD + adiabatic):

$$-\mathcal{E} = \nabla \cdot \mathbf{F}(\rho, \mathbf{u}) + \mathbf{S}(\nabla \cdot \mathbf{u}, \rho, \mathbf{u})$$

[SG, 2010]

**F:** energy flux vector

**S:** compressible source



# **Conclusion**

Thank you !