Baryon Acoustic Oscillations

BAO

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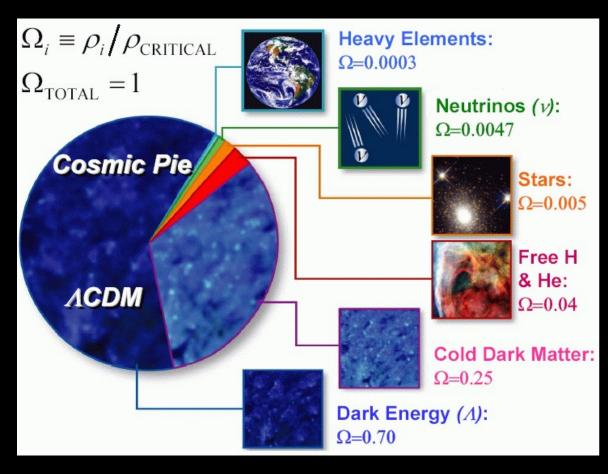
Outline

- Dark energy context
- The BAO probe
- How standard is the BAO ruler?

Material borrowed mainly from these sources

- http://cdm.berkeley.edu/doku.php?id=baopages
- http://cmb.as.arizona.edu/~eisenstein/acousticpeak/
- http://mwhite.berkeley.edu/BAO

The Big Problems: Dark Energy and Dark Matter



The confirmation of Dark Energy points to major holes in our understanding of fundamental physics

95% of the Universe is in forms unknown to us

1998 Science breakthrough of the year



Probing Dark Energy

• Dark energy is probed through how it influences the expansion rate of the universe H(z) and the rate growth of structure g(z)

$$H^{2}(z) = H^{2}_{0} \left[\Omega_{M} (1+z)^{3} + \Omega_{R} (1+z)^{4} + \Omega_{K} (1+z)^{2} + \Omega_{DE} (1+z)^{3(1+w)} \right]$$
matter radiation curvature dark energy

g(z) in general a complicated function of cosmological parameters

Probing Dark Energy

Best observational probes (DETF)

- Weak lensing (geometrical & growth)
- Baryon acoustic oscillations (geometrical)
- Supernovae (geometrical)
- Clusters of galaxies (growth & geometrical)

Probing Dark Energy

• Geometric test: integrals over H(z):

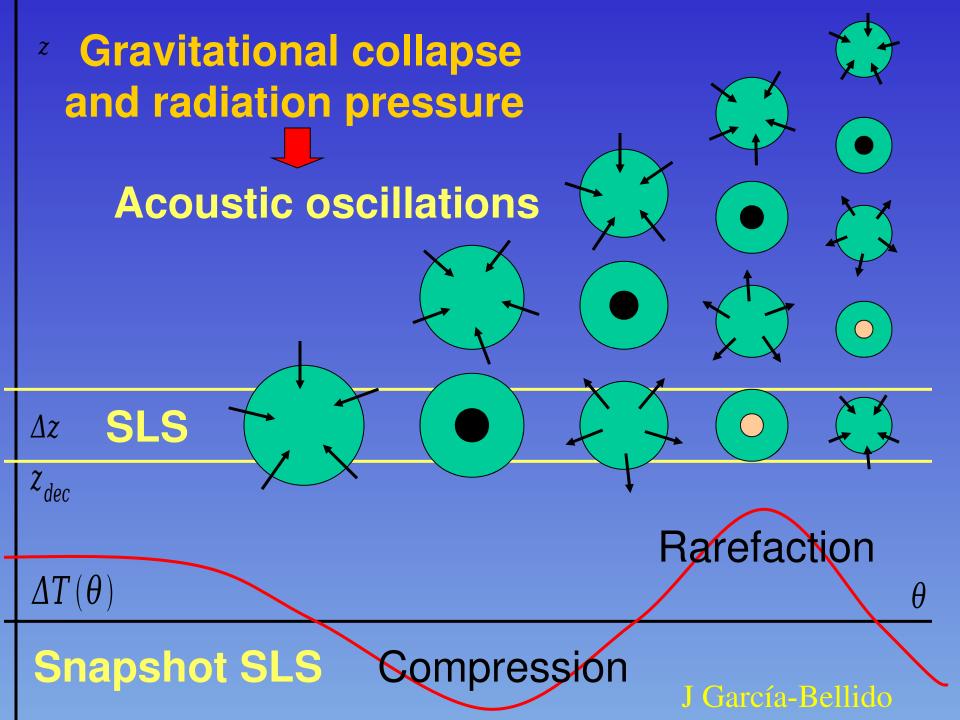
Comoving distance $r(z) = F[\int dz/H(z)]$ Standard Candles Supernovae $D_L(z) = (1+z) r(z)$ Standard Rulers Baryon Oscillations $D_A(z) = (1+z)^{-1} r(z)$ Standard Population Clusters $dV/dzd\Omega = r^2(z)/H(z)$

Growth of Structure test: g(z)

Clusters, Weak lensing, clustering

- DE equation of state: $w = P/\rho$
- DE parameterization: $w(z) = w_0 + w_a z/(1+z)$

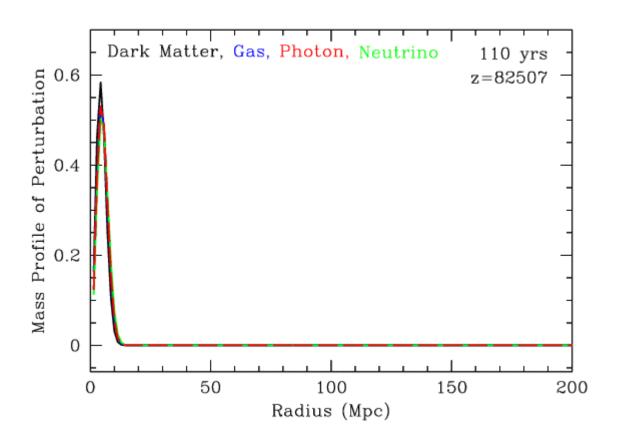
- The early universe composed of photons, baryons and dark matter
- Photons and baryons tightly coupled
- The early universe very homogeneous except tiny fluctuations
- Evolution:
 - as it expands it becomes cooler and less dense
 - fluctuations grow due to gravity
- Acoustic waves are generated as the photon-baryon fluid is attracted and falls onto the overdensities: compressions and rarefactions



- These acoustic waves propagate until the universe becomes cool enough for the electrons and protons to recombine and then the baryons and photons decouple
- The time when the baryons are "released" from the drag of the photons is known as the drag epoch, z_d
- From then on photons expand freely while the acoustic waves "freeze in" the baryons in a scale given by the size of the horizon at the drag epoch
- Progressively, baryons fall into dark matter potential wells but also dark matter is attracted to baryon overdendities

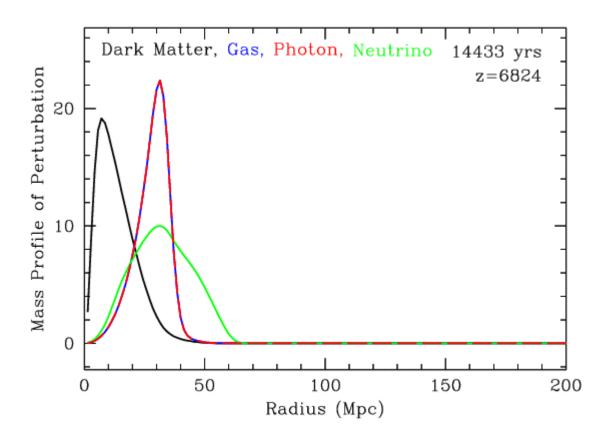
Description of one perturbation (Eisenstein)

- Four species: Dark matter, Baryons, Photons & Neutrinos
- Initial perturbations adiabatic: all species perturbed approximately same fractional amount

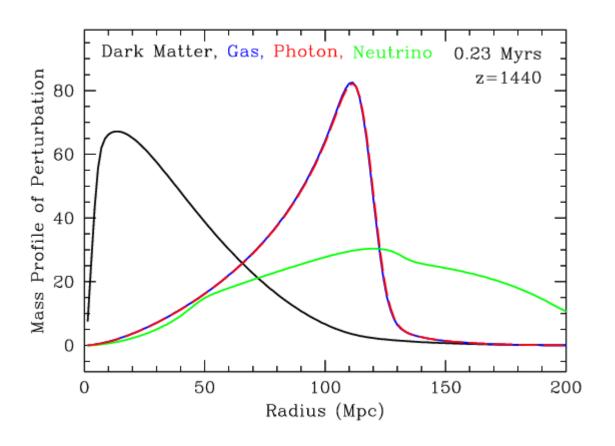


- Neutrinos do not interact and move too fast to be stopped by gravity, so they stream away
- Dark matter responds to gravity and falls onto the perturbation overdensity
- Perturbation dominated by photons and baryons as they are coupled. Perturbation is overdensity and overpressure.

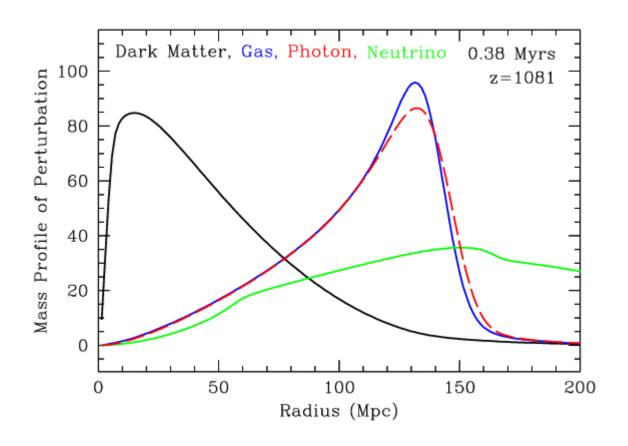
 Overpressure tries to equalize with surrounding resulting in an expanding sound wave moving at the speed of sound which is approximately 2/3 the speed of light
- The perturbation in photons & baryons is carried outward



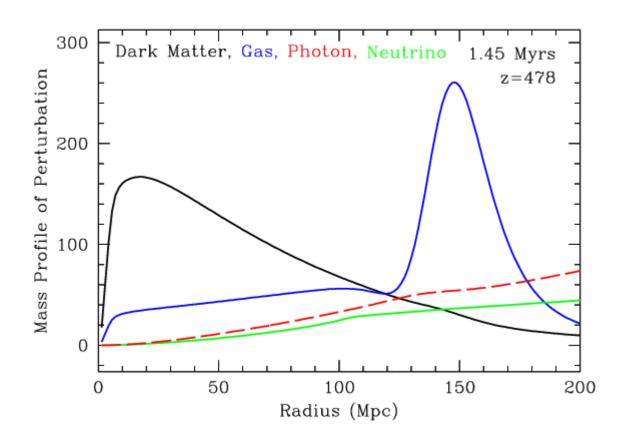
- The photons & baryons continue to expand
- Neutrinos spread out
- Dark matter continues to fall into perturbations, which grows



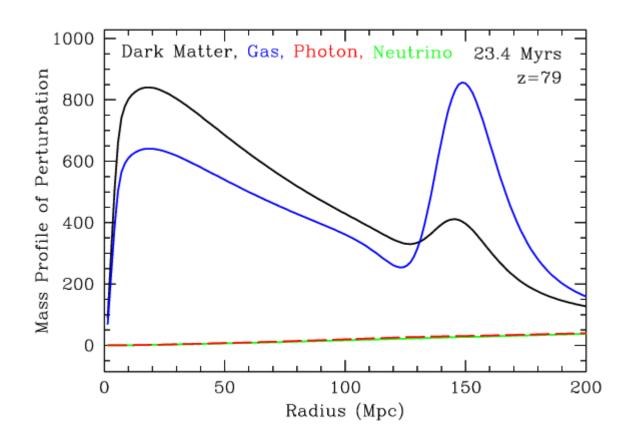
- As the expanding universe cools down, it reaches a point when the electrons and protons begin to combine
- Photons do not scatter as efficiently and start to decouple
- The sound speed drops and the pressure wave slows down



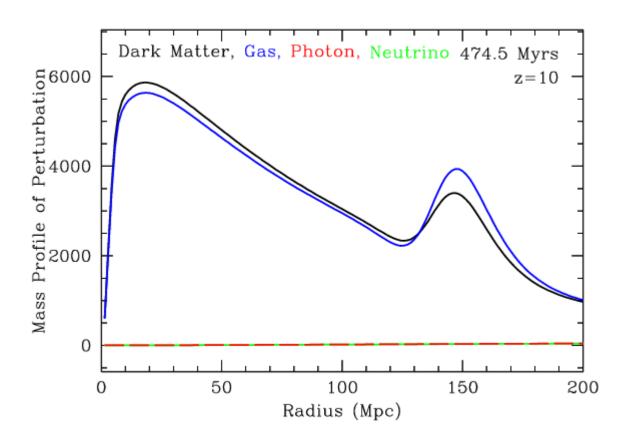
- The process continues until the photons completely decouple and then its perturbation smoothes out
- The sound speed of the baryon perturbation drops so much that the pressure wave stalls

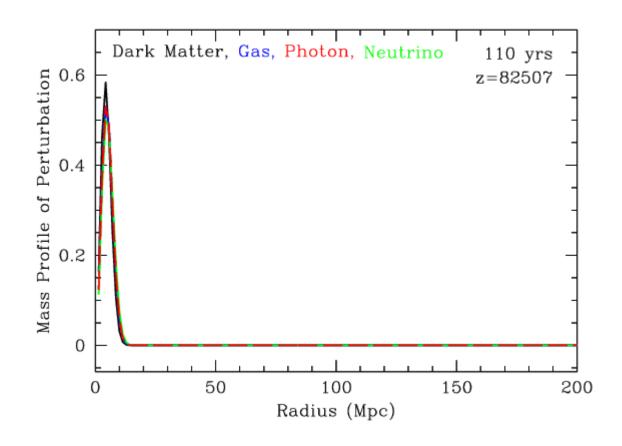


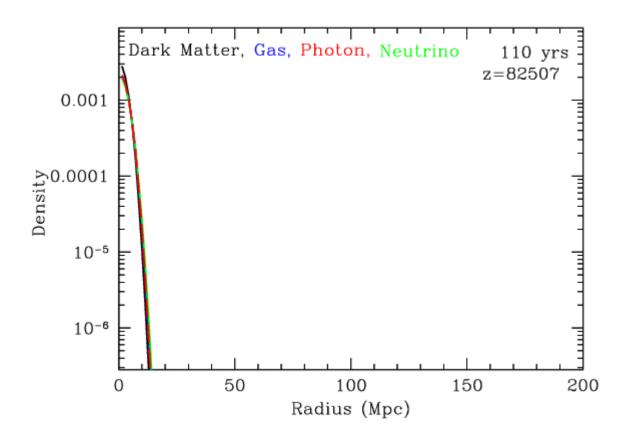
- We are left with the original dark matter perturbation surrounded by a baryon perturbation in a shell
- The two components attract each other and the perturbation start to mix



- Eventually the dark matter and baryon perturbations come together
- The acoustic peak perturbation is lower as the dark matter dominated in mass the baryons







• As galaxies form in matter (dark matter and baryons) overdensities, most of galaxies are at the original perturbation position, but there is a 1% enhancement of galaxies at the scale of the acoustic scale than can be seen in the galaxy correlation function

Photons and baryons decouple at the drag epoch

$$z_{d} = 1291 \frac{(\Omega_{0} h^{2})^{0.251}}{1 + 0.659(\Omega_{0} h^{2})^{0.828}} \left[1 + b_{1}(\Omega_{b} h^{2})^{b_{2}}\right],$$

$$b_{1} = 0.313(\Omega_{0} h^{2})^{-0.419} \left[1 + 0.607(\Omega_{0} h^{2})^{0.674}\right],$$

$$b_{2} = 0.238(\Omega_{0} h^{2})^{0.223},$$
(4)

• Photon-baryon fluid acoustic waves propagate at the sound speed: $c_s = 1/[3(1+R)]^{1/2}$

$$R \equiv 3\rho_b/4\rho_{\gamma} = 31.5\Omega_b h^2 \Theta_{2.7}^{-4} (z/10^3)^{-1} .$$
 (5)

• Sound horizon at the drag epoch is the comoving distance a wave can travel prior to z_d

$$s = \int_{0}^{t(z_d)} c_s(1+z)dt$$

$$= \frac{2}{3k_{eq}} \sqrt{\frac{6}{R_{eq}}} \ln \frac{\sqrt{1+R_d} + \sqrt{R_d+R_{eq}}}{1+\sqrt{R_{eq}}}, \qquad (6)$$

• The transition from a radiation to matter dominated universe happens at

$$z_{\rm eq} = 2.50 \times 10^4 \Omega_0 \, h^2 \Theta_{2.7}^{-4} \,,$$
 (2)

 At small scales, Silk damping: photons and baryons diffuse during decoupling

$$k_{\text{Silk}} = 1.6(\Omega_b h^2)^{0.52} (\Omega_0 h^2)^{0.73} [1 + (10.4\Omega_0 h^2)^{-0.95}]$$

$$\text{Mpc}^{-1}, \quad (7)$$

Sound horizon depends on

$$s = \int_0^{t_{\text{rec}}} c_s (1+z)dt = \int_{z_{\text{rec}}}^{\infty} \frac{c_s dz}{H(z)}$$

- epoch of recombination
- expansion of the universe
- baryon-to-photon ratio

• Sound horizon well determined by the CMB measurements of the acoustic peaks

	Quantity	Eq.	5-year WMAP
CMB CMB Matter Matter	$egin{array}{c} z_* \ r_s(z_*) \ z_d \ r_s(z_d) \end{array}$	(66) (6) (3) (6)	1090.51 ± 0.95 $146.8 \pm 1.8 \; \mathrm{Mpc}$ 1020.5 ± 1.6 $153.3 \pm 2.0 \; \mathrm{Mpc}$

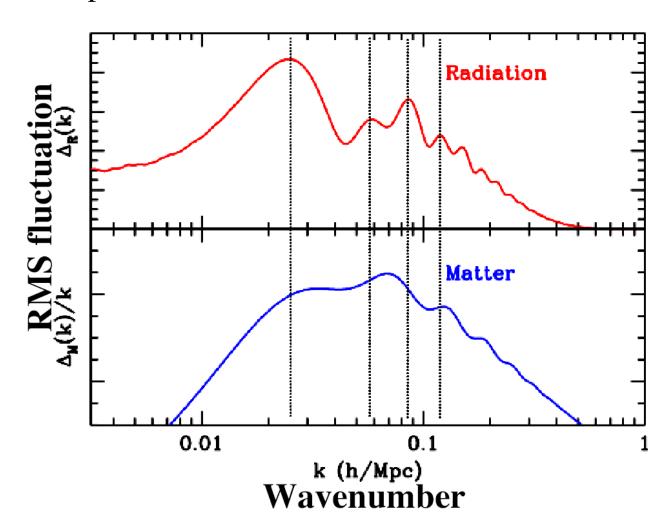
Power spectrum

$$T(k) = \frac{\Omega_b}{\Omega_0} T_b(k) + \frac{\Omega_c}{\Omega_0} T_c(k) . \tag{8}$$

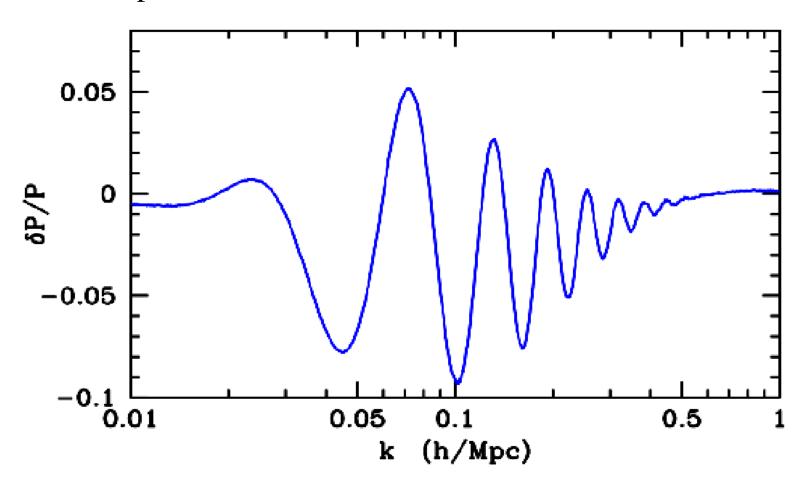
• The baryon part can be approximated as

$$T_b \to \alpha_b \frac{\sin(ks)}{ks} \mathcal{D}(k)$$
 (13)

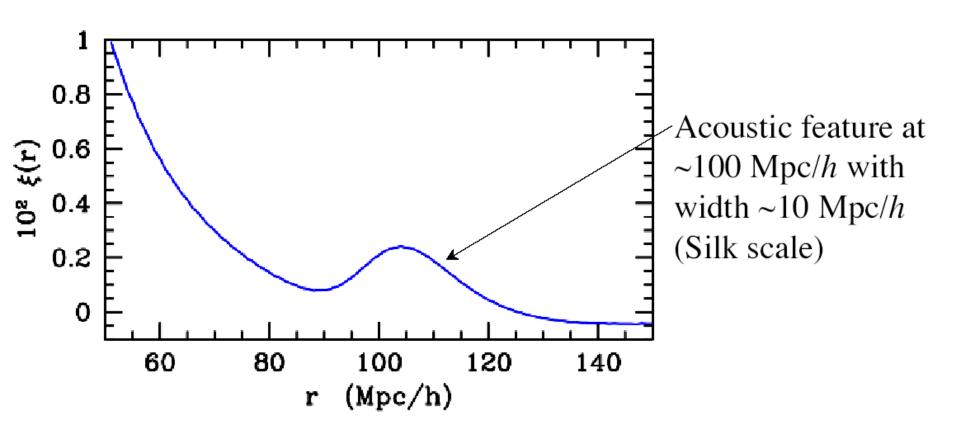
Power spectrum



Power spectrum



Correlation function



- The baryon acoustic oscillations provide a characteristic scale that is "frozen" in the galaxy distribution providing a standard ruler that can be measured as a function of redshift in either the galaxy correlation function or the galaxy power spectrum
- The BAO determination of the universe geometry is quite robust against systematics

The BAO standard ruler

• The BAO standard ruler provides a measurement of the angular diameter distance as a function of redshift

$$\Delta\theta = \ell/(d_A\,(1+z))$$

$$\Delta\theta : \text{angle subtended}$$

$$\ell : \text{intrinsic length}$$

$$d_A(z) = r(z)/(1+z)$$

$$d_A : \text{angular distance}$$

$$r(z) = \int_0^z \frac{dz'}{H(z')}$$

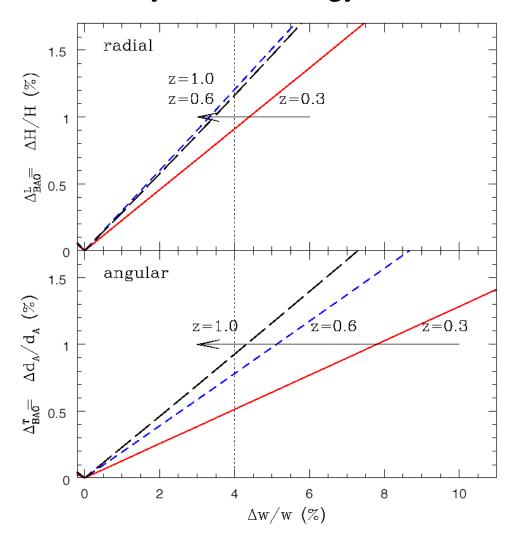
$$r(z) : \text{co-moving distance}$$

• The standard ruler can also be used in the line of sight and measure H(z) directly

$$\ell = \Delta z / H(z)$$
 Δz : redshift subtended

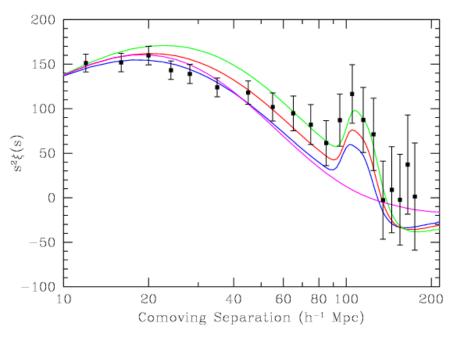
The BAO standard ruler

The BAO sensitivity to dark energy

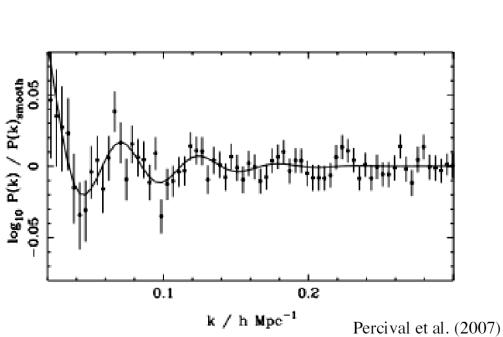


BAO measurements

The BAO scale has been measured in the SDSS and 2dFGRS



Eisenstein et al 2005



Is it a standard ruler?

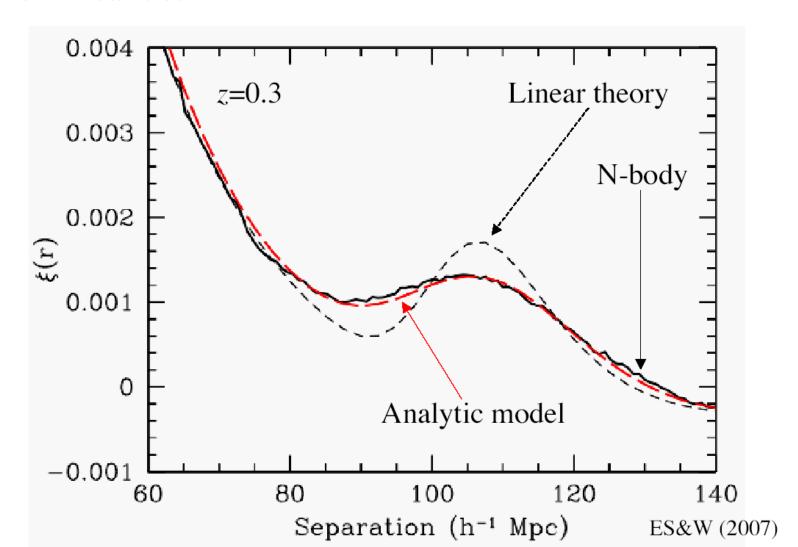
- So far, we have dealt with linear theory in real space
- However, we measure the non-linear power spectrum of galaxy tracers in redshift space
- Simulations help us to understand & treat these effects

Non-linearities

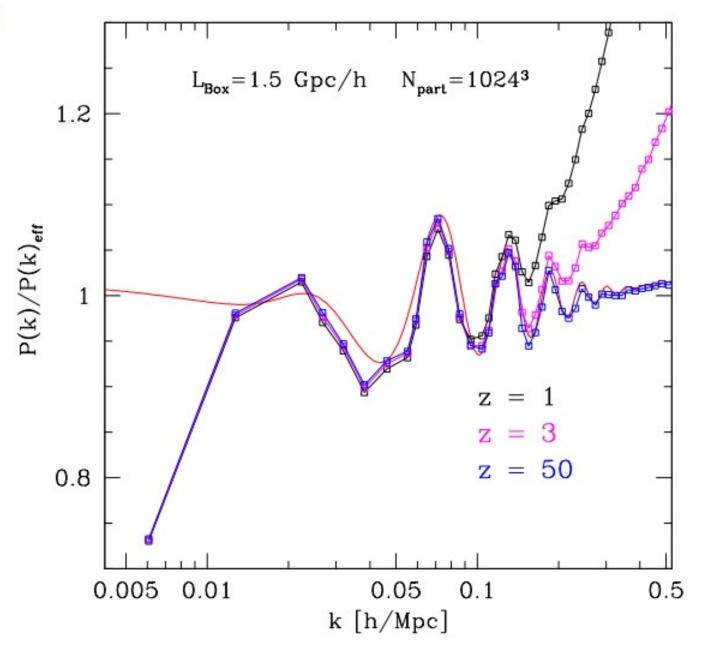
- As structure forms, neighbouring galaxies smear the baryon shell. This produces a broadening of the peak and non-linear power on small scales
- According to simulations it can be modelled

$$P_{nl} = P_{l} \exp(-k^{2} \Sigma_{nl})$$

Non-linearities



Baryon Acoustic Oscillations



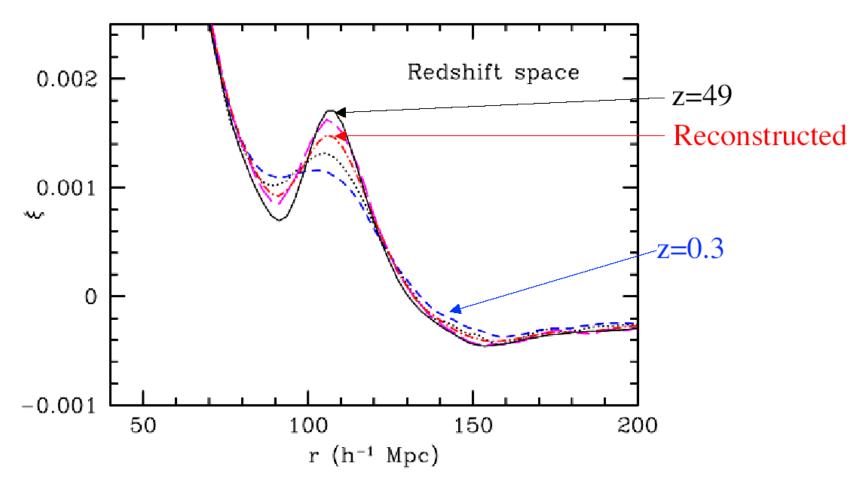
Non-linearities

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• Can it be corrected? Reconstruction (Eisenstein et al 07)

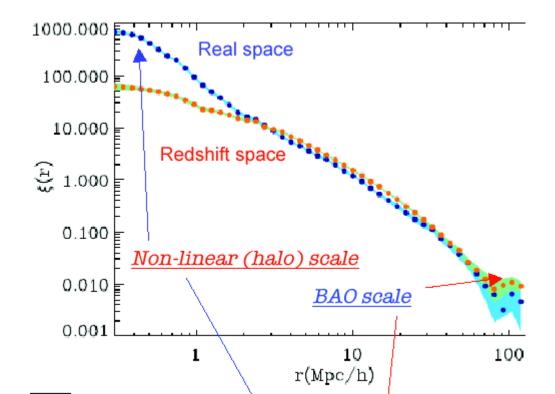
Non-linearities

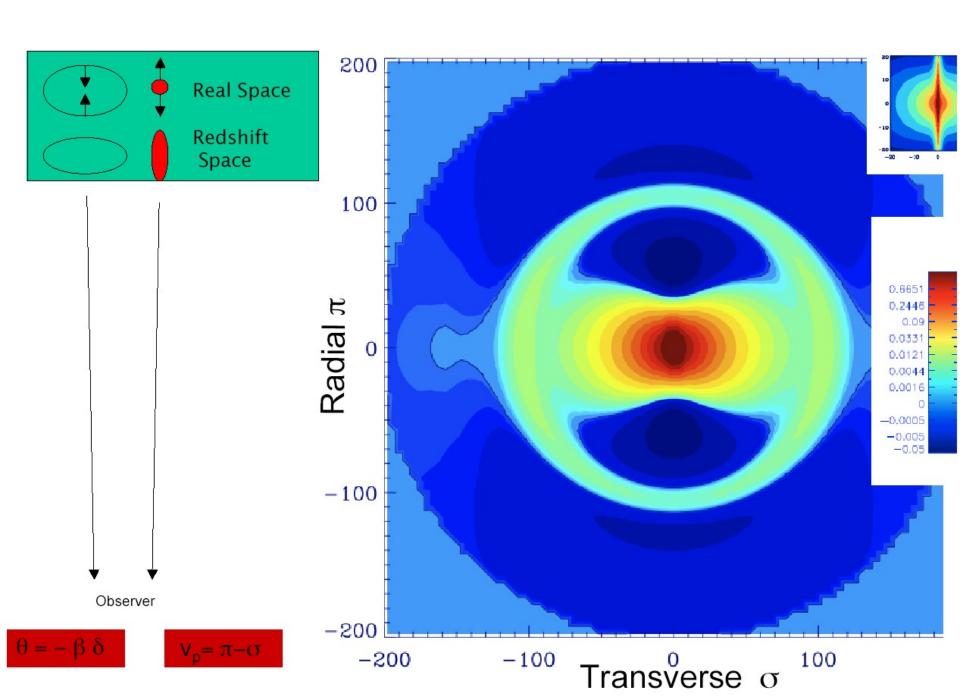


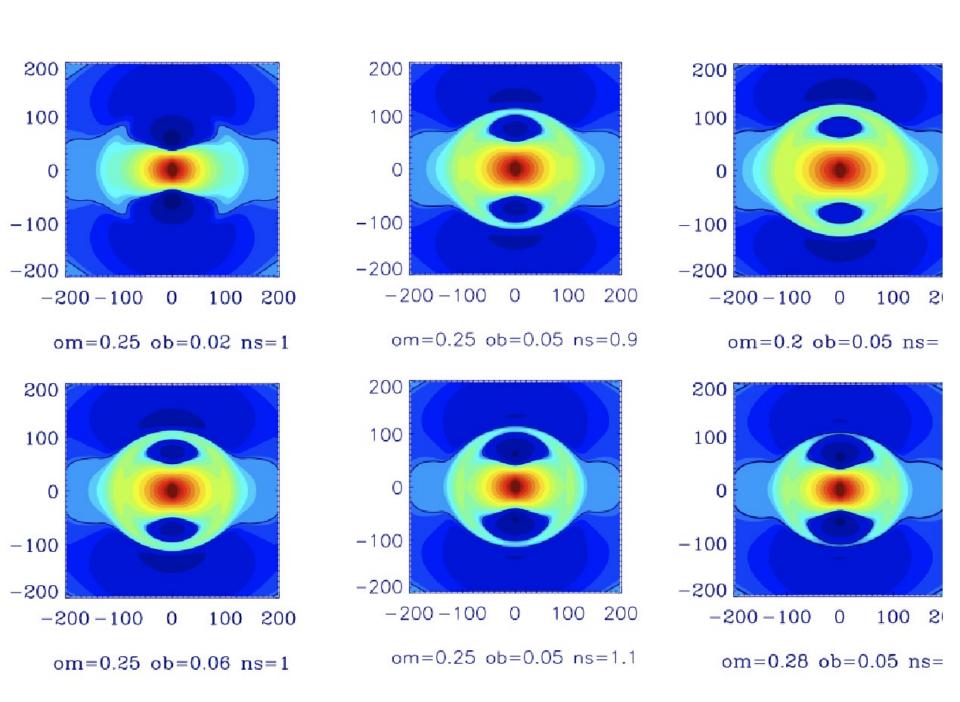
From Eisenstein et al. (2007)

Readshift space distortion

- The distortion depends on the density and velocity fields
- They can be modelled









project web: www.ice.cat/mice

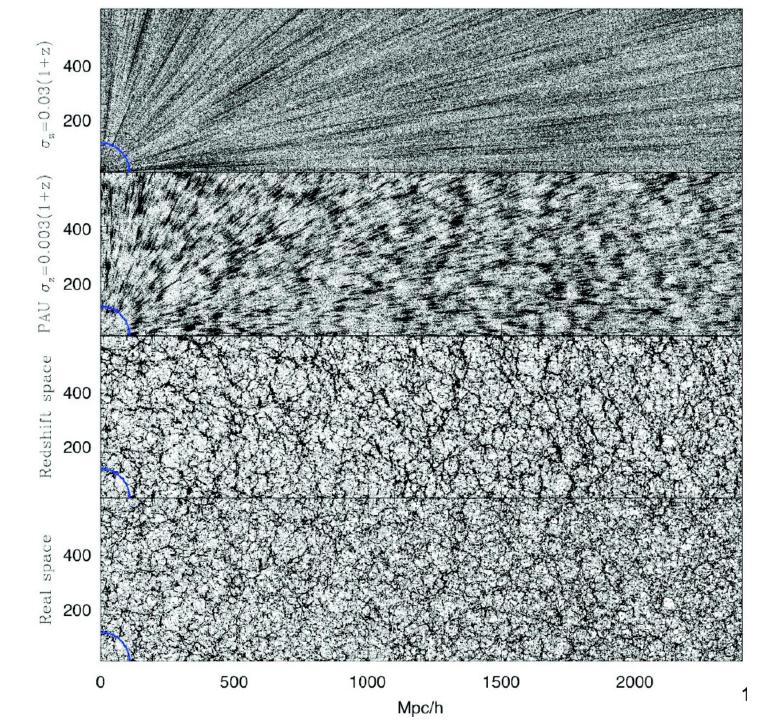
What is MICE?

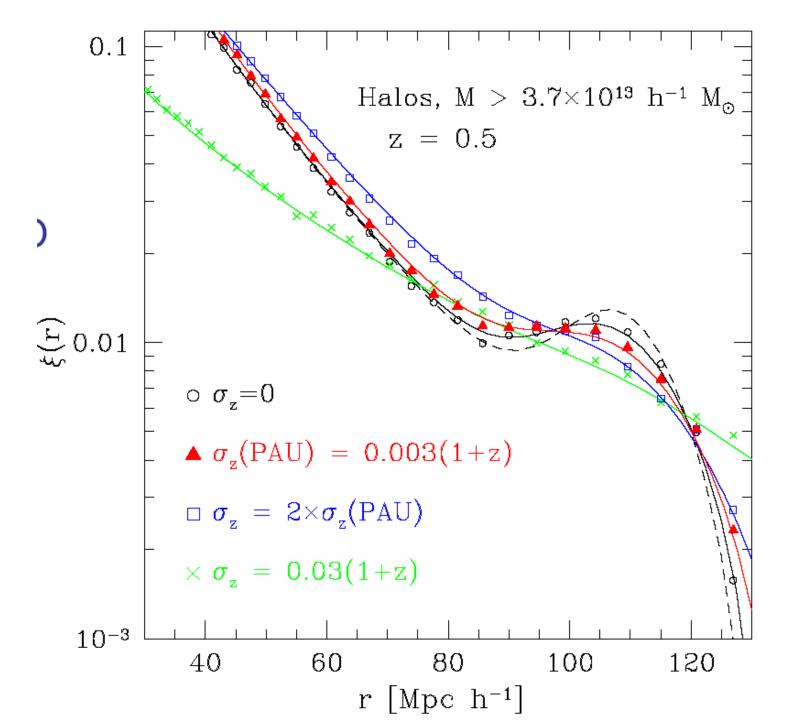
- ➤ Project to develop very large numerical simulations in cosmology using the **Marenostrum** supercomputer (Barcelona)
- **★** 10.000 processors, 20 TB RAM , 100 Teraflops
- **≰** GADGET N-body simulations with 10⁹-10¹⁰ dark-matter particles in volumes 1-500 Gpc³
 - ⇒ dynamical range of 5 orders of magnitude

➤Team:

Core (ICE): P.Fosalba, F.Castander, E.Gaztañaga Collaborators: V.Springel (MPA), C.Baugh (Durham), M.Manera (NYU), M.Crocce, A.Gonzalez, A.Cabré, J.Miralda-Escudé

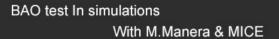


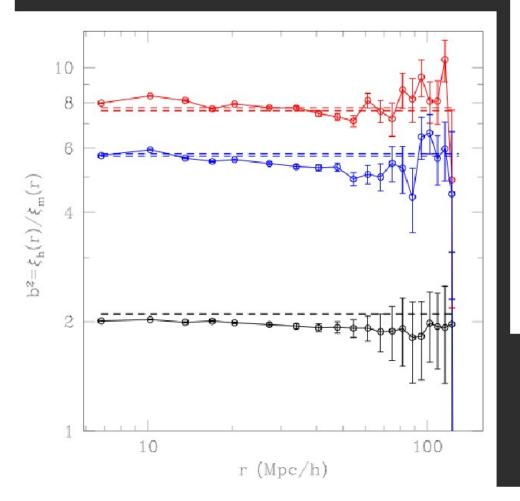


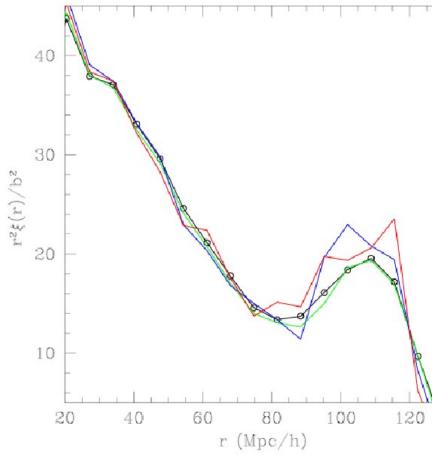


Galaxy bias

- Galaxies are a bias tracer of the underlying dark matter distribution
- Uncertainties in peak determination if bias is scale dependent
- The expectation is that at the BAO scale the effect of bias should depend on scale in a smooth manner







How robust is the BAO method?

- Recent claims that the BAO scale is not fixed but depends on non-linearities, bias,...
- e.g., Smith et al 2008, Sanchez et al 2007

Final Considerations

- What method is better suited: configuration space or Fourier space?
- Comparison observations-theory (simulations)
- Future surveys: sample variance limited $\frac{\sigma}{P} = \sqrt{\frac{2}{n_{\text{modes}}}} \left(1 + \frac{1}{P\bar{n}}\right)$
- Error determination -> simulations

Conclusions

- the BAO signature provides a standard ruler that can be used to measure the geometry of the universe
- \bullet it can measure both the angular diameter distance $D_{\mbox{\tiny A}}(z)$ and the expansion rate H(z)
- it has already been measured
- it is fairly robust against systematics