

Power Spectrum of the CLB

Gil Holder

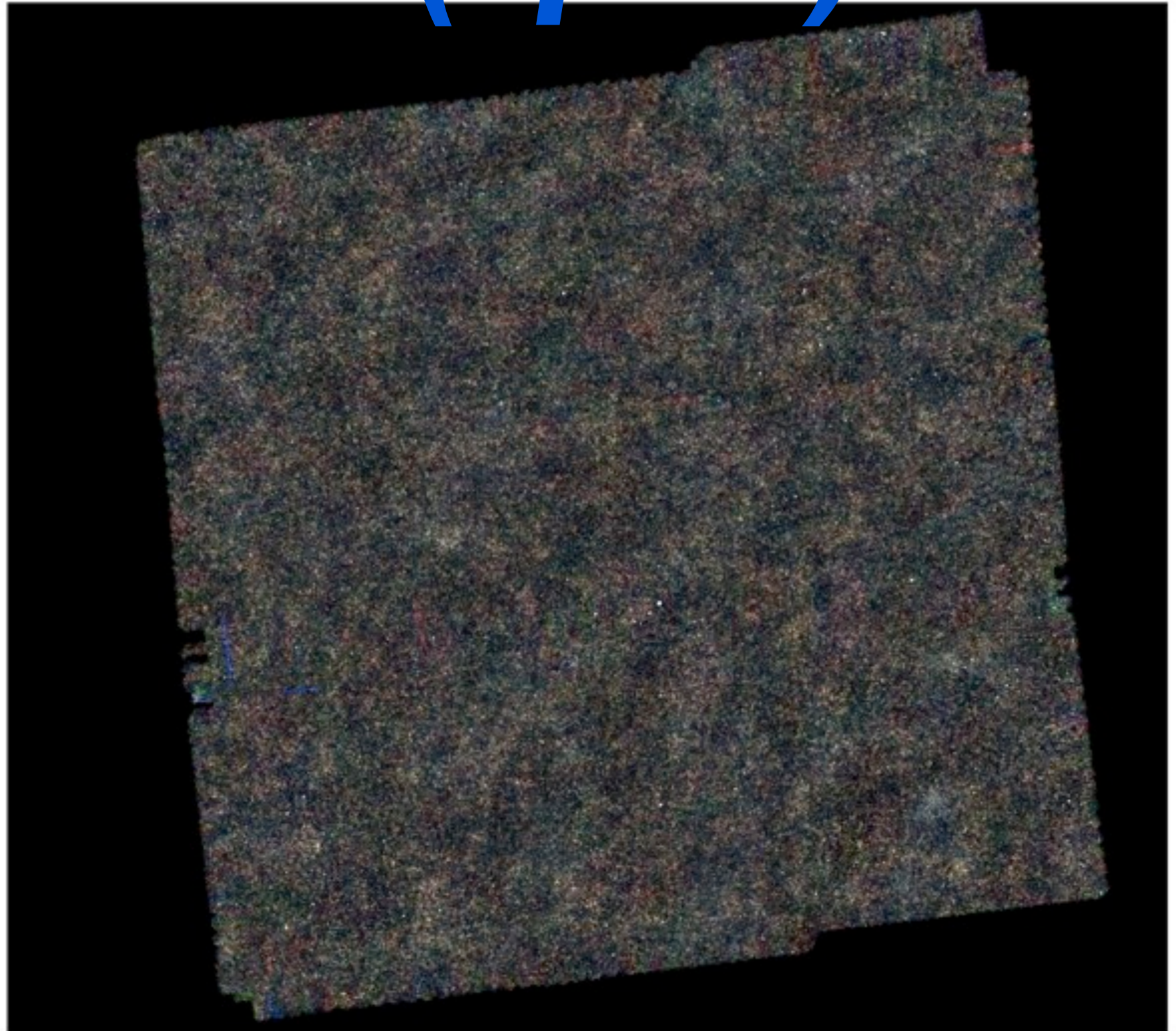
McGill University (Montreal, Canada)

Outline

- power spectra principles
 - ★ flat sky approximation
 - ★ cross-spectra
 - ★ diversion on cross-linking
- incomplete overview of CIB detections

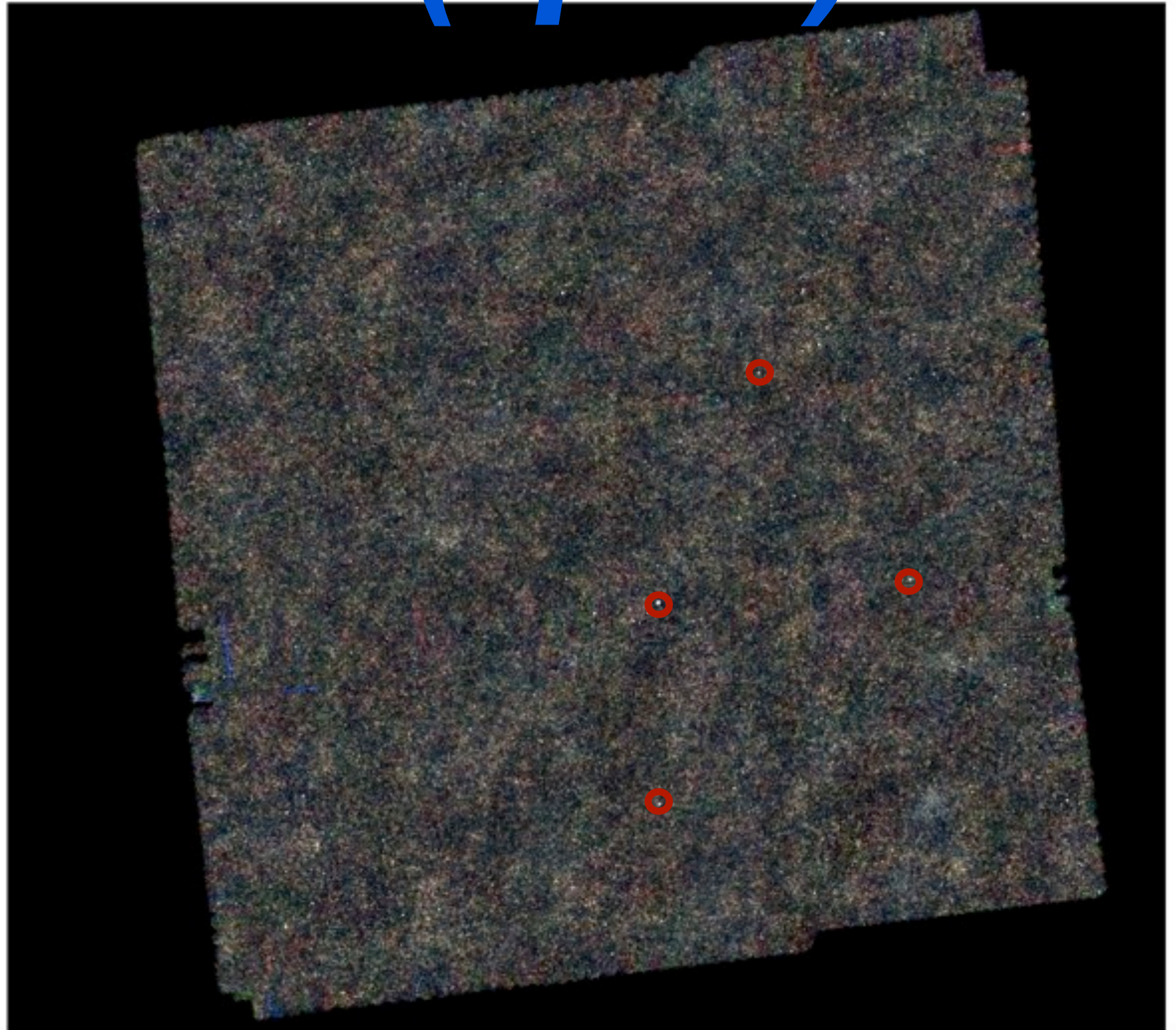
Herschel (Spire)

- 100 sq deg with full overlap with SPT deep field (23h30,-55d)
- 250,350,500 μm



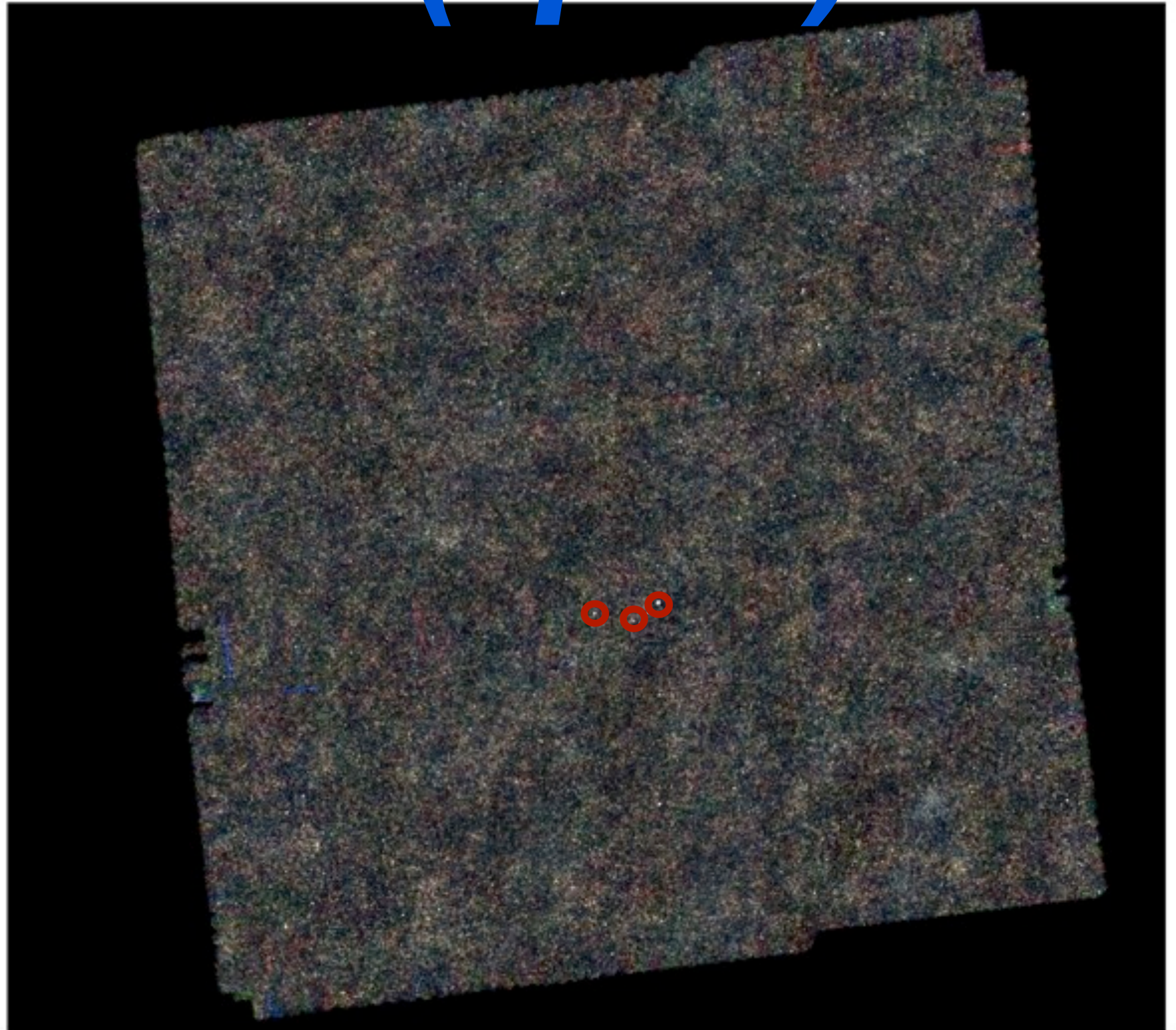
Herschel (Spire)

- can do source counts
- 1 pt probabilities



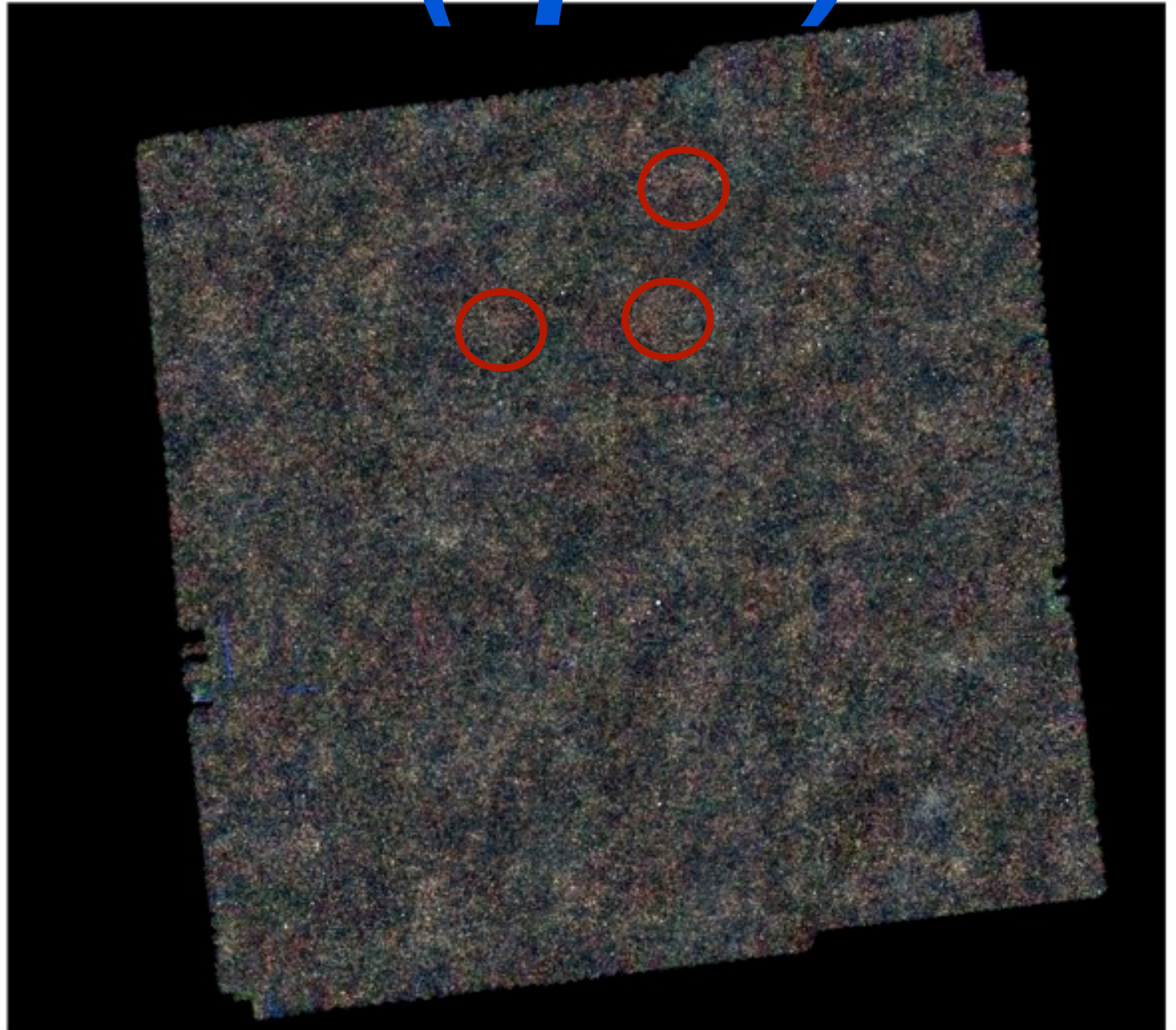
Herschel (Spire)

- are sources too close to each other?
- 2 pt probabilities



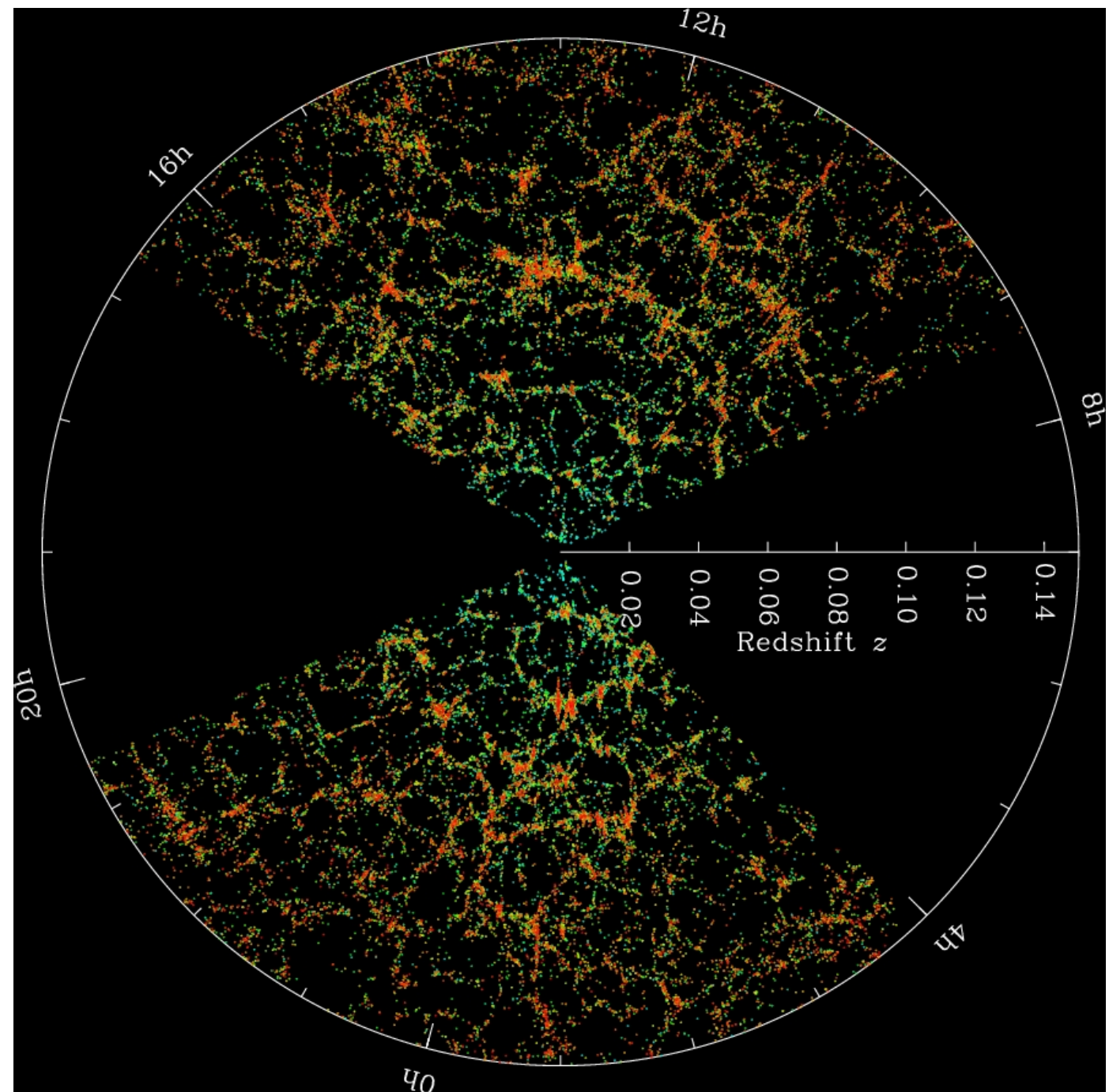
Herschel (Spire)

- what about large scale structure in the map?
- general plotchiness in the map seems to not be fully captured by bright source counts



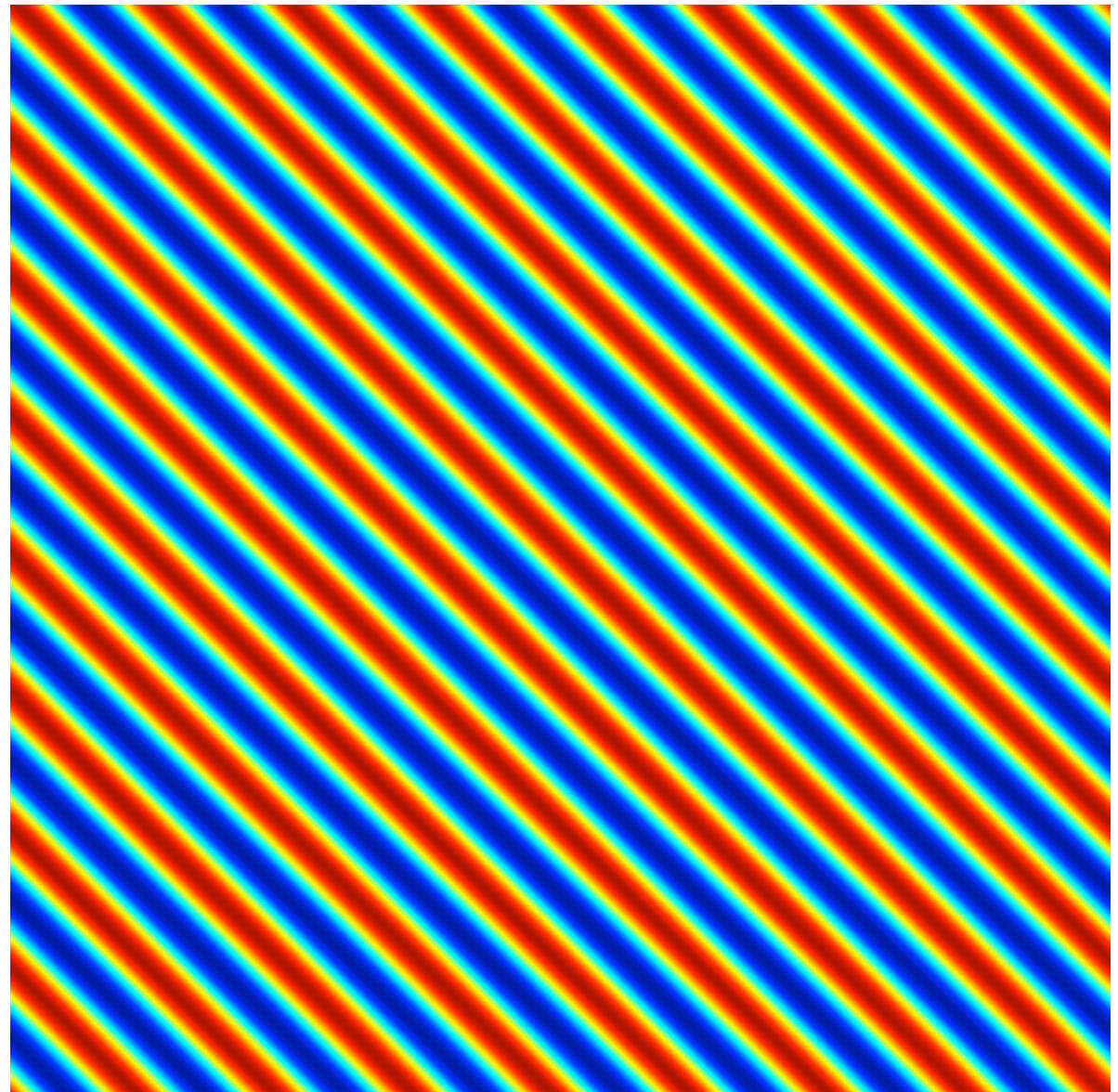
Clustering of Point Sources

- Radio and IR/submm sources presumably trace the large scale matter fluctuations
- Back of the envelope:
 - Power spectrum contribution: mean T^2 x projected clustering amplitude
 - Arcminute scales: few Mpc has clustering ~ 1 in 3D, divide by number of independent cells along line of sight $\Rightarrow 1e-3$
 - end result $dT/T \sim 10\%$



Poisson Power Spectrum

- single point source leads to a corrugation in Fourier space
- sum over N sources:
 $T^2 \propto N S^2$
- no preferred scales:
white noise

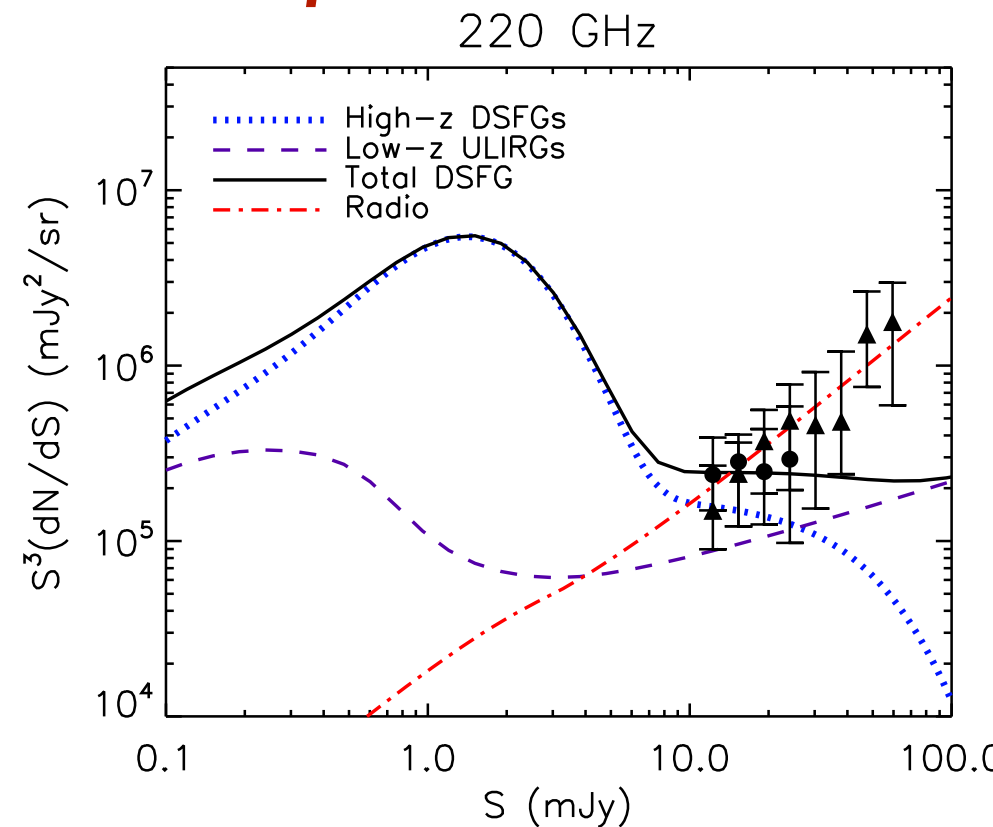
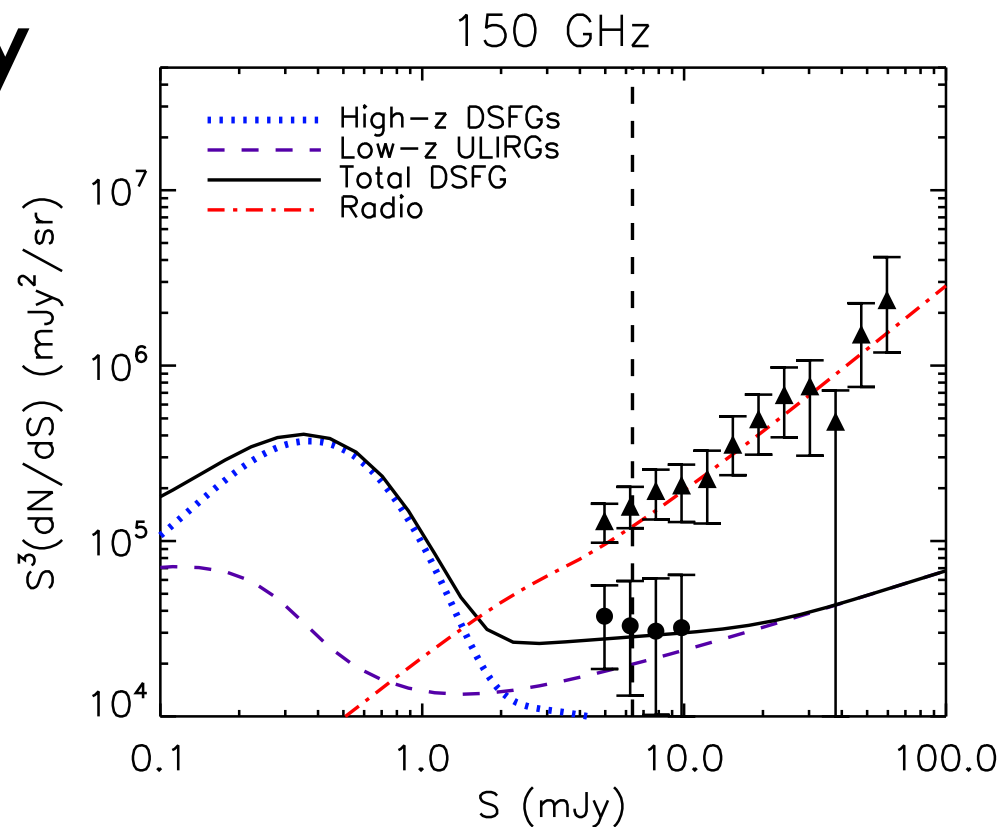


Real part of Fourier transform of 1 source

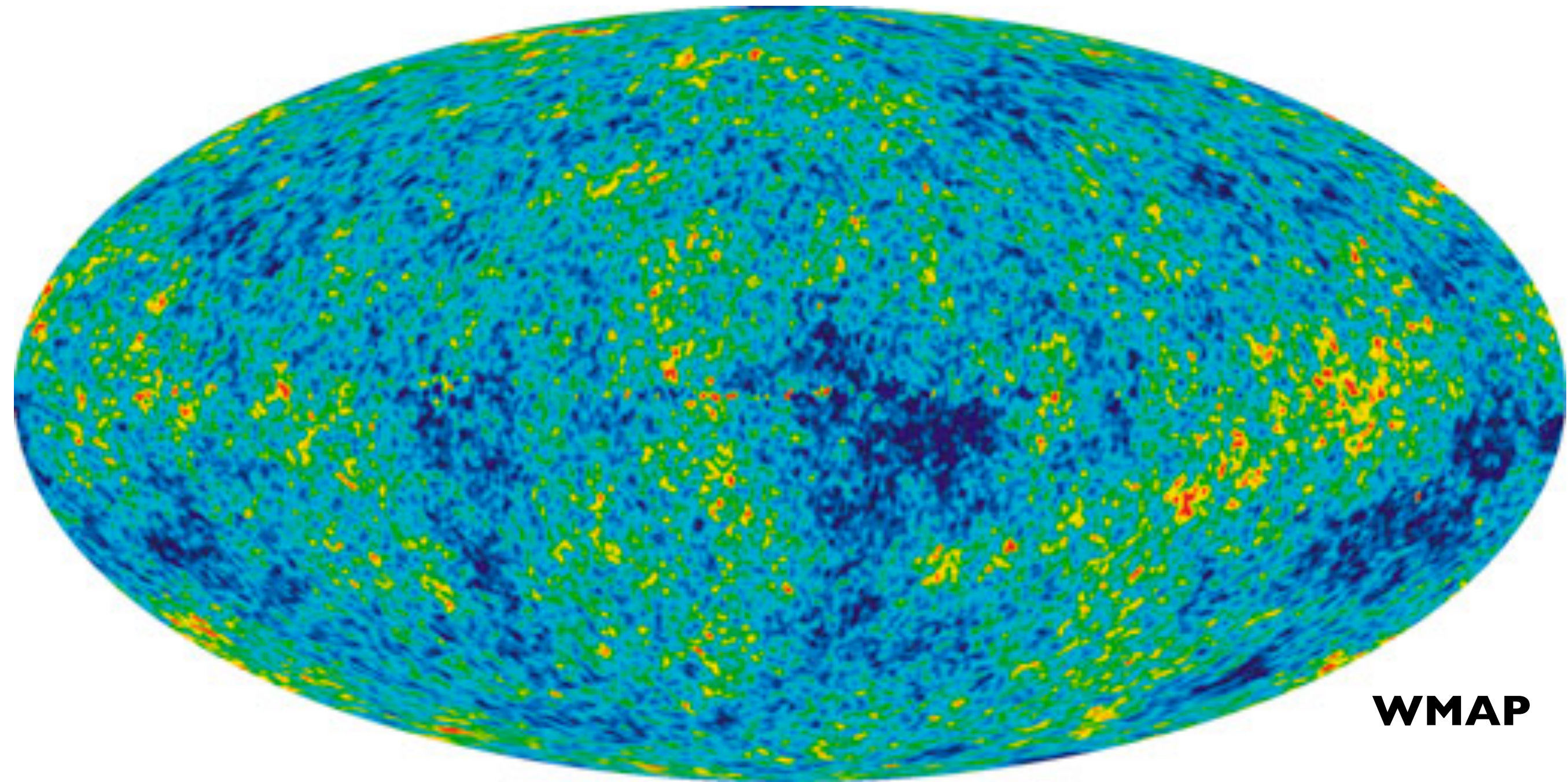
Which Sources?

- CIB mainly comes from dim sources (faint end of classic SMGs?)

Contributions to Poisson power



Hall et al 2010



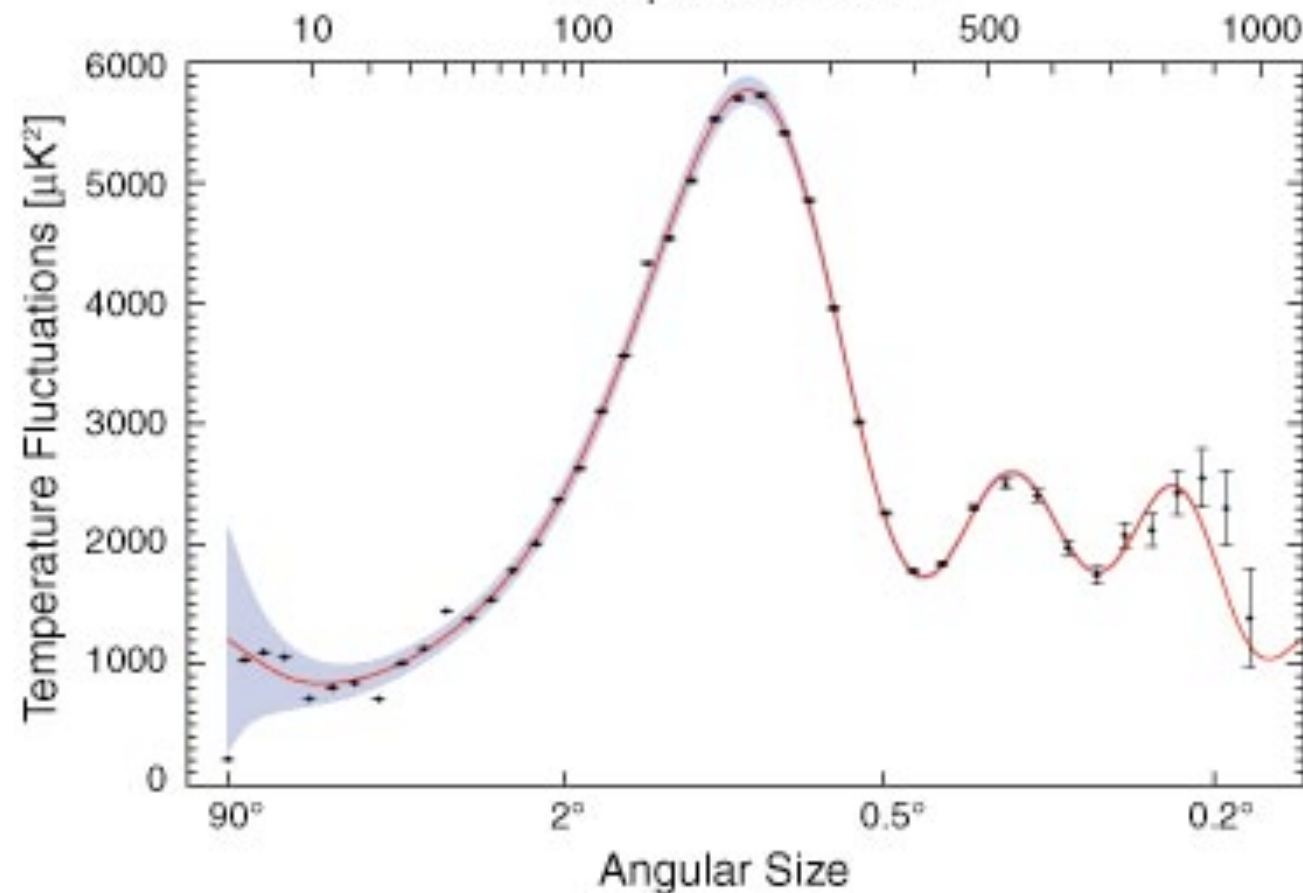
What do we do with CMB maps?

Spherical Harmonics

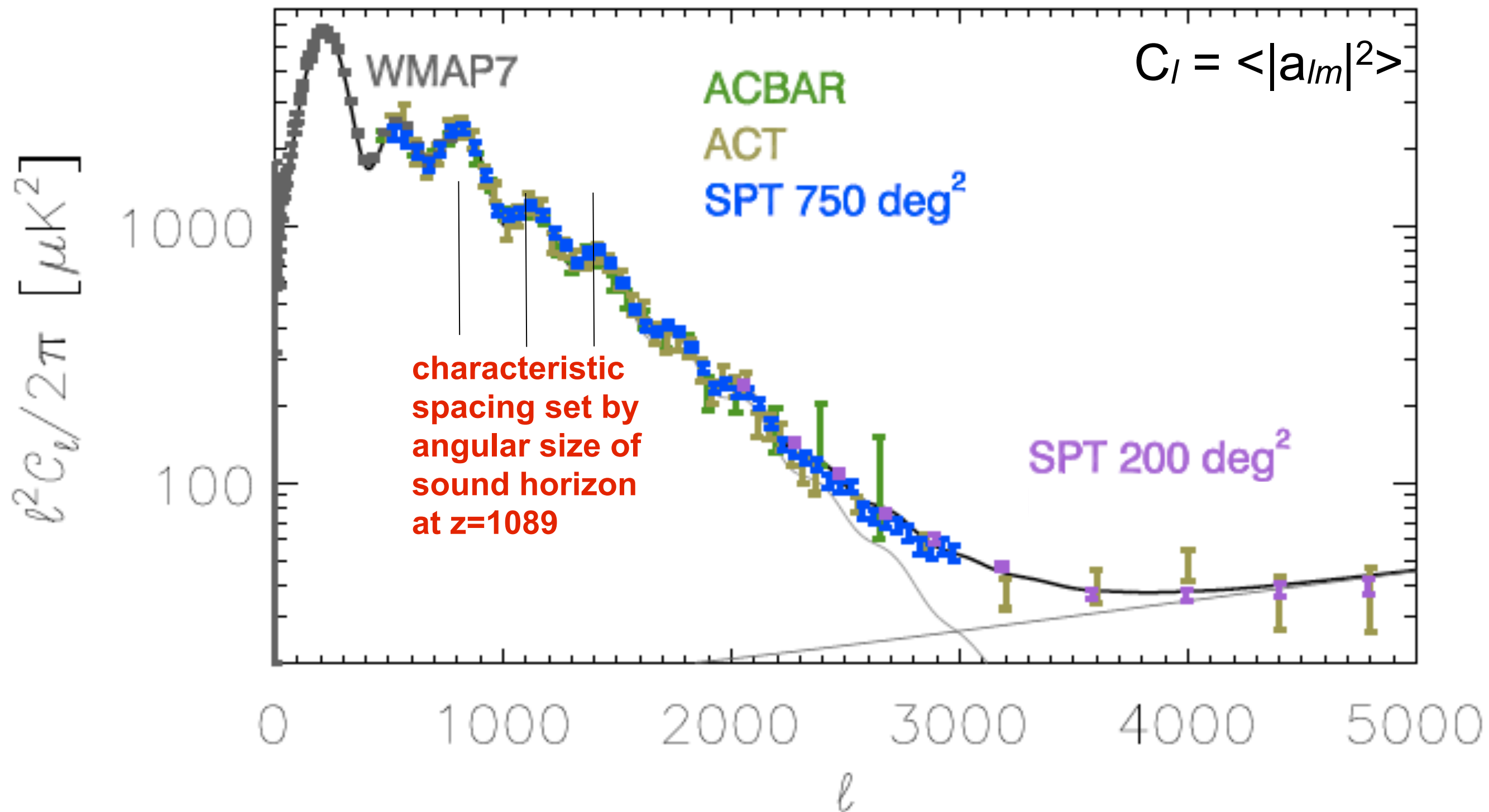
$$T(\theta, \phi) = \sum_{\ell m} a_{\ell m} Y_{\ell m}(\theta, \phi).$$

$$Y_{\ell}^m(\theta, \varphi) = \sqrt{\frac{(2\ell + 1)}{4\pi} \frac{(\ell - m)!}{(\ell + m)!}} P_{\ell}^m(\cos \theta) e^{im\varphi}$$

Multipole moment l



CMB Power Spectrum



Fourier Transforms

$$s(t) = \int_{-\infty}^{\infty} S(f) e^{i2\pi ft} df$$

Does one have to choose between Fourier coefficients and spherical harmonics?
(no)

From a_{lm} to $a_{k_x k_y}$

Equation satisfied by $P_{lm}(x)$:

$$(1 - x^2) y'' - 2xy' + \left(\ell[\ell + 1] - \frac{m^2}{1 - x^2} \right) y = 0,$$

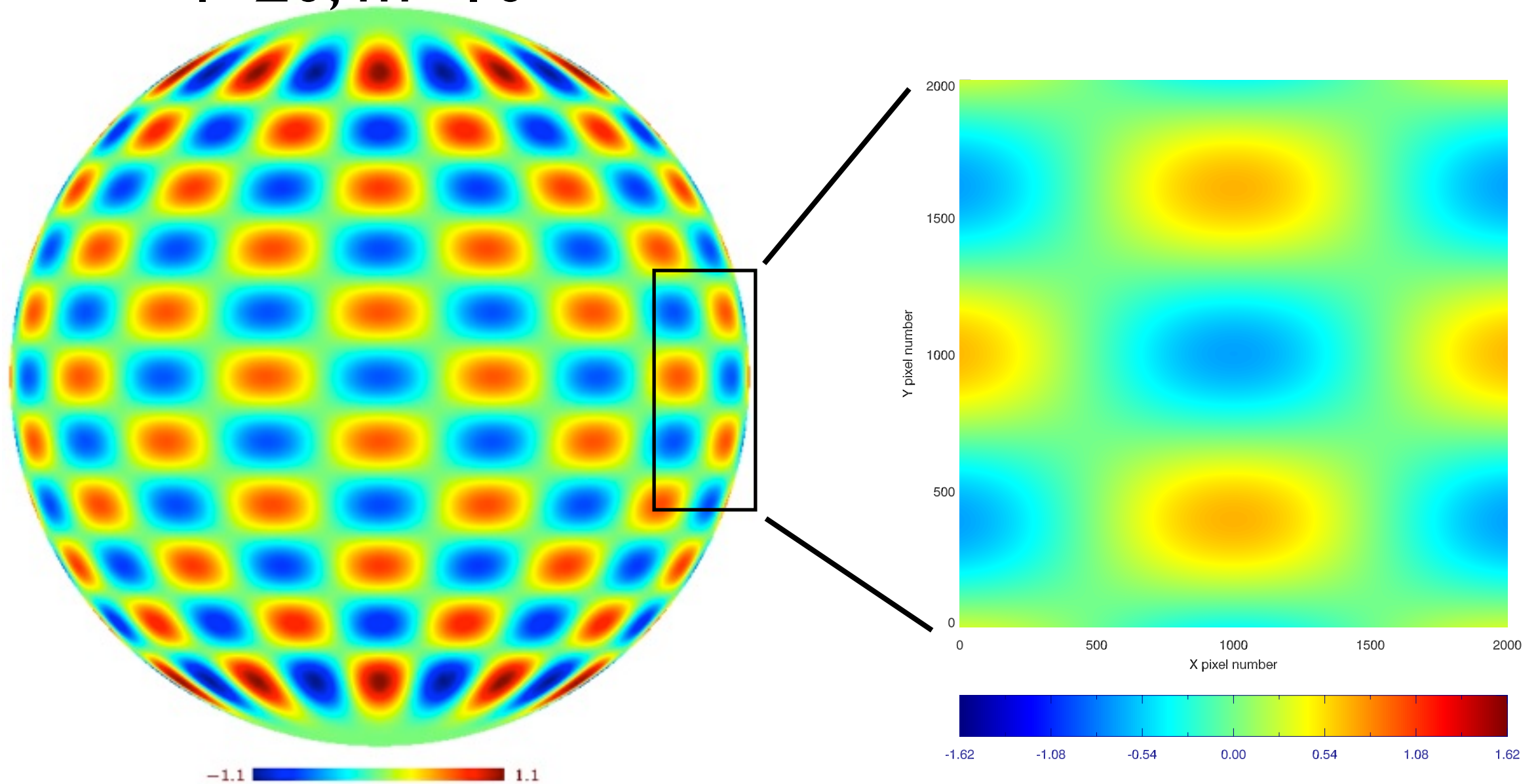
For $x \sim 0$:

$$(1 - \cancel{x^2}) y'' - 2\cancel{x}y' + \left(\ell[\ell + 1] - \frac{m^2}{1 - \cancel{x^2}} \right) y = 0,$$

Harmonic Oscillator with **$k = \ell(\ell + 1) - m^2$**
(Fourier modes!)

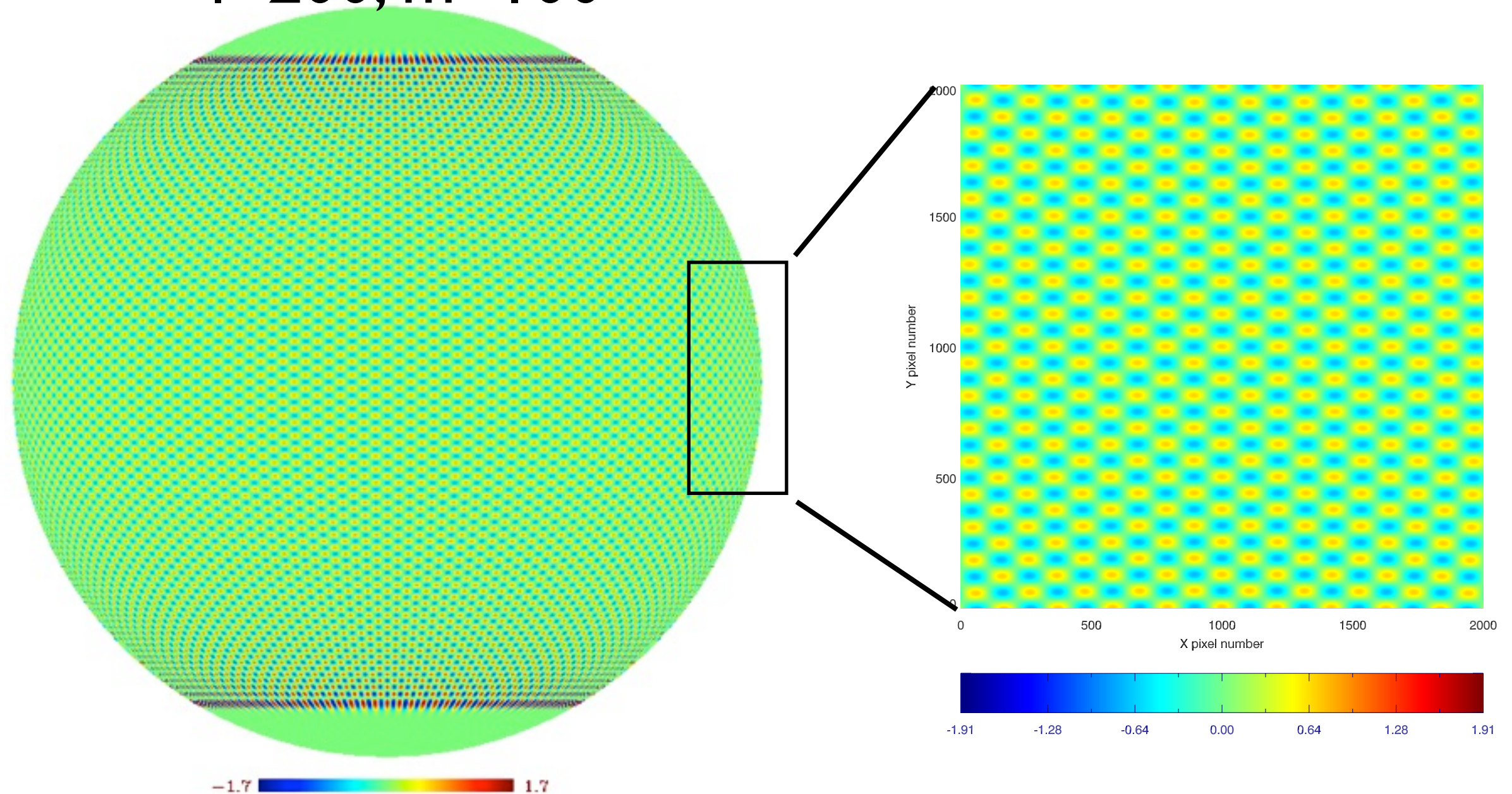
Projecting a_{lm}

on line processing :
 $l=20, m=10$



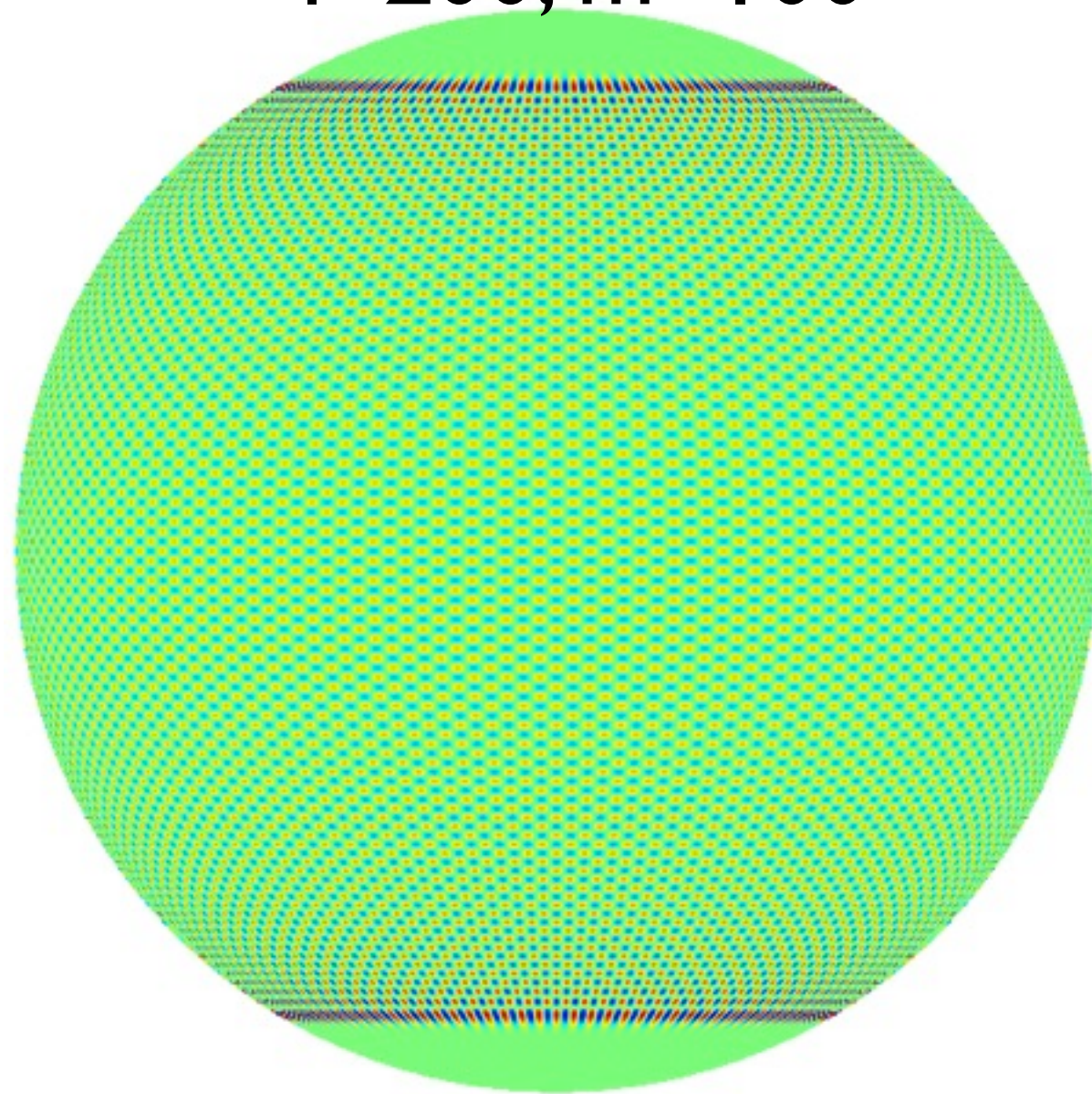
Projecting a_{lm}

on line processing :
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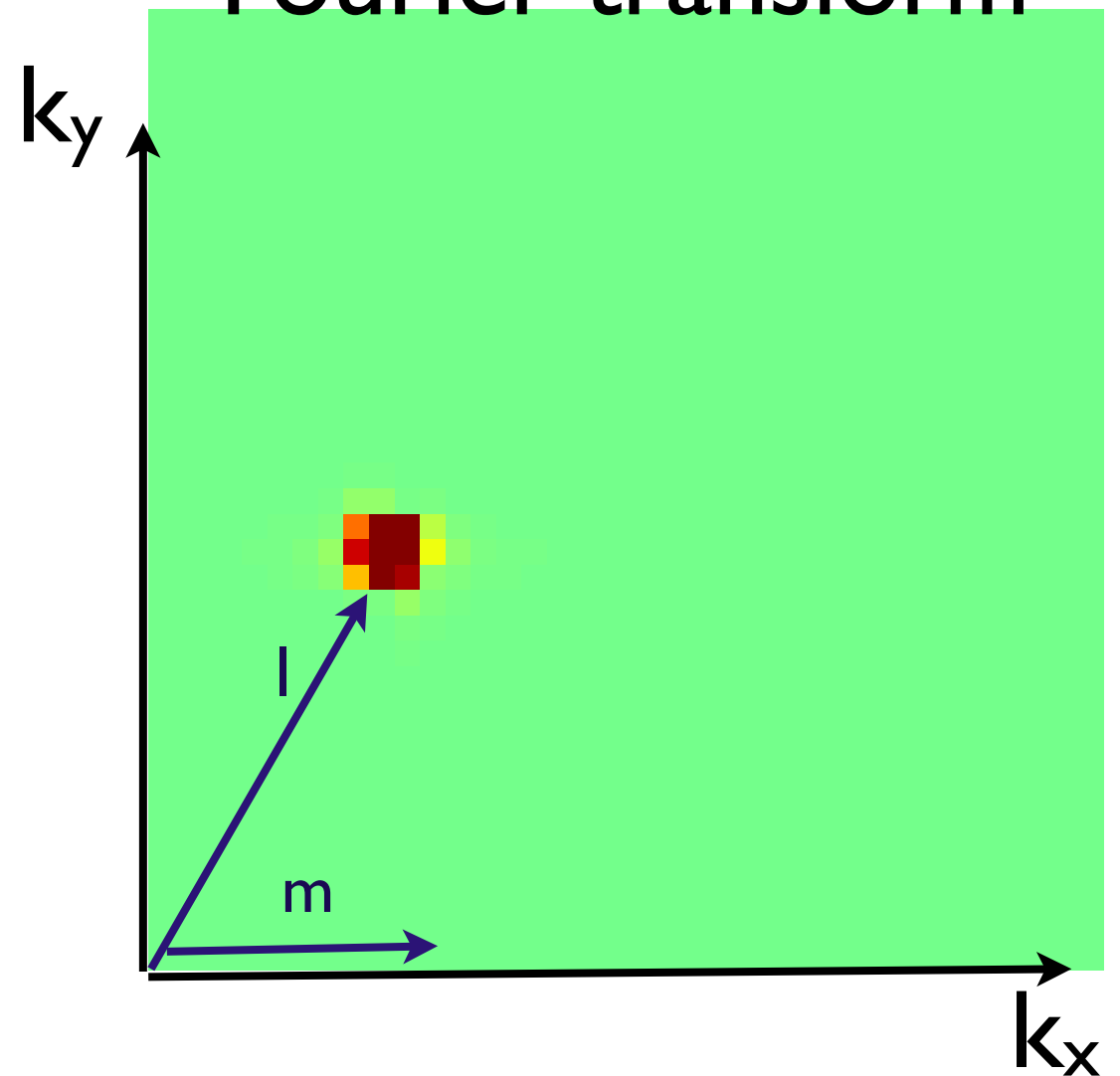
Projecting a_{lm}

on line processing :
 $l=200, m=100$



-1.7 1.7

Fourier transform



“Flat Sky” vs multipoles

- for chunk of map extracted at equator, the two are basically identical for all but the largest angles ($l < 30$)
- choice of pole shouldn't matter (statistical isotropy), so Fourier or multipole, up to you
- $l \sim 2l600 k_\varphi (\text{arcmin}^{-1})$

Power spectrum Uncertainties

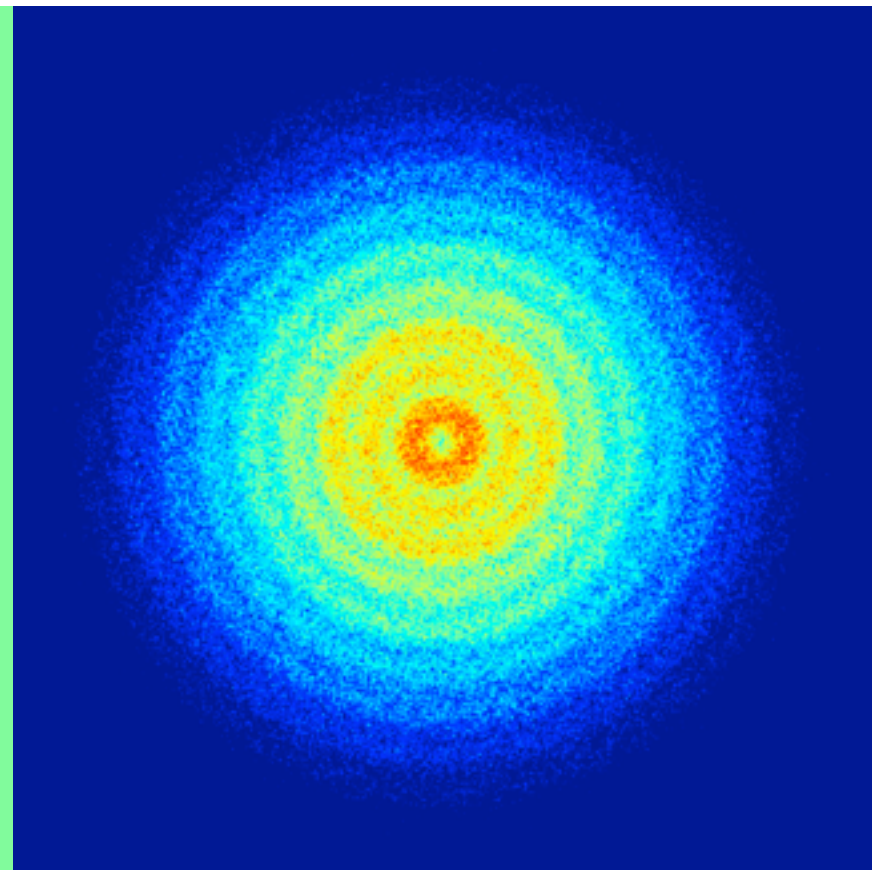
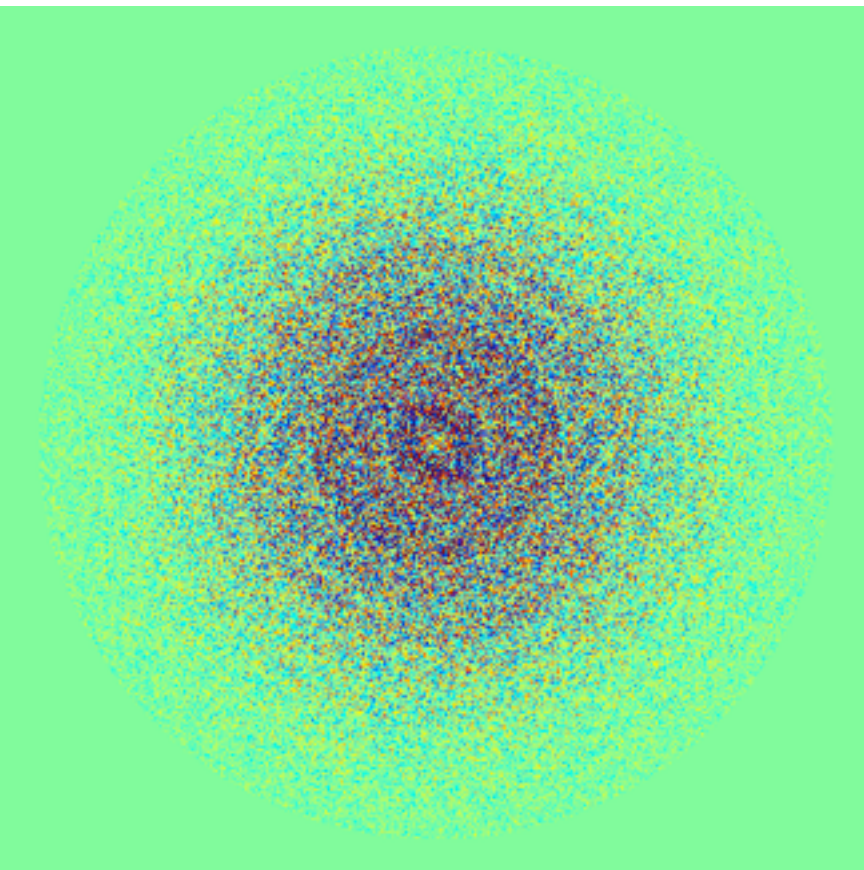
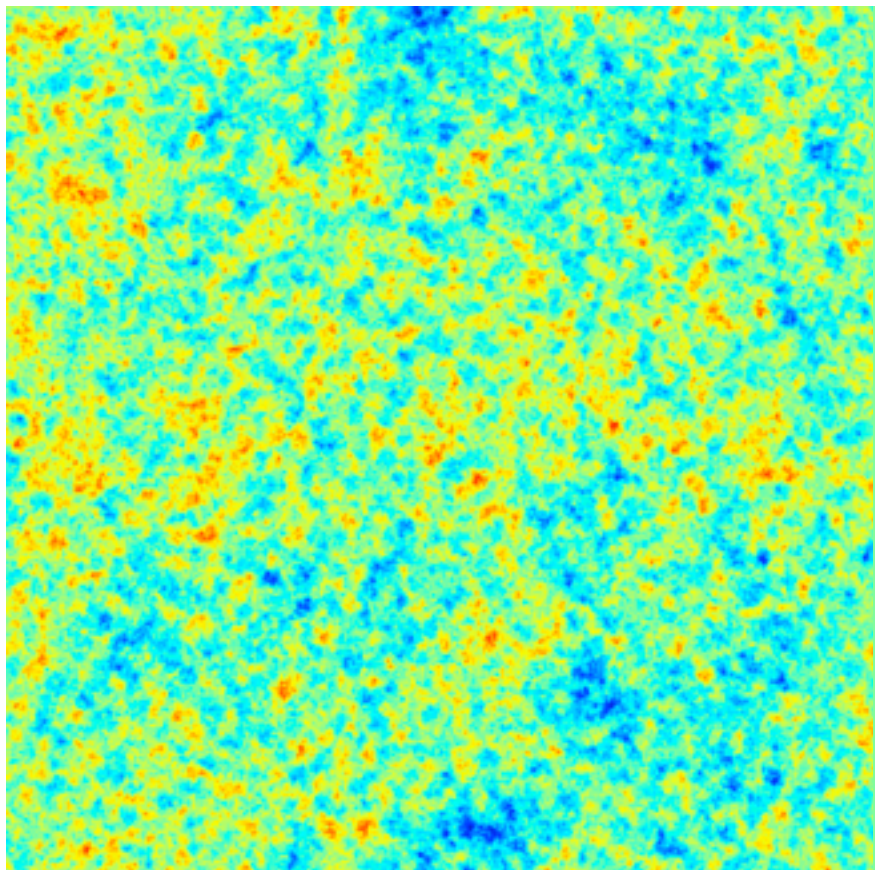
- fundamentally limited by number of independent measurements, noise
- $C_{l;\text{meas}} = C_{l;\text{true}} + C_{l;\text{noise}}$ *in any single map you can't tell the difference*
- $\text{Var}(C_l) \sim 2/n_{\text{meas}} C_l^2$ **“sample variance”**
- more modes means better measurement of $C_{l;\text{true}} + C_{l;\text{noise}}$
- lower noise gives better measure of $C_{l;\text{true}}$

Simulated CMB

Map

Real part of FFT

$\log_{10}(|\text{FFT}|^2)$

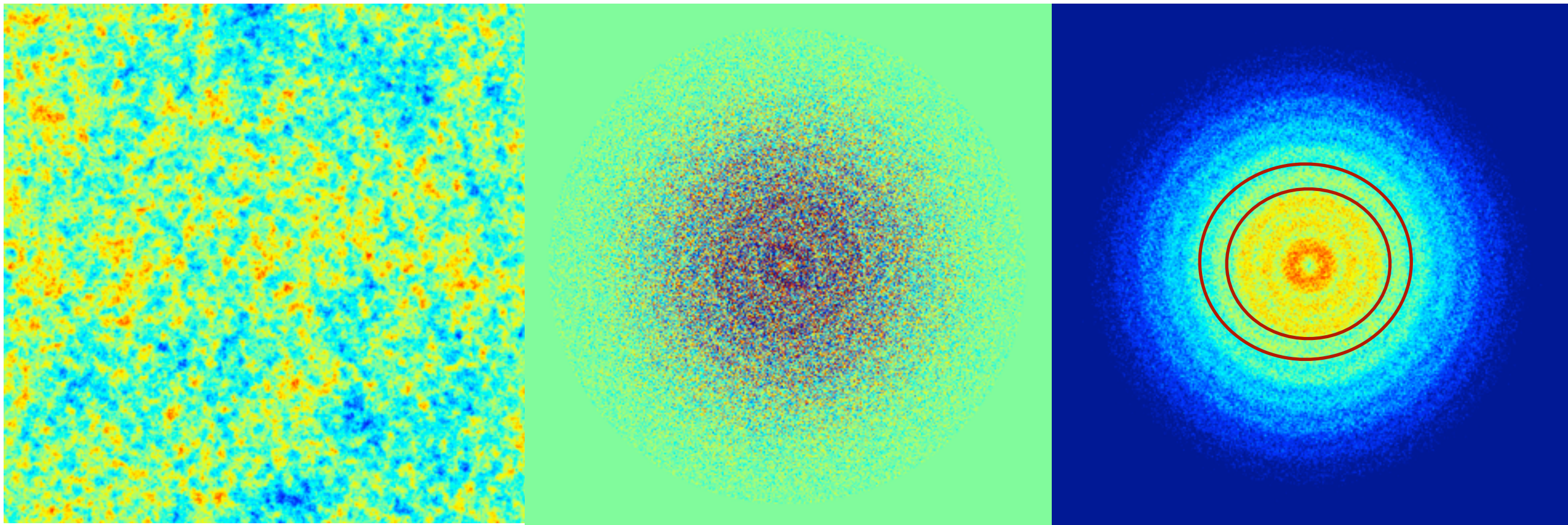


Simulated CMB

Map

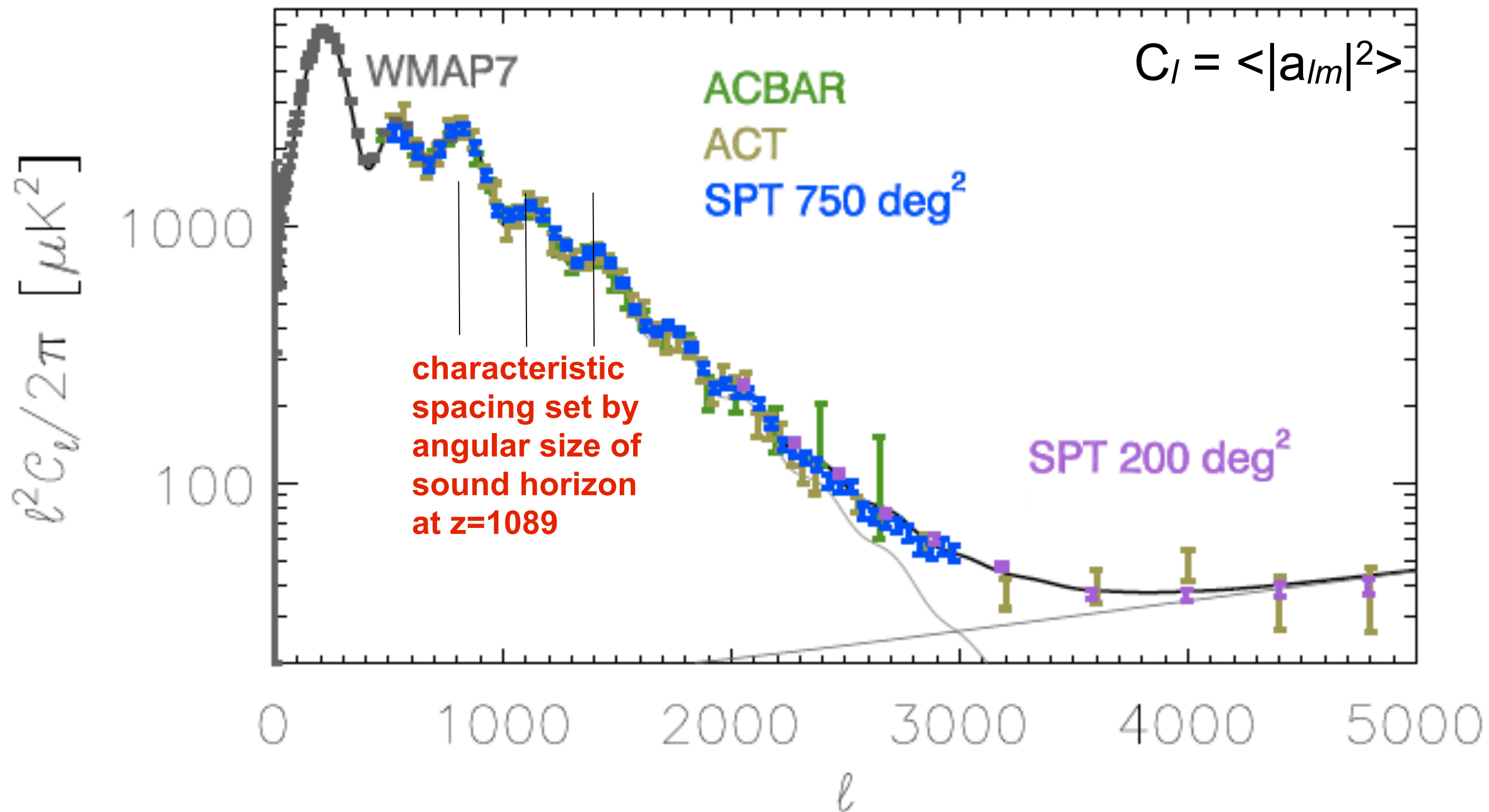
Real part of FFT

$\log_{10}(|\text{FFT}|^2)$



*# of independent
samples set by map size*

CMB Power Spectrum



Cross Spectra

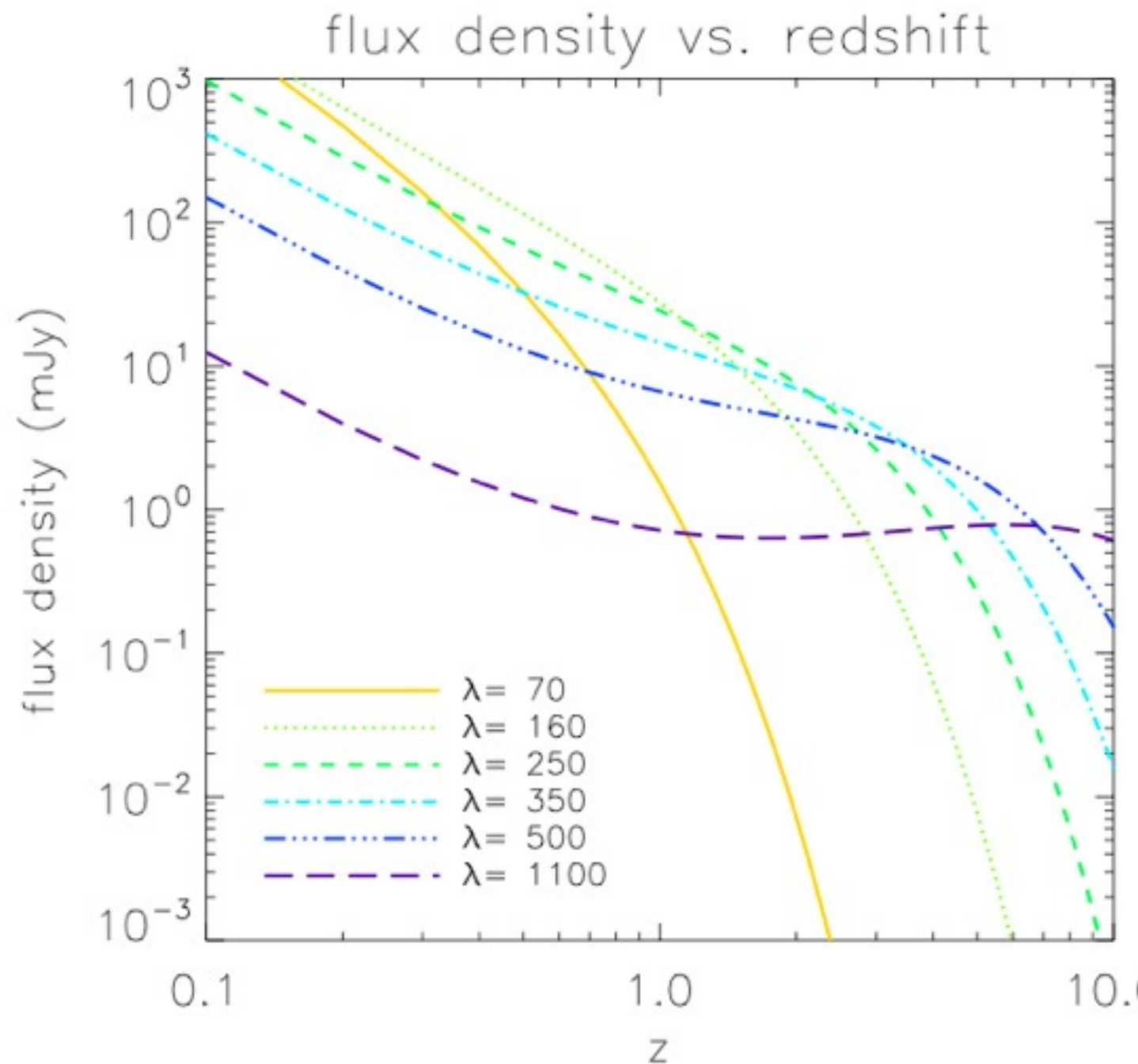
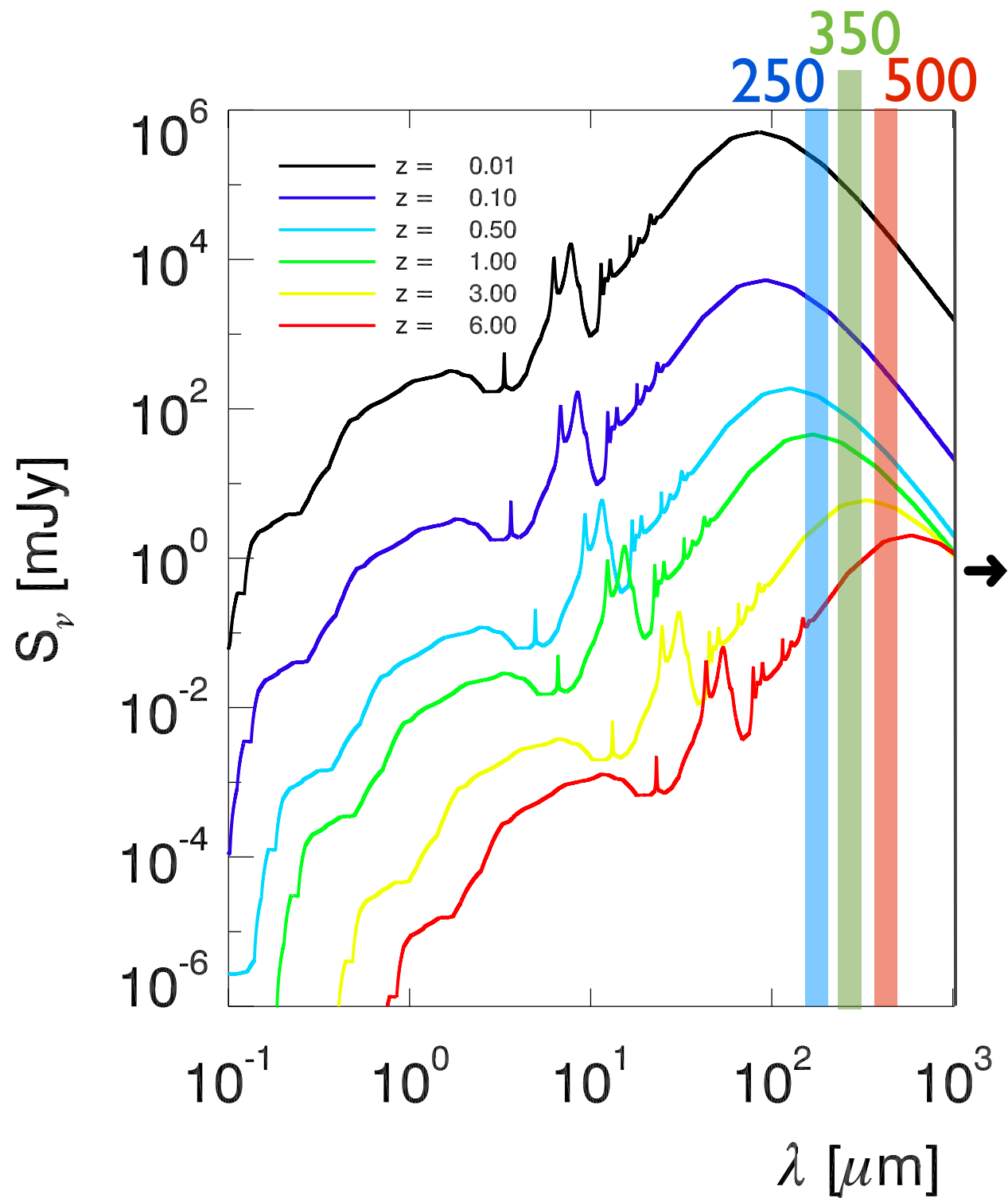
all quantities are Fourier space!

- $T_m = T + n$
- $\langle T_{1;m} T_{2;m} \rangle = \langle T_1 T_2 \rangle + \langle n_1 n_2 \rangle + \langle T_1 n_2 + T_2 n_1 \rangle$
- for $1=2$ (map auto power spectrum), $\langle n_1 n_2 \rangle = \sigma^2$
- if $1 \neq 2$, $\langle n_1 n_2 \rangle = 0$, so no bias
- quirks in your noise model don't affect cross spectrum!

Cross-Spectra in Action

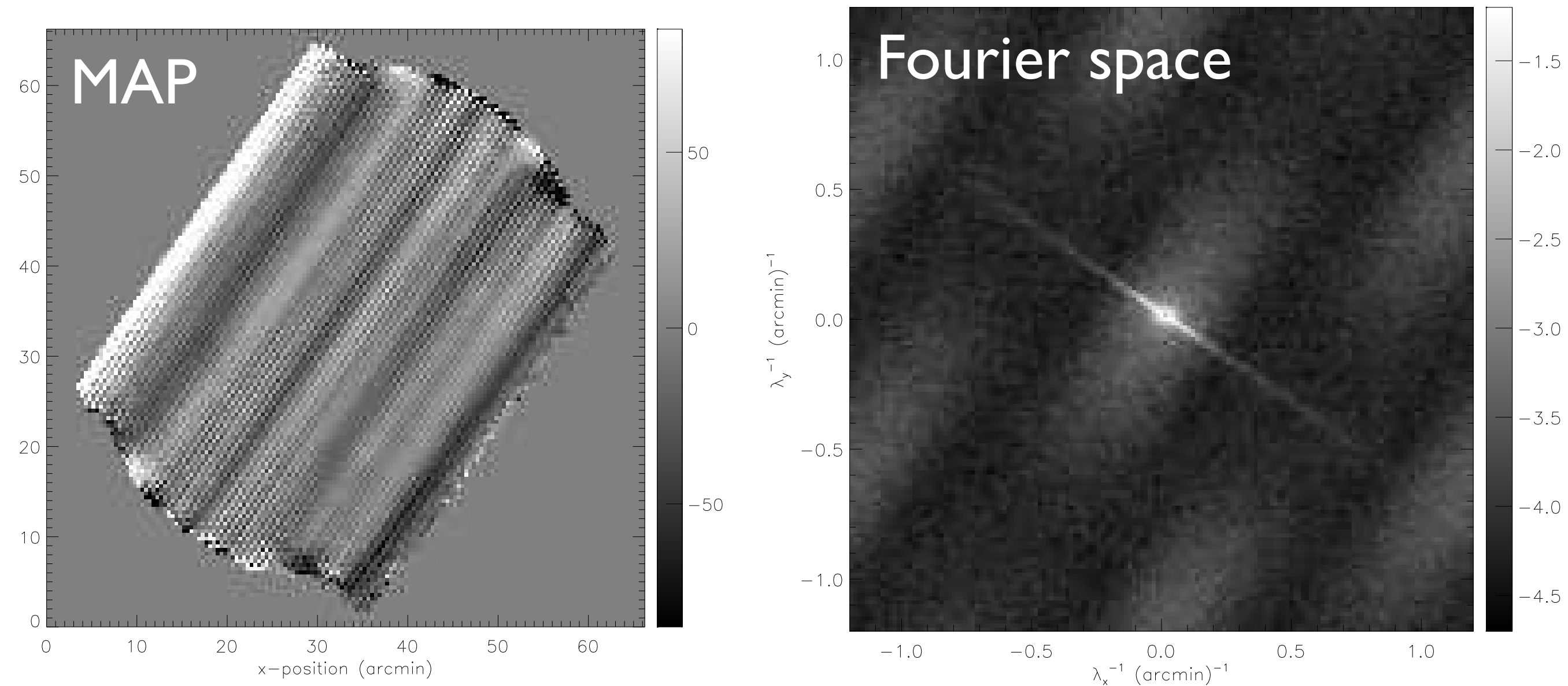
- if you have mapped a field 2 times, you can combine them: $T=(T_1+T_2)$
- the autospectrum of T is $\sim T_1T_1+T_2T_2+2T_1T_2$
- but $T_1T_1+T_2T_2$ have noise bias, so with a hit in sensitivity you can ignore these
- in the limit of large number of maps, hit in sensitivity goes to 0 while robust against problems of not knowing your detector noise

Multi-frequency is good!



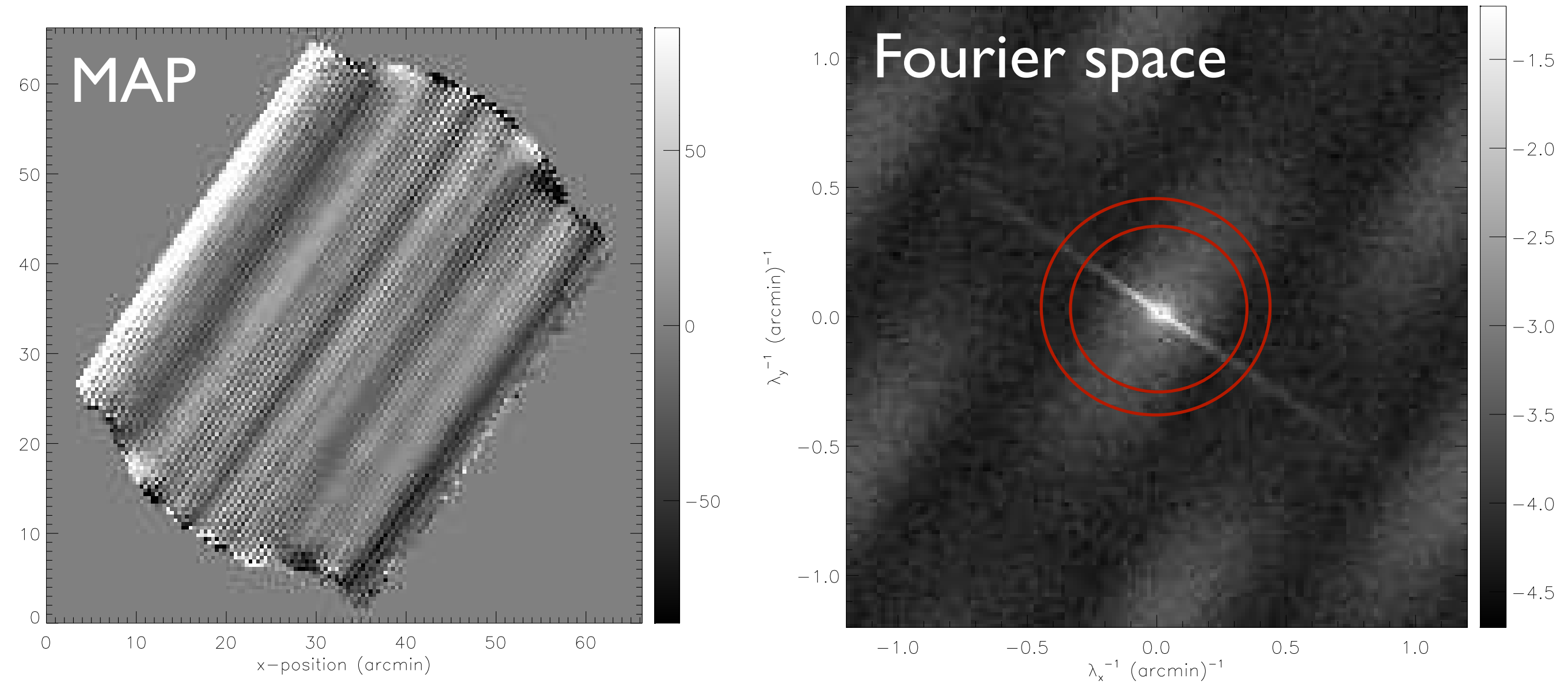
negative K -correction

Power Spectra: Anisotropic Noise

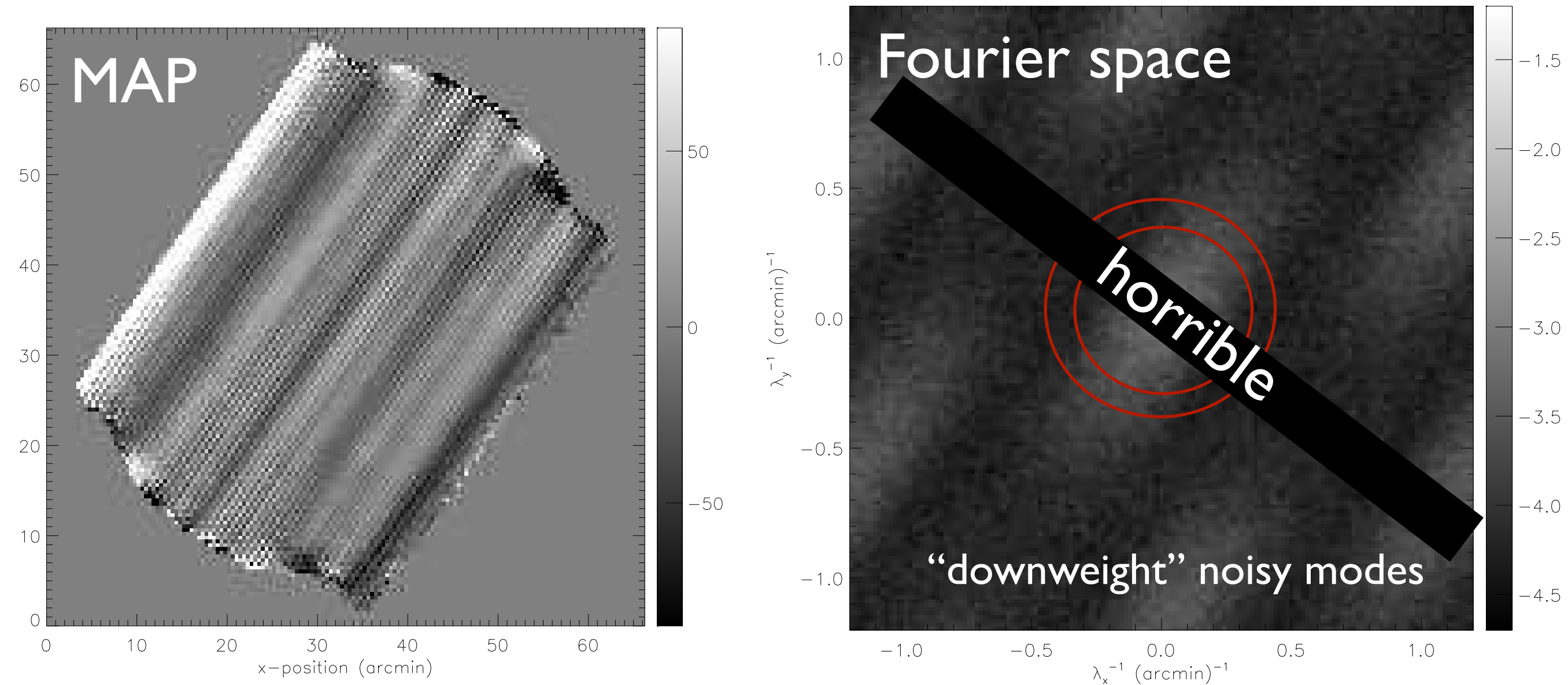


simulated BLAST noise map; Patanchon et al 07 | I.3463

Power Spectra: Anisotropic Noise

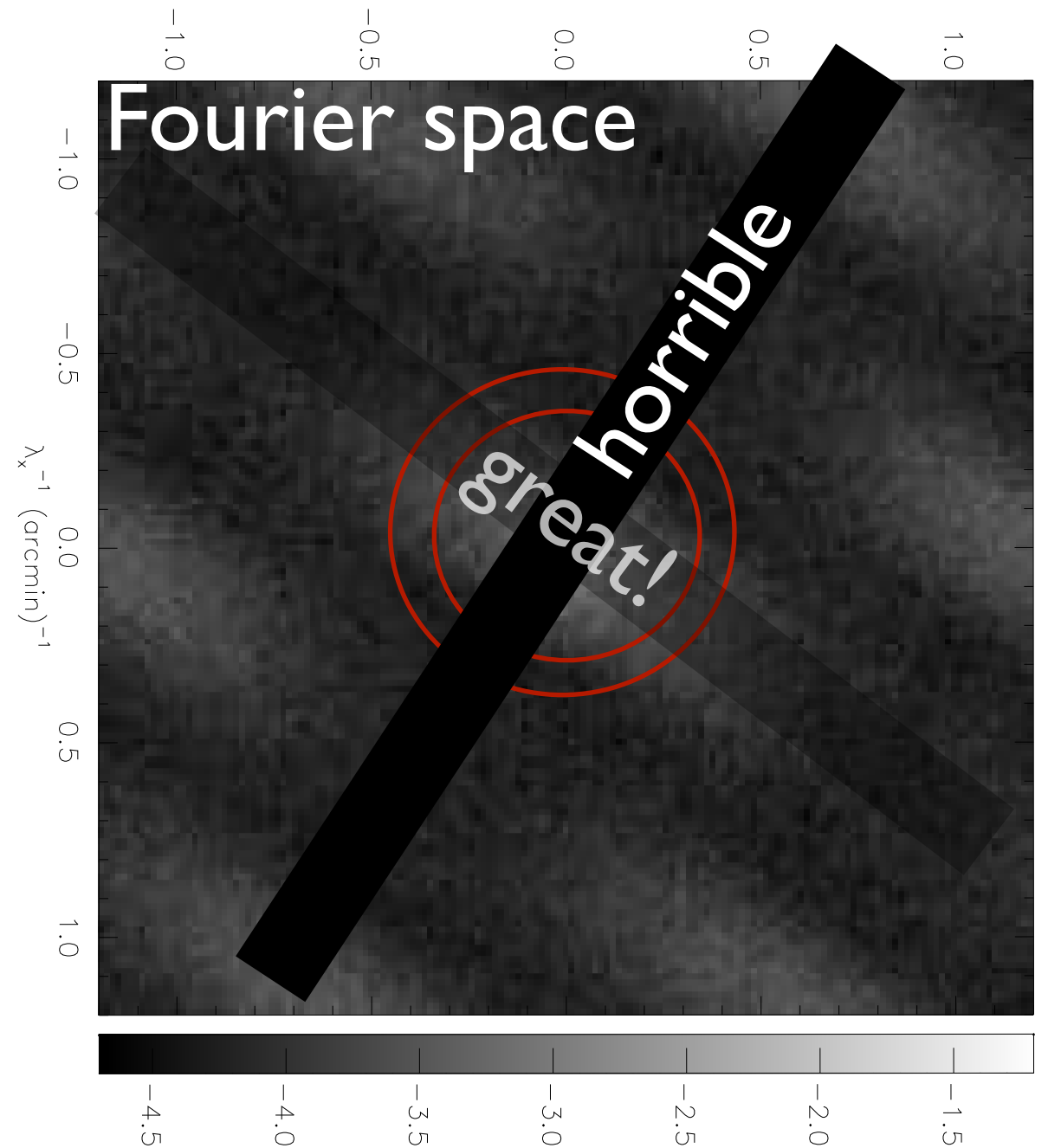
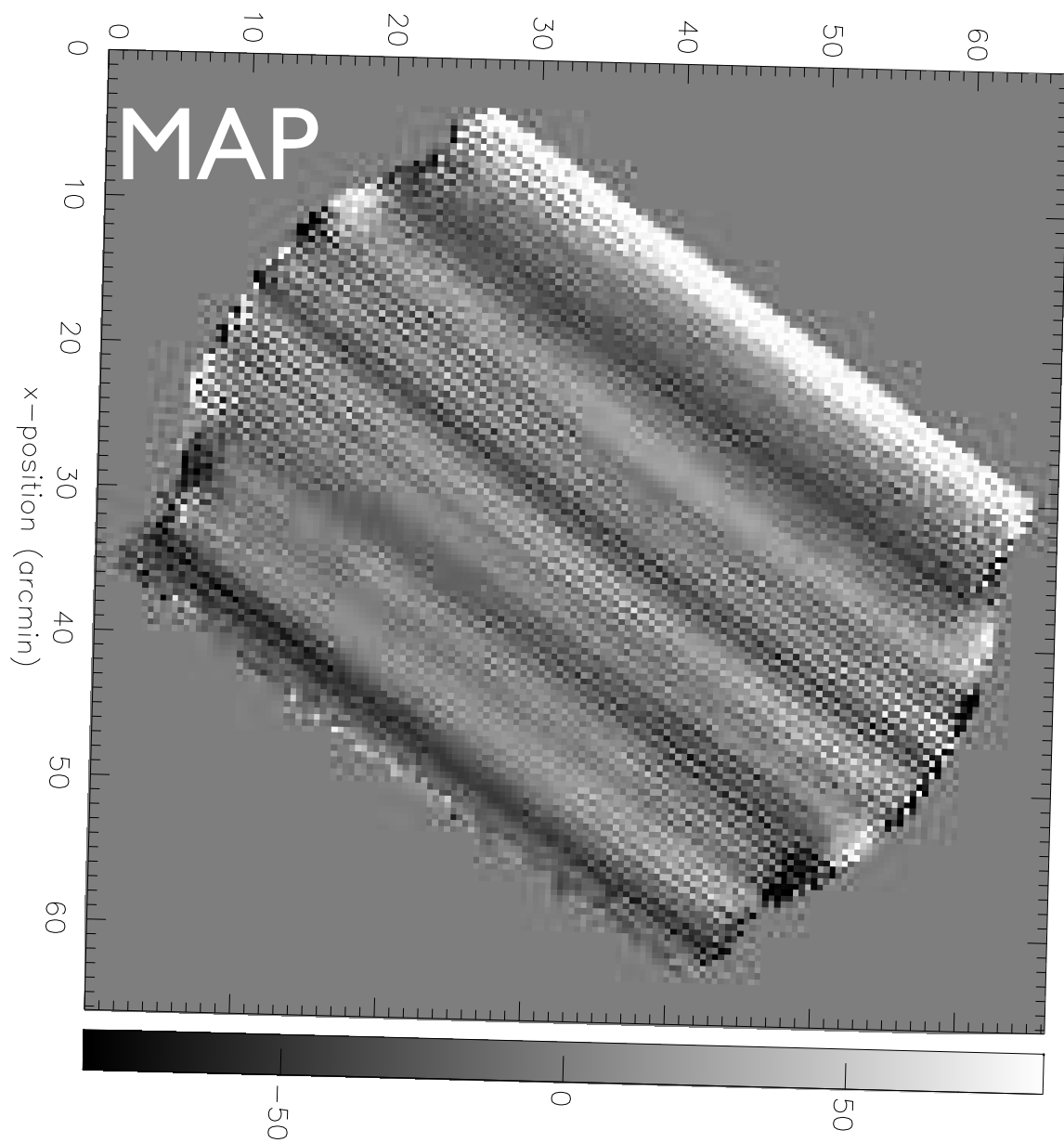


Power Spectra: Anisotropic Noise



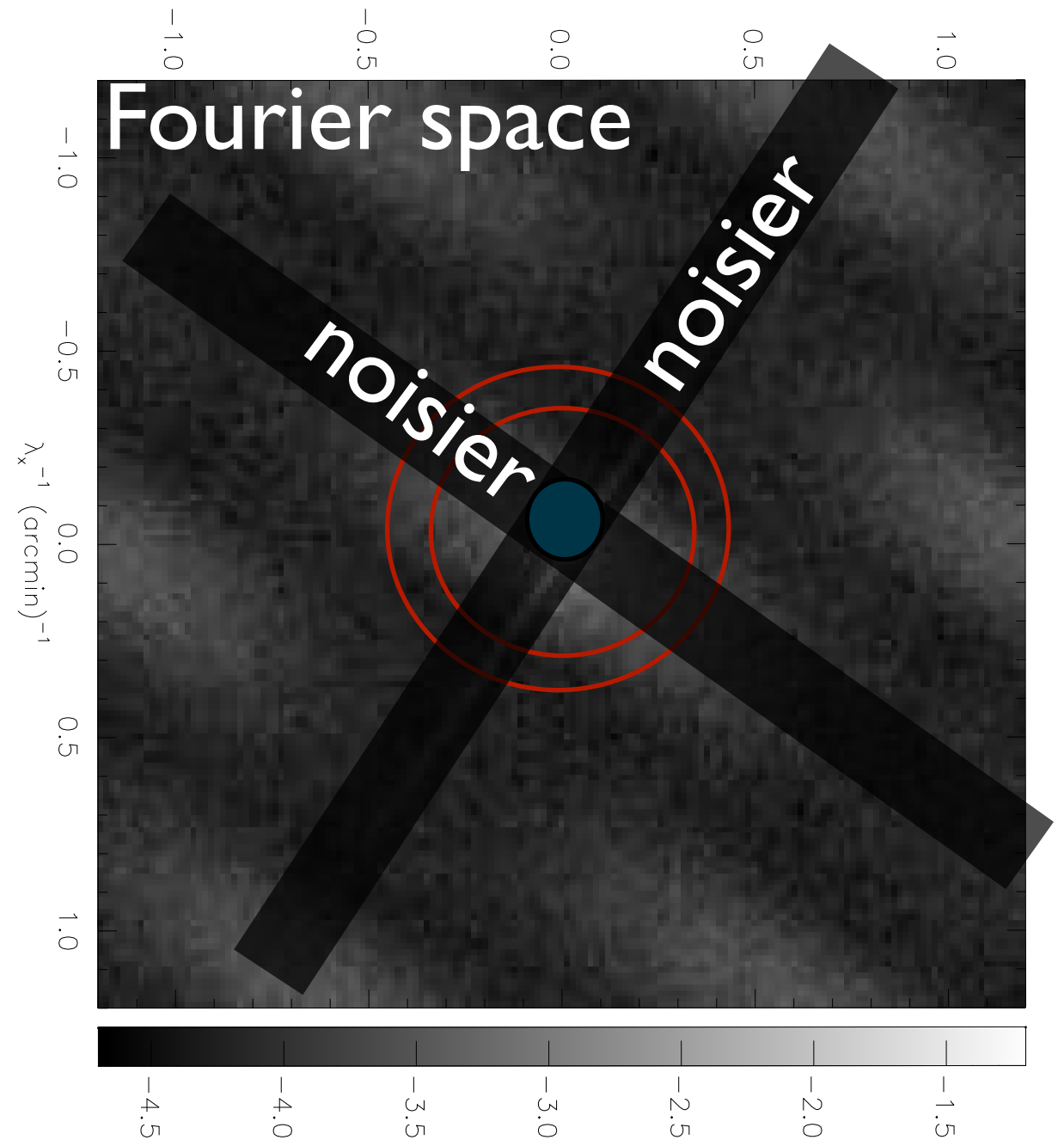
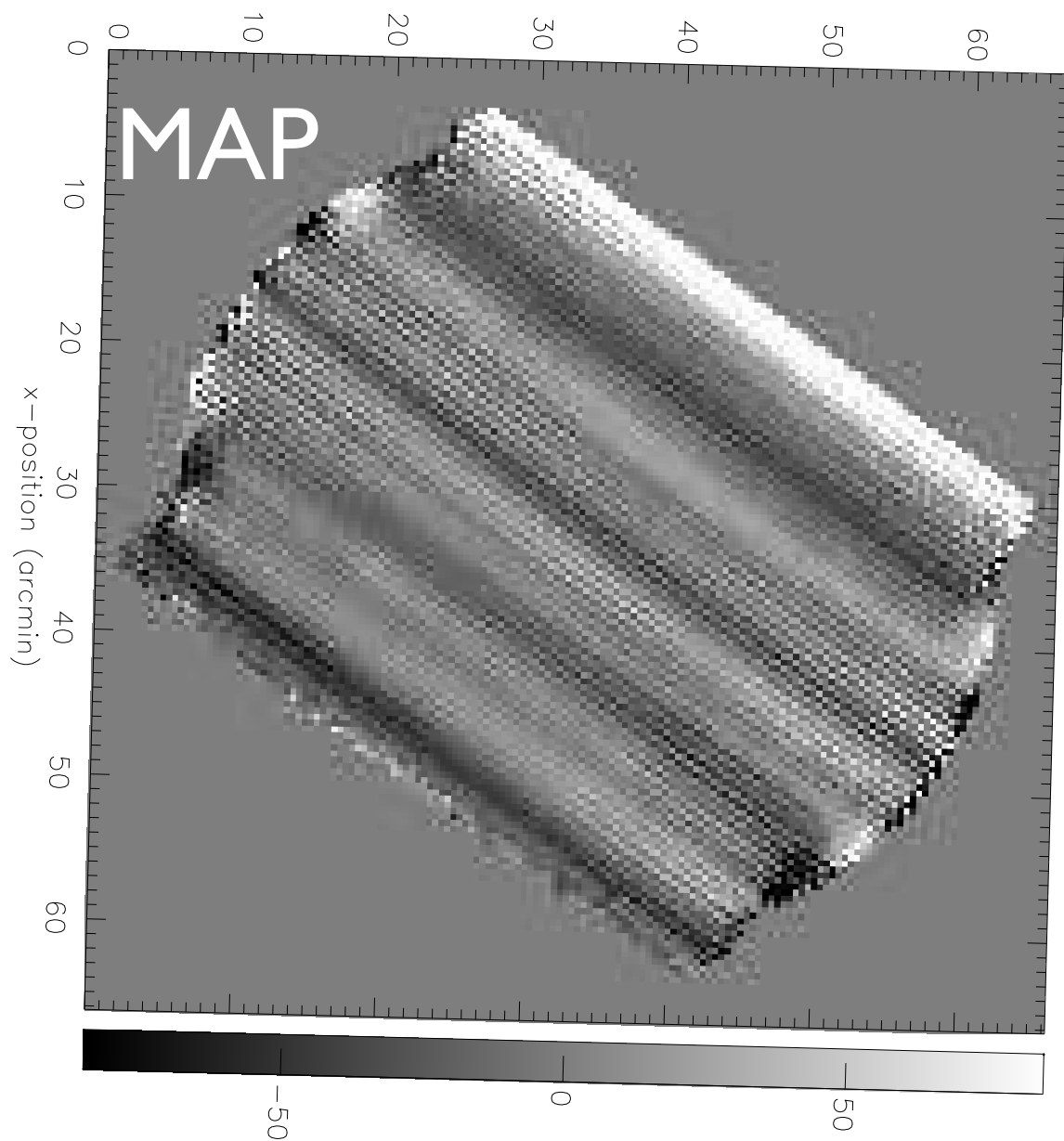
Cross Linking

Make a map with artifact at different angle to be able to reconstruct lost modes



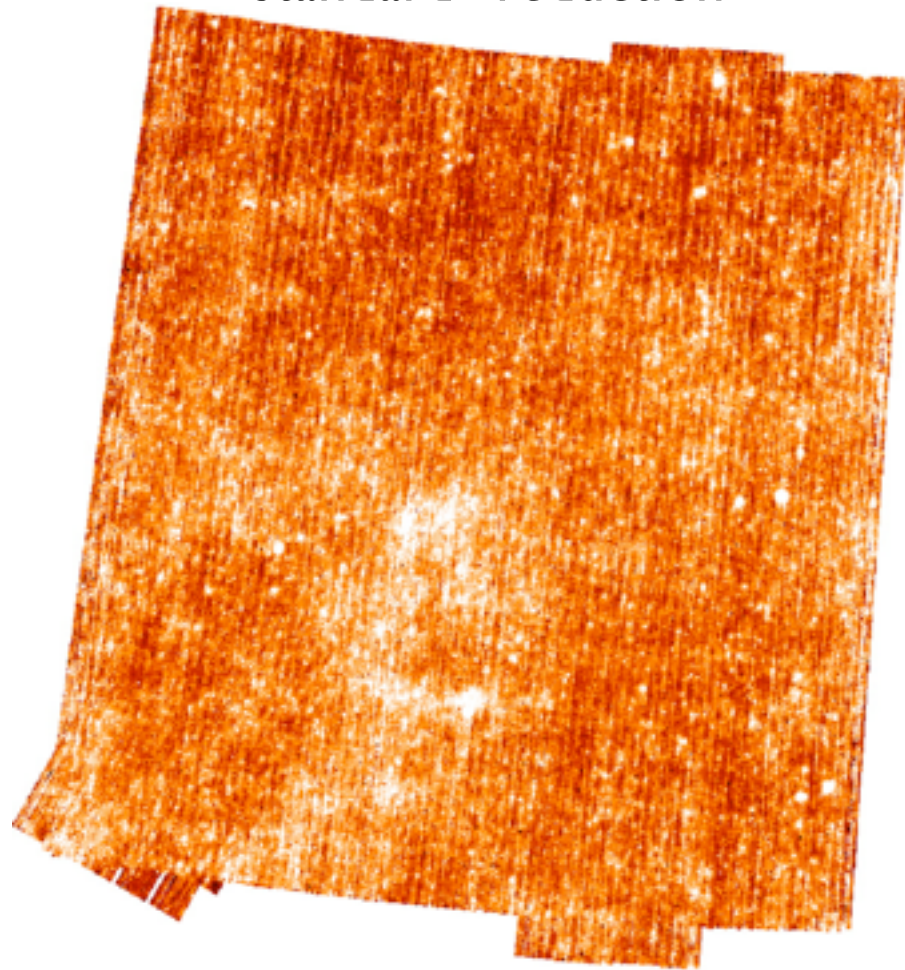
Cross Linking

Drawback: you now have $2N$ sort of noisy modes instead of N horrible ones

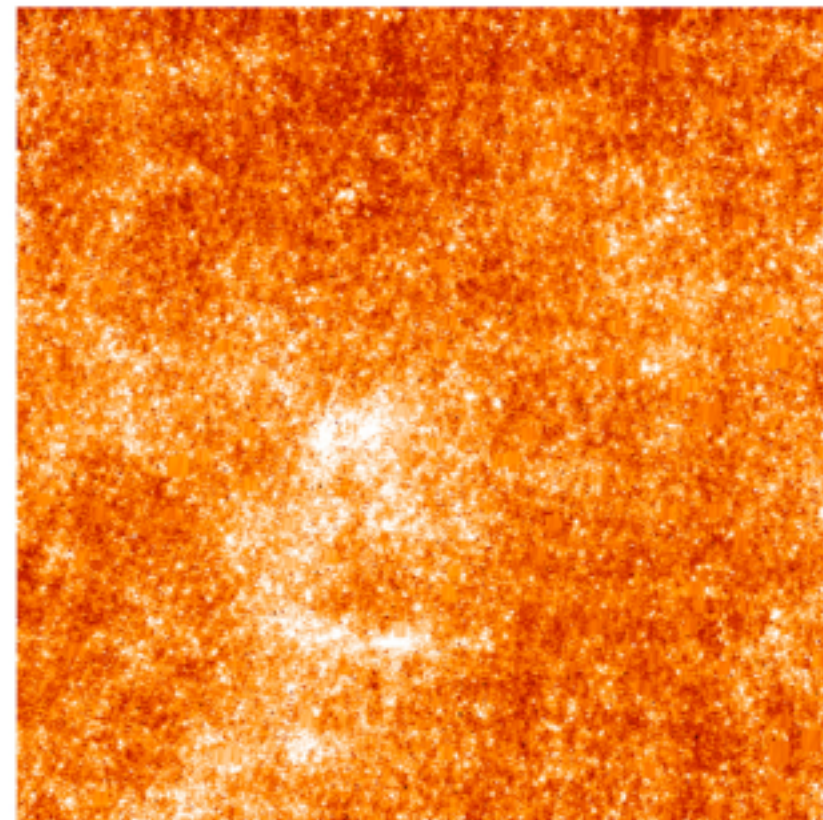


160 μm Maps

“standard” reduction



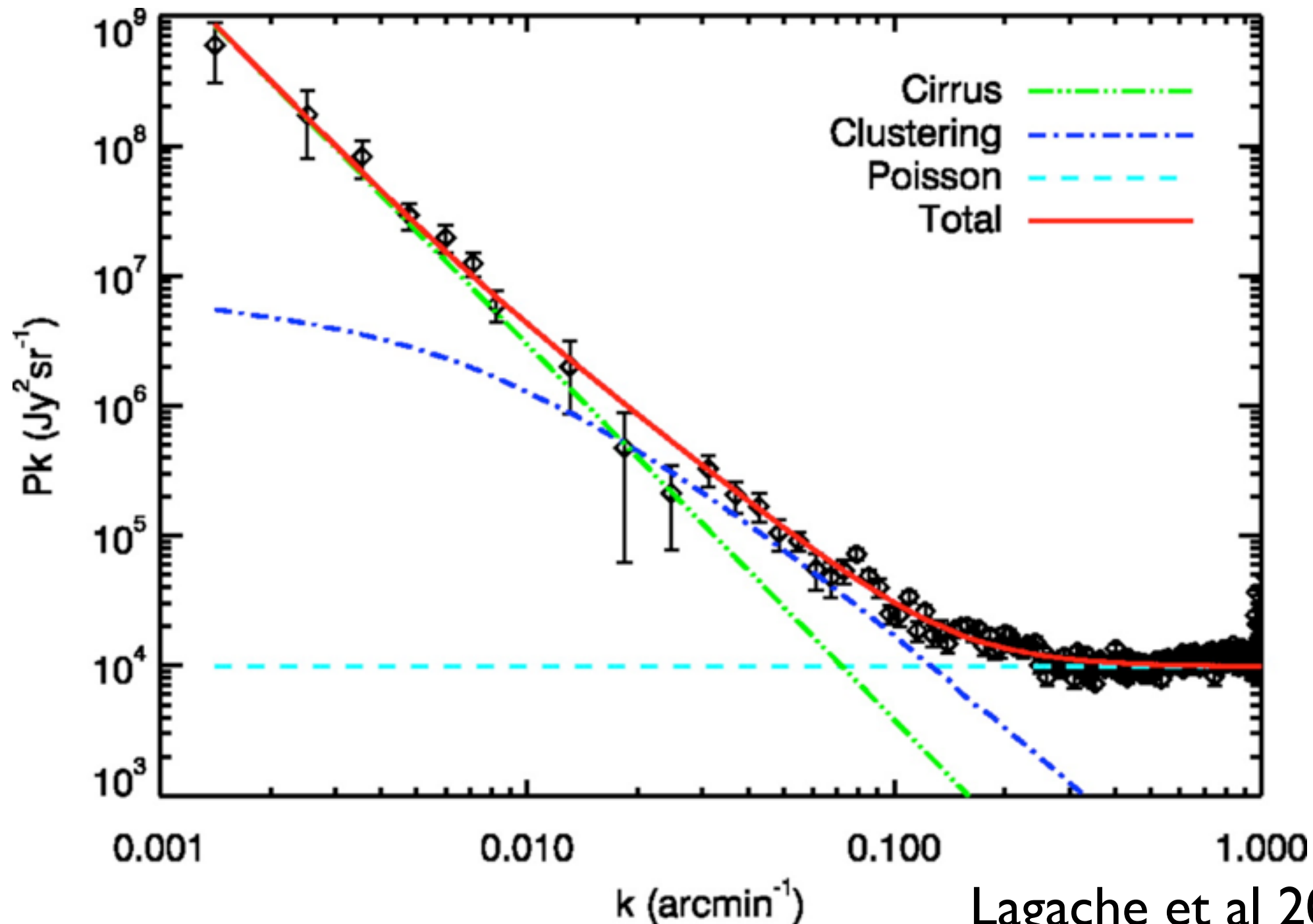
corrected map



10 deg²
with
Spitzer

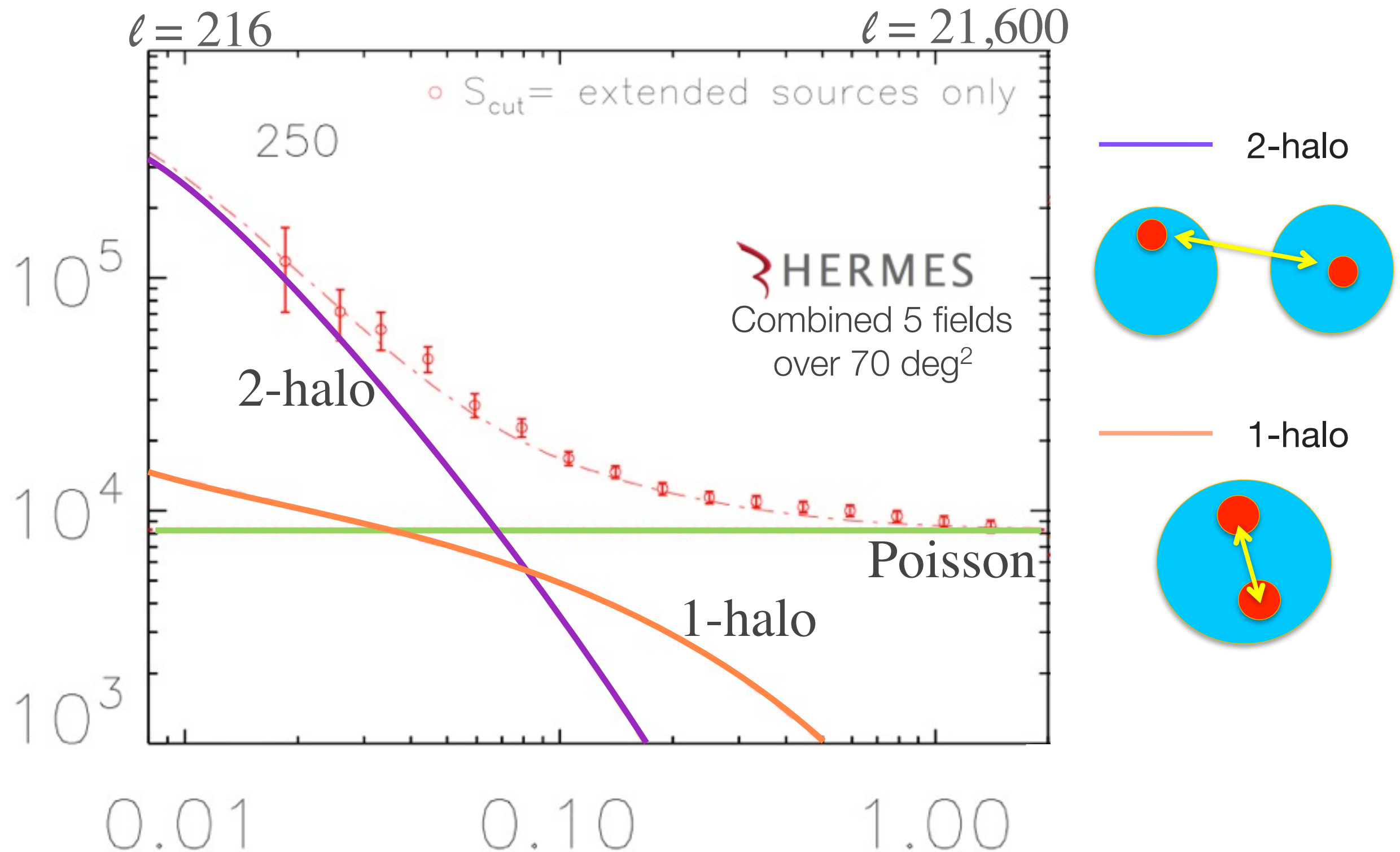
Lagache et al 2007

160 μ m Power Spectrum

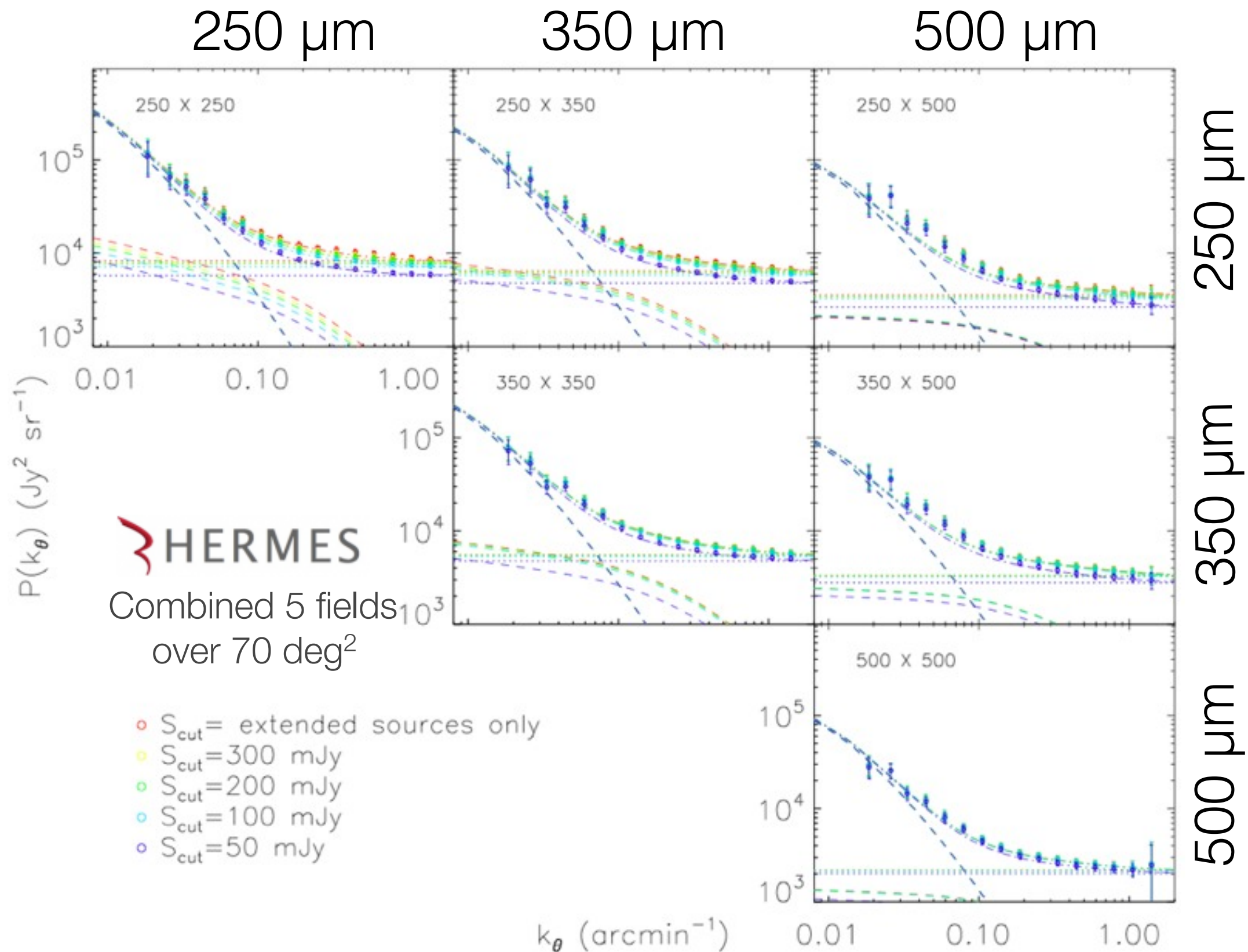


Lagache et al 2007

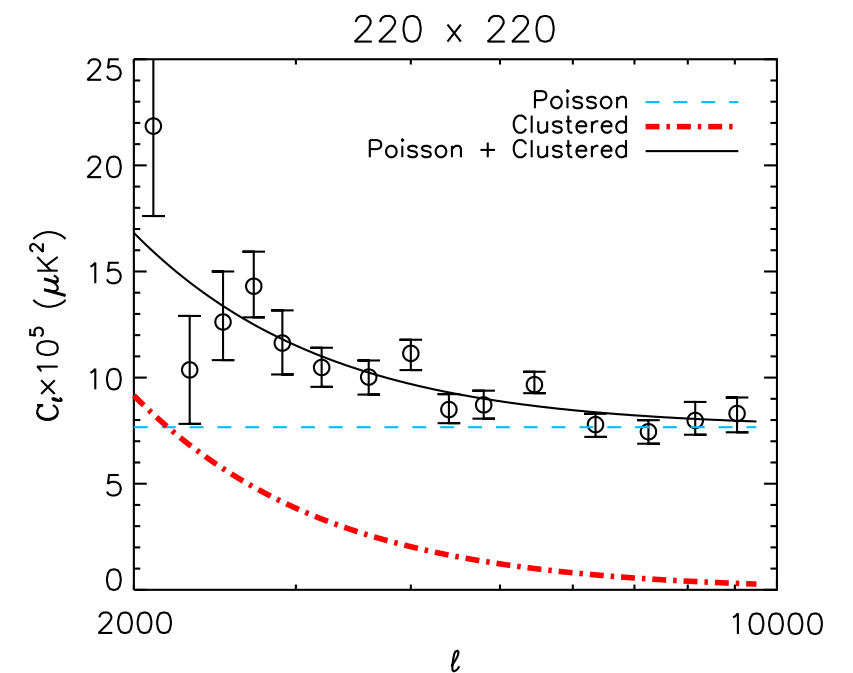
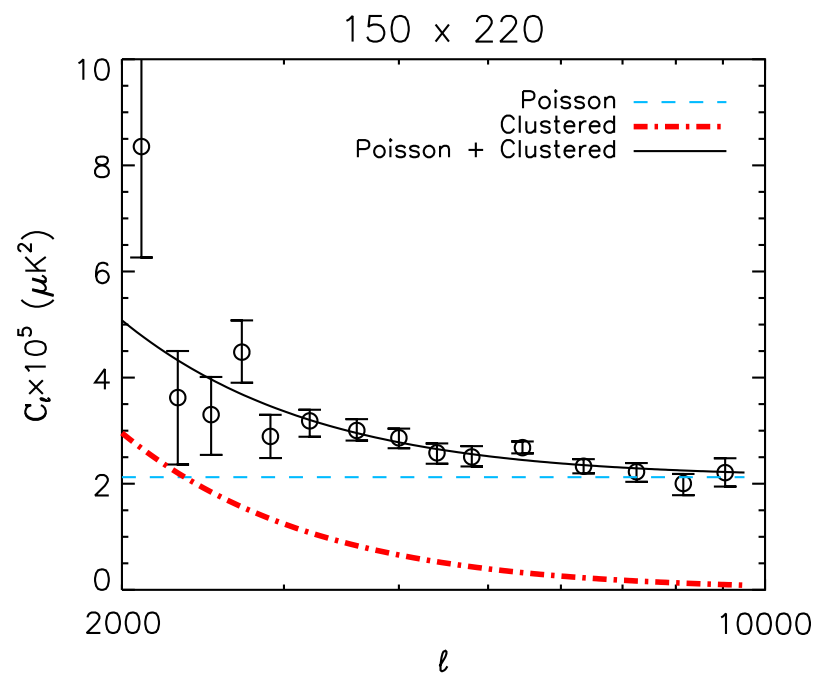
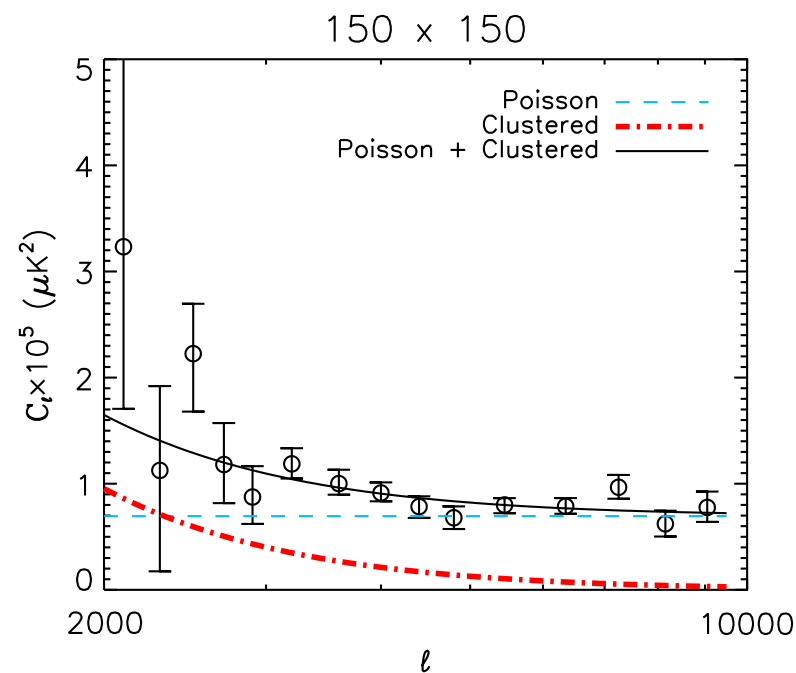
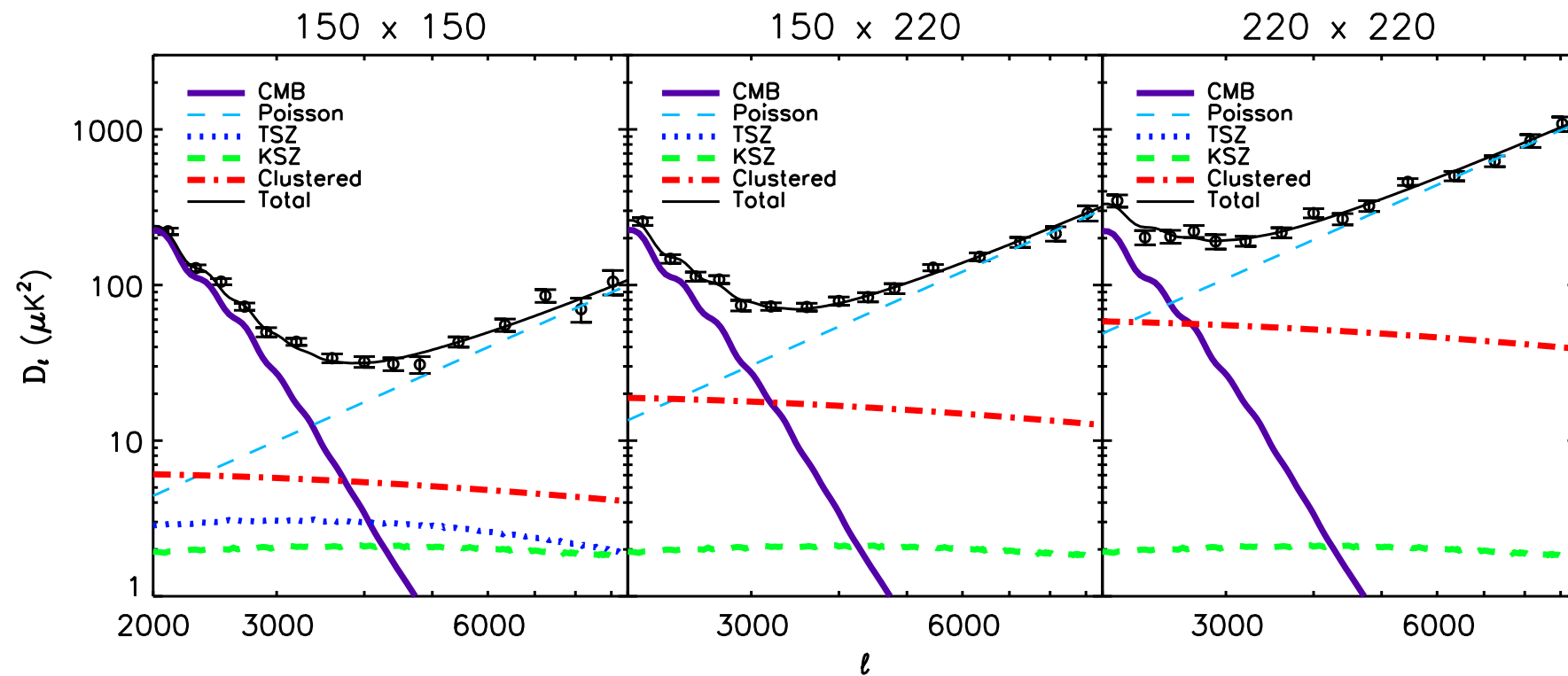
HerMES power spectra



HerMES power spectra

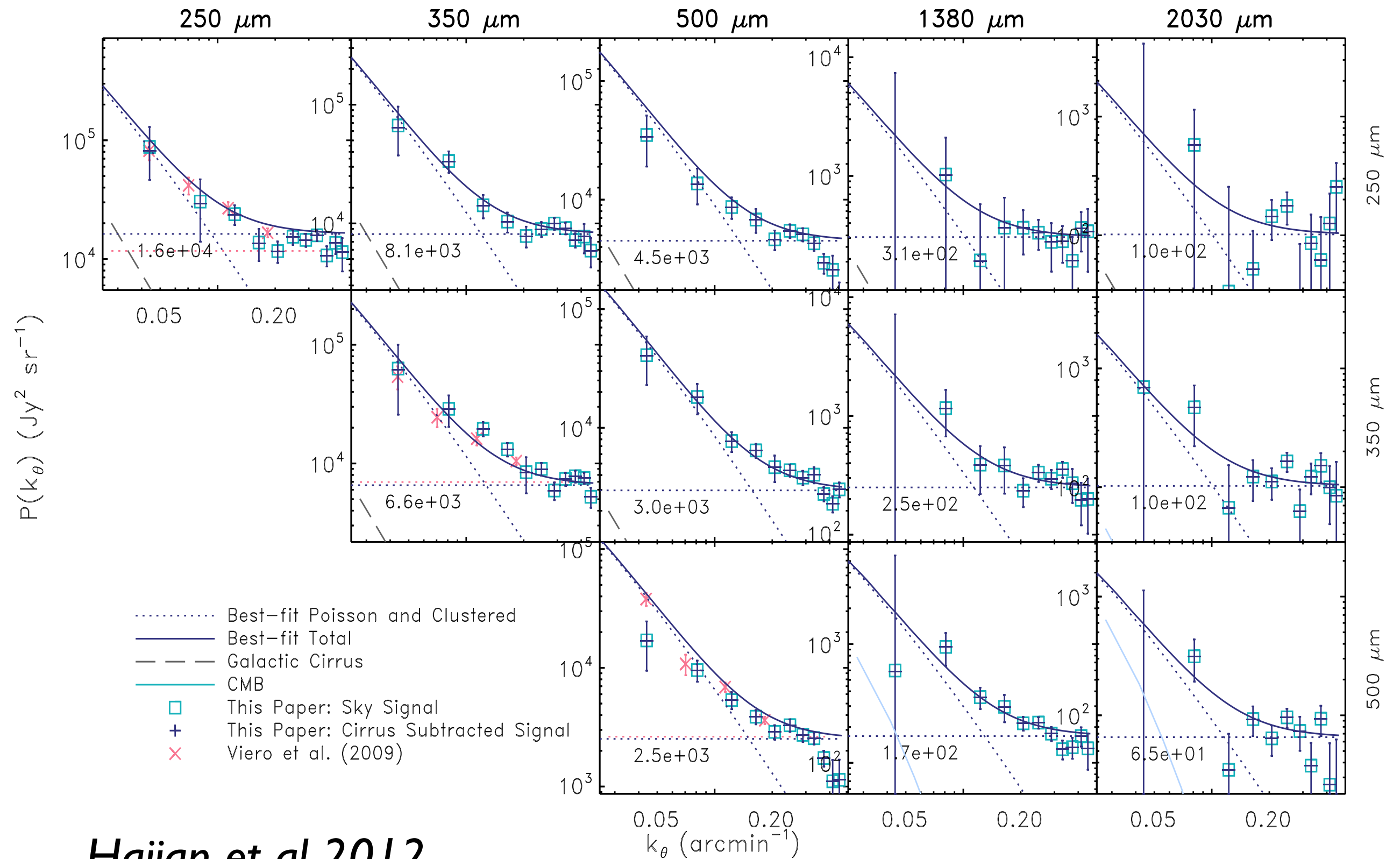


SPT CIB Detection



Hall et al 2010

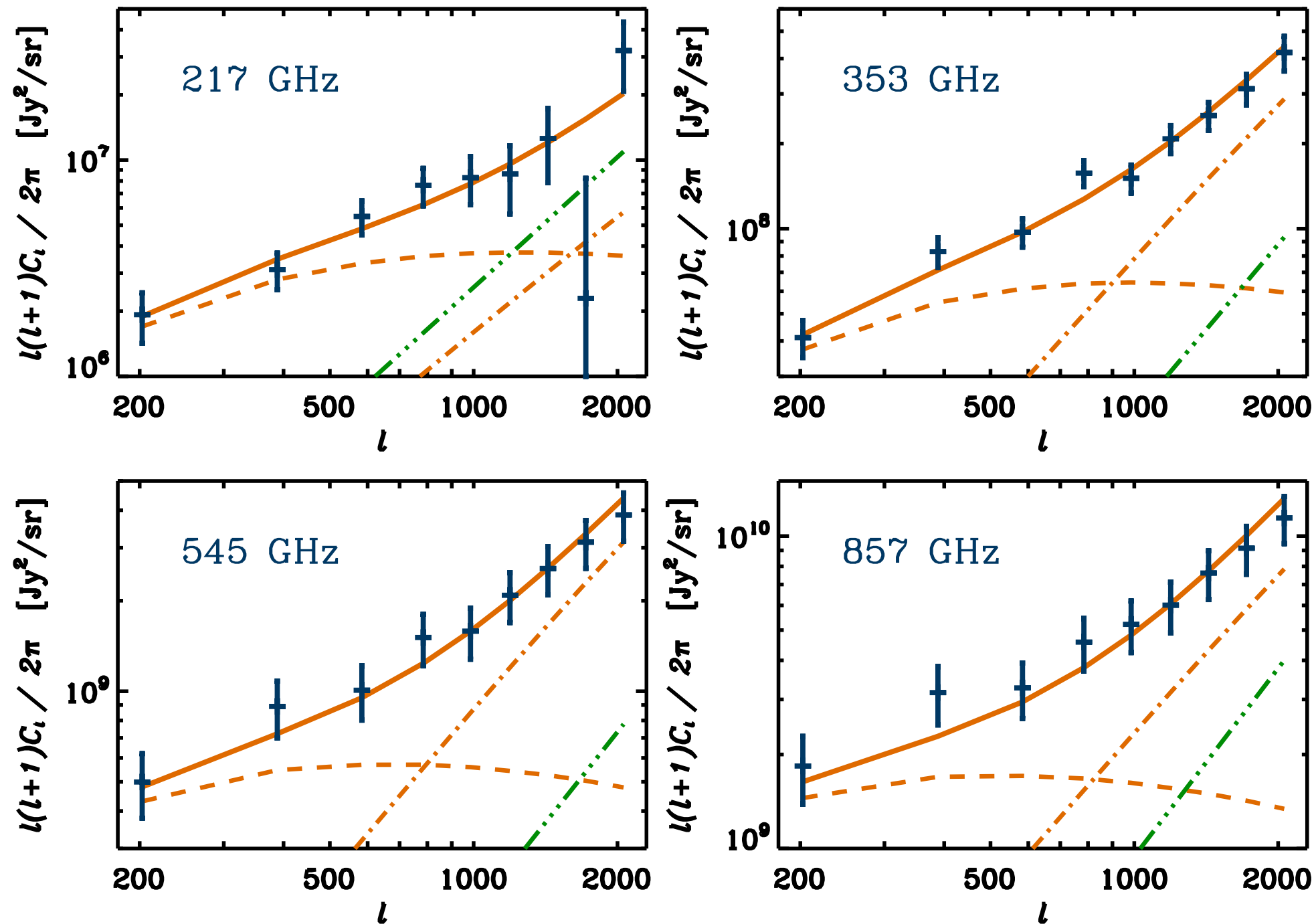
ACT X BLAST



Hajian et al 2012

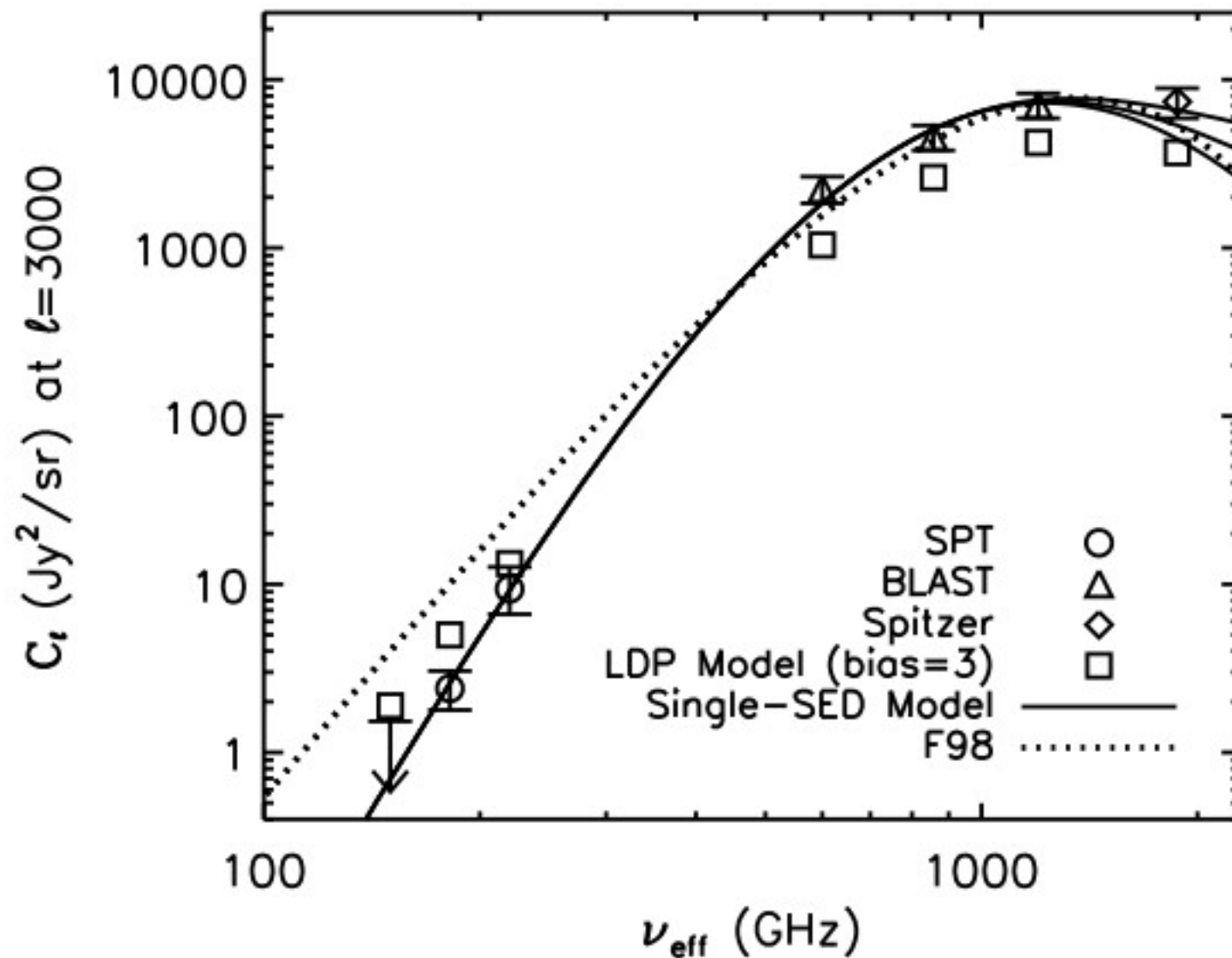
Planck CIB

Measurements



Planck collaboration 2012

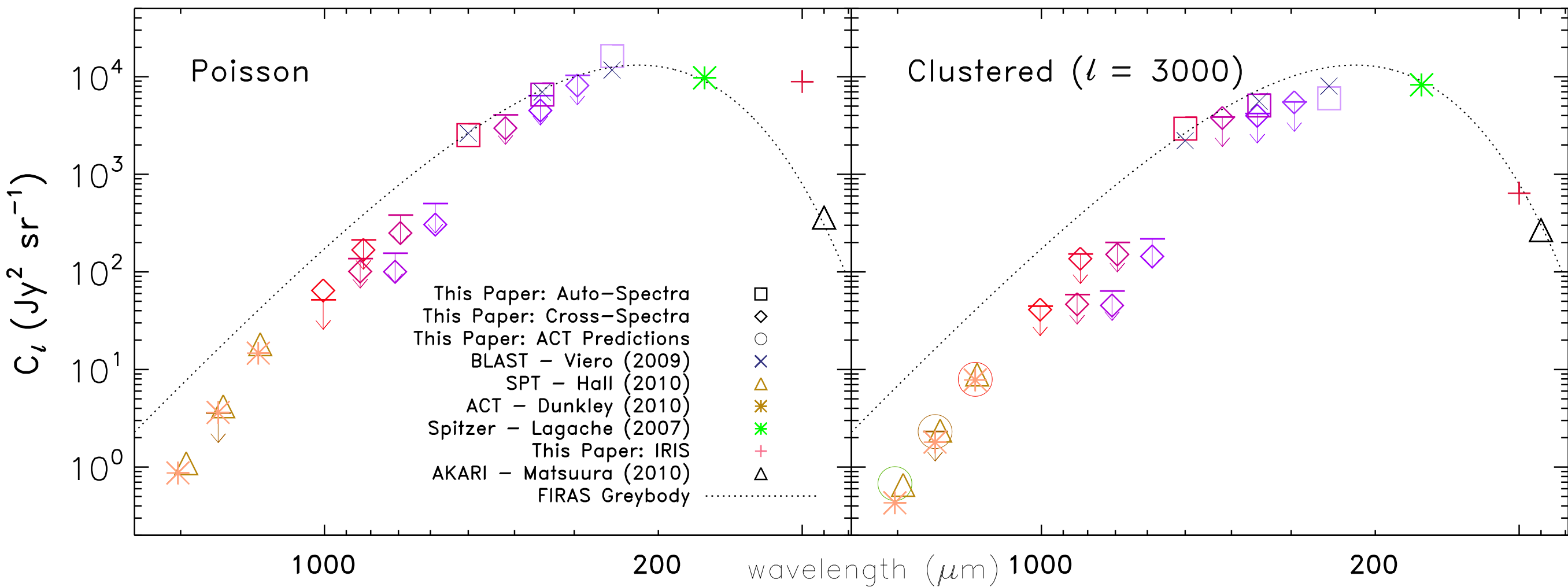
Frequency scaling of Dusty Galaxy Backgrounds



Single-SED model assumes all galaxies have same rest-frame properties ($T=34$ K, $\beta=2$) spread over a broad range in redshift (peaking at $z=2$)

First detection of clustered point source power from CIB sources in the mm bands

Frequency Scaling of the CIB



Summary

- lots more information in a map than just source counts (1 pt function)
- power spectrum of CIB map is providing measurements of clustering at many wavelengths (2 pt function)
 - SPT, ACT, Planck turn out to be CIB experiments